



$$I(J^P) = \frac{1}{2}(0^-)$$

## **D<sup>0</sup> MASS**

The fit includes  $D^\pm$ ,  $D^0$ ,  $D_s^\pm$ ,  $D^{*\pm}$ ,  $D^{*0}$ , and  $D_s^{*\pm}$  mass and mass difference measurements.

VALUE (MeV)	EVTS	DOCUMENT ID	TECN	COMMENT
<b>1864.5± 0.4 OUR FIT</b>		Error includes scale factor of 1.1.		
<b>1864.1± 1.0 OUR AVERAGE</b>				
1864.6± 0.3±1.0	641	BARLAG	90C ACCM	$\pi^-$ Cu 230 GeV
1852 ± 7	16	ADAMOVICH	87 EMUL	Photoproduction
1861 ± 4		DERRICK	84 HRS	$e^+ e^-$ 29 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
1856 ± 36	22	ADAMOVICH	84B EMUL	Photoproduction
1847 ± 7	1	FIORINO	81 EMUL	$\gamma N \rightarrow \bar{D}^0 +$
1863.8± 0.5		<sup>1</sup> SCHINDLER	81 MRK2	$e^+ e^-$ 3.77 GeV
1864.7± 0.6		<sup>1</sup> TRILLING	81 RVUE	$e^+ e^-$ 3.77 GeV
1863.0± 2.5	238	ASTON	80E OMEG	$\gamma p \rightarrow \bar{D}^0$
1860 ± 2	143	<sup>2</sup> AVERY	80 SPEC	$\gamma N \rightarrow D^{*+}$
1869 ± 4	35	<sup>2</sup> AVERY	80 SPEC	$\gamma N \rightarrow D^{*+}$
1854 ± 6	94	<sup>2</sup> ATIYA	79 SPEC	$\gamma N \rightarrow D^0 \bar{D}^0$
1850 ± 15	64	BALTAY	78C HBC	$\nu N \rightarrow K^0 \pi\pi$
1863 ± 3		GOLDHABER	77 MRK1	$D^0, D^+$ recoil spectra
1863.3± 0.9		<sup>1</sup> PERUZZI	77 MRK1	$e^+ e^-$ 3.77 GeV
1868 ± 11		PICCOLO	77 MRK1	$e^+ e^-$ 4.03, 4.41 GeV
1865 ± 15	234	GOLDHABER	76 MRK1	$K\pi$ and $K3\pi$

<sup>1</sup> PERUZZI 77 and SCHINDLER 81 errors do not include the 0.13% uncertainty in the absolute SPEAR energy calibration. TRILLING 81 uses the high precision  $J/\psi(1S)$  and  $\psi(2S)$  measurements of ZHOLENTZ 80 to determine this uncertainty and combines the PERUZZI 77 and SCHINDLER 81 results to obtain the value quoted. TRILLING 81 enters the fit in the  $D^\pm$  mass, and PERUZZI 77 and SCHINDLER 81 enter in the  $m_{D^\pm} - m_{D^0}$ , below.

<sup>2</sup> Error does not include possible systematic mass scale shift, estimated to be less than 5 MeV.

## **$m_{D^\pm} - m_{D^0}$**

The fit includes  $D^\pm$ ,  $D^0$ ,  $D_s^\pm$ ,  $D^{*\pm}$ ,  $D^{*0}$ , and  $D_s^{*\pm}$  mass and mass difference measurements.

VALUE (MeV)	DOCUMENT ID	TECN	COMMENT
<b>4.78±0.10 OUR FIT</b>	Error includes scale factor of 1.1.		
<b>4.74±0.28 OUR AVERAGE</b>			
4.7 ± 0.3	<sup>3</sup> SCHINDLER	81 MRK2	$e^+ e^-$ 3.77 GeV
5.0 ± 0.8	<sup>3</sup> PERUZZI	77 MRK1	$e^+ e^-$ 3.77 GeV

<sup>3</sup> See the footnote on TRILLING 81 in the  $D^0$  and  $D^\pm$  sections on the mass.

## $D^0$ MEAN LIFE

Measurements with an error  $> 10 \times 10^{-15}$  s have been omitted from the average.

VALUE ( $10^{-15}$ s)	EVTS	DOCUMENT ID	TECN	COMMENT
<b>410.1 ± 1.5 OUR AVERAGE</b>				
409.6 ± 1.1 ± 1.5	210k	LINK	02F FOCS	$\gamma$ nucleus, $\approx 180$ GeV
407.9 ± 6.0 ± 4.3	10k	KUSHNIR...	01 SELX	$K^- \pi^+, K^- \pi^+ \pi^+ \pi^-$
413 ± 3 ± 4	35k	AITALA	99E E791	$K^- \pi^+$
408.5 ± 4.1 ± 3.5	25k	BONVICINI	99 CLE2	$e^+ e^- \approx \Upsilon(4S)$
413 ± 4 ± 3	16k	FABRETTI	94D E687	$K^- \pi^+, K^- \pi^+ \pi^+ \pi^-$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
424 ± 11 ± 7	5118	FABRETTI	91 E687	$K^- \pi^+, K^- \pi^+ \pi^+ \pi^-$
417 ± 18 ± 15	890	ALVAREZ	90 NA14	$K^- \pi^+, K^- \pi^+ \pi^+ \pi^-$
388 ± 23 ± 21	641	<sup>4</sup> BARLAG	90C ACCM	$\pi^-$ Cu 230 GeV
480 ± 40 ± 30	776	ALBRECHT	88I ARG	$e^+ e^-$ 10 GeV
422 ± 8 ± 10	4212	RAAB	88 E691	Photoproduction
420 ± 50	90	BARLAG	87B ACCM	$K^-$ and $\pi^-$ 200 GeV

<sup>4</sup> BARLAG 90C estimate systematic error to be negligible.

## $D^0$ - $\bar{D}^0$ MIXING

Revised January 2006 by D. Asner (Carleton University)

Standard Model contributions to  $D^0$ - $\bar{D}^0$  mixing are strongly suppressed by CKM and GIM factors. Thus the observation of  $D^0$ - $\bar{D}^0$  mixing might be evidence for physics beyond the Standard Model. See Burdman and Shipsey [1] for a review of  $D^0$ - $\bar{D}^0$  mixing, Ref. [2] for a compilation of mixing predictions, and Ref. [3] for later predictions.

**Formalism:** The time evolution of the  $D^0$ - $\bar{D}^0$  system is described by the Schrödinger equation

$$i \frac{\partial}{\partial t} \begin{pmatrix} D^0(t) \\ \bar{D}^0(t) \end{pmatrix} = \left( \mathbf{M} - \frac{i}{2} \boldsymbol{\Gamma} \right) \begin{pmatrix} D^0(t) \\ \bar{D}^0(t) \end{pmatrix}, \quad (1)$$

where the  $\mathbf{M}$  and  $\boldsymbol{\Gamma}$  matrices are Hermitian, and *CPT* invariance requires that  $M_{11} = M_{22} \equiv M$  and  $\Gamma_{11} = \Gamma_{22} \equiv \Gamma$ . The off-diagonal elements of these matrices describe the dispersive and absorptive parts of  $D^0$ - $\bar{D}^0$  mixing.

The two eigenstates  $D_1$  and  $D_2$  of the effective Hamiltonian matrix  $(\mathbf{M} - \frac{i}{2}\Gamma)$  are given by

$$|D_{1,2}\rangle = p|D^0\rangle \pm q|\overline{D}^0\rangle. \quad (2)$$

The corresponding eigenvalues are

$$\lambda_{1,2} \equiv m_{1,2} - \frac{i}{2}\Gamma_{1,2} = \left(M - \frac{i}{2}\Gamma\right) \pm \frac{q}{p} \left(M_{12} - \frac{i}{2}\Gamma_{12}\right), \quad (3)$$

where  $m_1$  and  $\Gamma_1$  are the mass and width of the  $D_1$ , etc., and

$$\left|\frac{q}{p}\right|^2 = \frac{M_{12}^* - \frac{i}{2}\Gamma_{12}^*}{M_{12} - \frac{i}{2}\Gamma_{12}}. \quad (4)$$

We define reduced mixing amplitudes  $x$  and  $y$  by

$$x \equiv 2M_{12}/\Gamma = (m_1 - m_2)/\Gamma = \Delta m/\Gamma \quad (5)$$

and

$$y \equiv \Gamma_{12}/\Gamma = (\Gamma_1 - \Gamma_2)/2\Gamma = \Delta\Gamma/2\Gamma, \quad (6)$$

where  $\Gamma \equiv (\Gamma_1 + \Gamma_2)/2$ . The mixing rate,  $R_M$ , is approximately  $(x^2 + y^2)/2$ . In Eq. (5) and Eq. (6), the middle relation holds only in the limit of  $CP$  conservation, in which case the subscripts 1 and 2 denote the  $CP$ -even and  $CP$ -odd eigenstates.

The parameters  $x$  and  $y$  are measured in several ways. The most precise constraints are obtained using the time-dependence of  $D$  decays. Since  $D^0$ - $\overline{D}^0$  mixing is a small effect, the identification tag of the initial particle as a  $D^0$  or a  $\overline{D}^0$  must be extremely accurate. The usual tag is the charge of the distinctive slow pion in the decay sequence  $D^{*+} \rightarrow D^0\pi^+$  or  $D^{*-} \rightarrow \overline{D}^0\pi^-$ . In current experiments, the probability of mistagging is about 0.1%. Another tag of comparable accuracy is identification of one of the  $D$ 's produced from  $\psi(3770) \rightarrow D^0\overline{D}^0$ . Time-dependent analyses are not possible at symmetric

charm threshold facilities (the  $D^0$  and  $\bar{D}^0$  do not travel far enough). However, the quantum coherent  $D^0\bar{D}^0$   $C = -1$  state provides time-integrated sensitivity [4, 5].

**Time-Dependent Analyses:** We extend the formalism of this *Review*'s note on “ $B^0-\bar{B}^0$  Mixing” [6]. In addition to the “right-sign” instantaneous decay amplitudes  $\bar{A}_f \equiv \langle f | H | \bar{D}^0 \rangle$  and  $A_{\bar{f}} \equiv \langle \bar{f} | H | D^0 \rangle$  for  $CP$  conjugate final states  $f$  and  $\bar{f}$ , we include the “wrong-sign” amplitudes  $\bar{A}_{\bar{f}} \equiv \langle \bar{f} | H | \bar{D}^0 \rangle$  and  $A_f \equiv \langle f | H | D^0 \rangle$ .

It is usual to normalize the wrong-sign decay distributions to the integrated rate of right-sign decays and to express time in units of the precisely measured  $D^0$  mean lifetime,  $\bar{\tau}_{D^0} = 1/\Gamma = 2/(\Gamma_1 + \Gamma_2)$ . Starting from a pure  $|D^0\rangle$  or  $|\bar{D}^0\rangle$  state at  $t = 0$ , the time-dependent rates of production of the wrong-sign final states relative to the integrated right-sign states are then

$$r(t) = \frac{|\langle f | H | D^0(t) \rangle|^2}{|\bar{A}_f|^2} = \left| \frac{q}{p} \right|^2 \left| g_+(t) \chi_f^{-1} + g_-(t) \right|^2 \quad (7)$$

and

$$\bar{r}(t) = \frac{|\langle \bar{f} | H | \bar{D}^0(t) \rangle|^2}{|A_{\bar{f}}|^2} = \left| \frac{p}{q} \right|^2 \left| g_+(t) \chi_{\bar{f}} + g_-(t) \right|^2, \quad (8)$$

where

$$\chi_f \equiv q\bar{A}_f/pA_f, \quad \chi_{\bar{f}} \equiv q\bar{A}_{\bar{f}}/pA_{\bar{f}}, \quad (9)$$

and

$$g_{\pm}(t) = \frac{1}{2} (e^{-iz_1 t} \pm e^{-iz_2 t}), \quad z_{1,2} = \frac{\lambda_{1,2}}{\Gamma}. \quad (10)$$

Note that a change in the convention for the relative phase of  $D^0$  and  $\bar{D}^0$  would cancel between  $q/p$  and  $\bar{A}_f/A_f$  and leave  $\chi_f$  invariant.

We expand  $r(t)$  and  $\bar{r}(t)$  to second order in time for modes where the ratio of decay amplitudes  $R_D = |A_f/\bar{A}_f|^2$  is very small.

**Semileptonic decays:** In semileptonic  $D$  decays,  $A_f = \bar{A}_f = 0$  in the Standard Model. Then in the limit of weak mixing, where  $|ix + y| \ll 1$ ,  $r(t)$  is given by

$$r(t) = |g_-(t)|^2 \left| \frac{q}{p} \right|^2 \approx \frac{e^{-t}}{4} (x^2 + y^2) t^2 \left| \frac{q}{p} \right|^2. \quad (11)$$

For  $\bar{r}(t)$  one replaces  $q/p$  here with  $p/q$ . In the limit of  $CP$  conservation,  $r(t) = \bar{r}(t)$ , and the time-integrated mixing rate relative to the time-integrated right-sign decay rate is

$$R_M = \int_0^\infty r(t) dt = \left| \frac{q}{p} \right|^2 \frac{x^2 + y^2}{2 + x^2 - y^2} \approx \frac{1}{2} (x^2 + y^2). \quad (12)$$

Table 1 summarizes results from semileptonic decays.

**Table 1:** Results for  $R_M$  in  $D^0$  semileptonic decays.

Year	Exper.	Final state(s)	$R_M$ (90 (95)% C.L.)
2005	Belle <sup>a</sup>	$K^{(*)+} e^- \bar{\nu}_e$	$< 1.0 \times 10^{-3}$
2005	CLEO <sup>b</sup>	$K^{(*)+} e^- \bar{\nu}_e$	$< 7.8 \times 10^{-3}$
2004	BABAR <sup>c</sup>	$K^{(*)+} e^- \bar{\nu}_e$	$< 4.2(4.6) \times 10^{-3}$
2002	FOCUS [7]	$K^+ \mu^- \bar{\nu}_\mu$	$< 1.01(1.31) \times 10^{-3}$
1996	E791 <sup>d</sup>	$K^+ \ell^- \bar{\nu}_\ell$	$< 5.0 \times 10^{-3}$

See the end of the  $D^0$  listings for these references: <sup>a</sup>BITENC 05, <sup>b</sup>CAWLFIELD 05, <sup>c</sup>AUBERT 04, <sup>d</sup>AITALA 96C.

**Wrong-sign decays to hadronic non- $CP$  eigenstates:**

Consider the final state  $f = K^+ \pi^-$ , where  $A_f$  is doubly Cabibbo-suppressed. The ratio of decay amplitudes is

$$\frac{A_f}{\bar{A}_f} = -\sqrt{R_D} e^{-i\delta}, \quad \left| \frac{A_f}{\bar{A}_f} \right| \sim O(\tan^2 \theta_c), \quad (13)$$

where  $R_D$  is the doubly Cabibbo-suppressed (DCS) decay rate relative to the Cabibbo-favored (CF) rate, the minus sign originates from the sign of  $V_{us}$  relative to  $V_{cd}$ , and  $\delta$  is the phase difference between DCS and CF processes not attributed to the first-order electroweak spectator diagram.

We characterize the violation of  $CP$  in the mixing amplitude, the decay amplitude, and the interference between mixing and decay, by real-valued parameters  $A_M$ ,  $A_D$ , and  $\phi$ . We adopt a parametrization similar to that of Nir [8] and CLEO [GODANG 00] and express these quantities in a way that is convenient to describe the three types of  $CP$  violation:

$$\left| \frac{q}{p} \right| = 1 + A_M, \quad (14)$$

$$\chi_f^{-1} \equiv \frac{p A_f}{q \bar{A}_f} = \frac{-\sqrt{R_D}(1 + A_D)}{(1 + A_M)} e^{-i(\delta + \phi)}, \quad (15)$$

$$\chi_{\bar{f}} \equiv \frac{q \bar{A}_{\bar{f}}}{p A_{\bar{f}}} = \frac{-\sqrt{R_D}(1 + A_M)}{(1 + A_D)} e^{-i(\delta - \phi)}. \quad (16)$$

In general,  $\chi_{\bar{f}}$  and  $\chi_f^{-1}$  are independent complex numbers. To leading order,

$$\begin{aligned} r(t) = e^{-t} \times & \left[ R_D(1 + A_D)^2 \right. \\ & \left. + \sqrt{R_D}(1 + A_M)(1 + A_D)y'_- t + \frac{(1 + A_M)^2 R_M}{2} t^2 \right] \end{aligned} \quad (17)$$

and

$$\begin{aligned} \bar{r}(t) = & e^{-t} \times \left[ \frac{R_D}{(1+A_D)^2} \right. \\ & \left. + \frac{\sqrt{R_D}}{(1+A_D)(1+A_M)} y'_+ t + \frac{R_M}{2(1+A_M)^2} t^2 \right]. \end{aligned} \quad (18)$$

Here

$$y'_\pm \equiv y' \cos \phi \pm x' \sin \phi = y \cos(\delta \mp \phi) - x \sin(\delta \mp \phi) \quad (19)$$

$$y' \equiv y \cos \delta - x \sin \delta, \quad x' \equiv x \cos \delta + y \sin \delta, \quad (20)$$

and  $R_M$  is the mixing rate relative to the time-integrated right-sign rate.

The three terms in Eq. (17) and Eq. (18) probe the three fundamental types of  $CP$  violation. In the limit of  $CP$  conservation,  $A_M$ ,  $A_D$ , and  $\phi$  are all zero, and then

$$r(t) = \bar{r}(t) = e^{-t} \left( R_D + \sqrt{R_D} y' t + \frac{1}{2} R_M t^2 \right), \quad (21)$$

and the time-integrated wrong-sign rate relative to the integrated right-sign rate is

$$R = \int_0^\infty r(t) dt = R_D + \sqrt{R_D} y' + R_M. \quad (22)$$

The ratio  $R$  is the most readily accessible experimental quantity. Table 2 gives recent measurements of  $R$  in  $D^0 \rightarrow K^+ \pi^-$  decay. The average of these results,  $R = (0.376 \pm 0.009)\%$ , is about two standard deviations from the average of earlier, less precise results,  $R = (0.81 \pm 0.23)\%$ , which we have omitted.

**Table 2:** Results for  $R$  in  $D^0 \rightarrow K^+ \pi^-$ .

Year	Exper.	Technique	$R(\times 10^{-3})$	$A_D(\%)$
2006	Belle <sup>a</sup>	$e^+ e^- \rightarrow \Upsilon(4S)$	$3.77 \pm 0.08 \pm 0.05$	—
2005	FOCUS <sup>b</sup>	$\gamma$ BeO	$4.29 \pm 0.63 \pm 0.28$	$18.0 \pm 14.0 \pm 4.1$
2003	BABAR <sup>c</sup>	$e^+ e^- \rightarrow \Upsilon(4S)$	$3.57 \pm 0.22 \pm 0.27$	$9.5 \pm 6.1 \pm 8.3$
2000	CLEO <sup>d</sup>	$e^+ e^- \rightarrow \Upsilon(4S)$	$3.32_{-0.65}^{+0.63} \pm 0.40$	$2_{-20}^{+19} \pm 1$

See the end of the  $D^0$  listings for these references: <sup>a</sup>ZHANG 06,  
<sup>b</sup>LINK 05, <sup>c</sup>AUBERT 03Z, <sup>d</sup>GODANG 00.

**Table 3:** Results from studies of the time dependence  $r(t)$ .

Year	Exper.	$y'$ (95% C.L.)	$x'^2/2$ (95% C.L.)
2006	Belle <sup>a</sup>	$-2.8 < y' < 2.1$ %	$< 0.036$ %
2005	FOCUS <sup>b</sup>	$-11.2 < y' < 6.7$ %	$< 0.40$ %
2003	BABAR <sup>c</sup>	$-5.6 < y' < 3.9$ %	$< 0.11$ %
2000	CLEO <sup>d</sup>	$-5.8 < y' < 1.0$ %	$< 0.041$ %

See the end of the  $D^0$  listings for these references: <sup>a</sup>ZHANG 06,  
<sup>b</sup>LINK 05, <sup>c</sup>AUBERT 03Z, <sup>d</sup>GODANG 00.

The contributions to  $R$ —allowing for  $CP$  violation—can be extracted by fitting the  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^- \pi^+$  decay rates. Table 2 gives the constraints on  $A_D$  with  $x' = y' = 0$ . Table 3 summarizes the results for  $y'$  and  $x'^2/2$ . Figure 1 shows the two-dimensional allowed regions. No meaningful constraints on  $A_M$  and  $\phi$  have been reported.

Extraction of the amplitudes  $x$  and  $y$  from the results in Table 3 requires knowledge of the relative strong phase  $\delta$ , a subject of theoretical discussion [4,9–11]. In most cases, it appears difficult for theory to accommodate  $\delta > 25^\circ$ , although the judicious placement of a  $K\pi$  resonance could allow  $\delta$  to be as large as  $40^\circ$ .

A quantum interference effect that provides useful sensitivity to  $\delta$  arises in the decay chain  $\psi(3770) \rightarrow D^0 \bar{D}^0 \rightarrow (f_{cp})(K^+ \pi^-)$ , where  $f_{cp}$  denotes a  $CP$  eigenstate from  $D^0$  decay, such as  $K^+ K^-$  [1, 16]. Here, the amplitude triangle relation

$$\sqrt{2} A(D_\pm \rightarrow K^- \pi^+) = A(D^0 \rightarrow K^- \pi^+) \pm A(\bar{D}^0 \rightarrow K^- \pi^+), \quad (23)$$

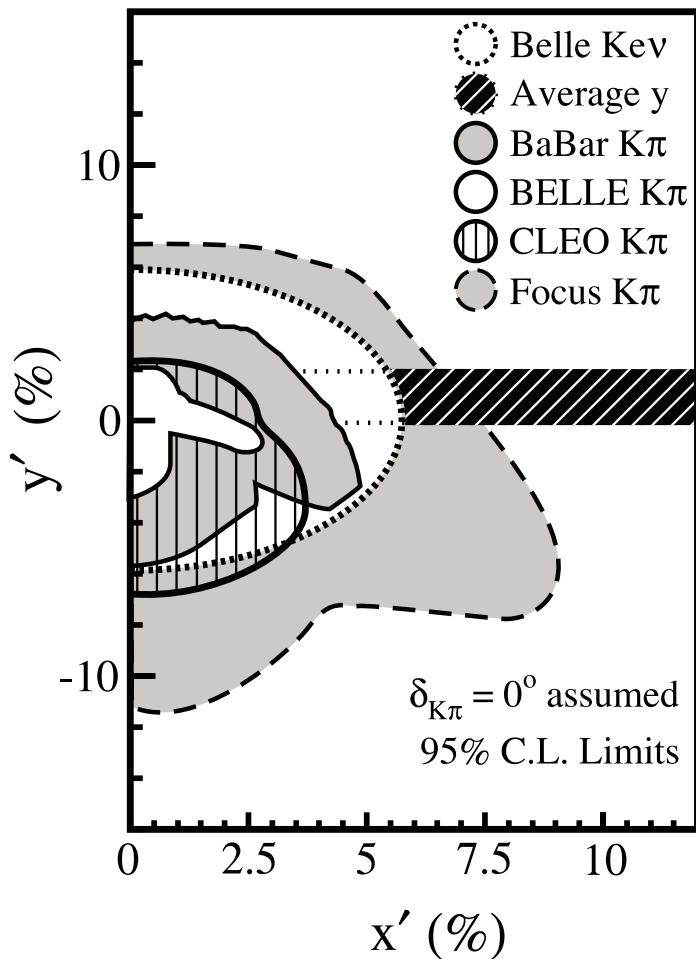
where  $D_\pm$  denotes a  $CP$  eigenstate, implies that

$$\cos \delta = \frac{B(D_+ \rightarrow K^- \pi^+) - B(D_- \rightarrow K^- \pi^+)}{2\sqrt{R_D} B(D^0 \rightarrow K^- \pi^+)}, \quad (24)$$

neglecting  $CP$  violation and exploiting  $R_D \ll \sqrt{R_D}$ .

The strong phase  $\delta$  might also be determined by constructing amplitude quadrangles from a complete set of branching fraction measurements of the other DCS  $D$  decays to two pseudoscalars [12]. This analysis would have to assume that the amplitudes from both  $\Delta I = 1$  and  $\Delta I = 0$  that populate the total  $I = 1/2$   $K\pi$  state have the same strong phase relative to the amplitude that populates the total  $I = 3/2$   $K\pi$  state.

The Dalitz-plot analyses of DCS  $D$  decays to a pseudoscalar and a vector allow the measurement of the relative strong phase between some amplitudes, providing additional constraints to the amplitude quadrangle [13] and thus the determination of the strong phase difference between the relevant DCS and CF amplitudes. In  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ , the DCS and CF decay amplitudes populate the same Dalitz plot, which allows direct



**Figure 1:** Allowed regions in the  $x'y'$  plane. The allowed region for  $y$  is the average of the results from E791<sup>a</sup>, FOCUS<sup>b</sup>, CLEO<sup>c</sup>, BABAR<sup>d</sup>, and Belle<sup>e</sup>. Also shown is the limit from  $D^0 \rightarrow K^{(*)}\ell\nu$  from Belle<sup>f</sup> and limits from  $D \rightarrow K\pi$  from CLEO<sup>g</sup>, BABAR<sup>h</sup>, Belle<sup>i</sup> and FOCUS<sup>j</sup>. The CLEO, BABAR and Belle results allow  $CP$  violation in the decay and mixing amplitudes, and in the interference between these two processes. The FOCUS result does not allow  $CP$  violation. We assume  $\delta = 0$  to place the  $y$  results. A non-zero  $\delta$  would rotate the  $D^0 \rightarrow CP$  eigenstates confidence region clockwise about the origin by  $\delta$ . All results are consistent with the absence of mixing. See the end of the  $D^0$  listings for these references: <sup>a</sup>AITALA 99E, <sup>b</sup>LINK 00, <sup>c</sup>CSORNA 02, <sup>d</sup>AUBERT 03P, <sup>e</sup>ABE 02I, <sup>f</sup>BITENC 05, <sup>g</sup>GODANG 00, <sup>h</sup>AUBERT 03Z, <sup>i</sup>ZHANG 06, <sup>j</sup>LINK 05. See full-color version on color pages at end of book.

measurement of the relative strong phase. CLEO has measured the relative phase between  $D^0 \rightarrow K^*(892)^+\pi^-$  and  $D^0 \rightarrow K^*(892)^-\pi^+$  to be  $(189 \pm 10 \pm 3^{+15}_{-5})^\circ$  [MURAMATSU 02], consistent with the  $180^\circ$  expected from Cabibbo factors and a small strong phase.

There are several results for  $R$  measured in multibody final states with nonzero strangeness. Here  $R$ , defined in Eq. (22), becomes an average over the Dalitz space, weighted by experimental efficiencies and acceptance. Table 4 summarizes the results.

**Table 4:** Results for  $R$  in  $D^0 \rightarrow K^{(*)+}\pi^-(n\pi)$ .

Year	Exper.	$D^0$ final state	$R(\%)$
2005	Belle <sup>a</sup>	$K^+\pi^-\pi^+\pi^-$	$0.320 \pm 0.019^{+0.018}_{-0.013}$
2005	Belle <sup>a</sup>	$K^+\pi^-\pi^0$	$0.229 \pm 0.017^{+0.013}_{-0.009}$
2002	CLEO <sup>b</sup>	$K^{*+}\pi^-$	$0.5 \pm 0.2^{+0.6}_{-0.1}$
2001	CLEO <sup>c</sup>	$K^+\pi^-\pi^+\pi^-$	$0.41^{+0.12}_{-0.11} \pm 0.04$
2001	CLEO <sup>d</sup>	$K^+\pi^-\pi^0$	$0.43^{+0.11}_{-0.10} \pm 0.07$
1998	E791 <sup>e</sup>	$K^+\pi^-\pi^+\pi^-$	$0.68^{+0.34}_{-0.33} \pm 0.07$

See the end of the  $D^0$  listings for these references: <sup>a</sup>TIAN 05, <sup>b</sup>MURAMATSU 02, <sup>c</sup>DYTMAN 01, <sup>d</sup>BRANDENBURG 01, <sup>e</sup>AITALA 98.

For multibody final states, Eqs. (13)–(22) apply to one point in the Dalitz space. Although  $x$  and  $y$  do not vary across the space, knowledge of the resonant substructure is needed to extrapolate the strong phase difference  $\delta$  from point to point. Both the sign and magnitude of  $x$  and  $y$  may be measured using the time-dependent resonant substructure of multibody  $D^0$  decays. CLEO has performed a time-dependent Dalitz-plot

analysis of  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ , and reports  $(-4.5 < x < 9.3)\%$  and  $(-6.4 < y < 3.6)\%$  at the 95% confidence level, without phase or sign ambiguity [ASNER 05], as shown in Figure 2.

**Decays to  $CP$  Eigenstates:** When the final state  $f$  is a  $CP$  eigenstate, there is no distinction between  $f$  and  $\bar{f}$ , and then  $A_f = A_{\bar{f}}$  and  $\bar{A}_{\bar{f}} = \bar{A}_f$ . We denote final states with  $CP$  eigenvalues  $\pm 1$  by  $f_{\pm}$ . In analogy with Eqs. (7)–(8), the decay rates to  $CP$  eigenstates are then

$$\begin{aligned} r_{\pm}(t) &= \frac{|\langle f_{\pm} | H | D^0(t) \rangle|^2}{|\bar{A}_{\pm}|^2} \\ &= \frac{1}{4} \left| h_{\pm}(t) \left( \frac{A_{\pm}}{\bar{A}_{\pm}} \pm \frac{q}{p} \right) + h_{\mp}(t) \left( \frac{A_{\pm}}{\bar{A}_{\pm}} \mp \frac{q}{p} \right) \right|^2, \\ &\propto \frac{1}{|p|^2} \left| h_{\pm}(t) + \eta_{\pm} h_{\mp}(t) \right|^2, \end{aligned} \quad (25)$$

and

$$\bar{r}_{\pm}(t) = \frac{|\langle f_{\pm} | H | \bar{D}^0(t) \rangle|^2}{|A_{\pm}|^2} \propto \frac{1}{|q|^2} \left| h_{\pm}(t) - \eta_{\pm} h_{\mp}(t) \right|^2, \quad (26)$$

where

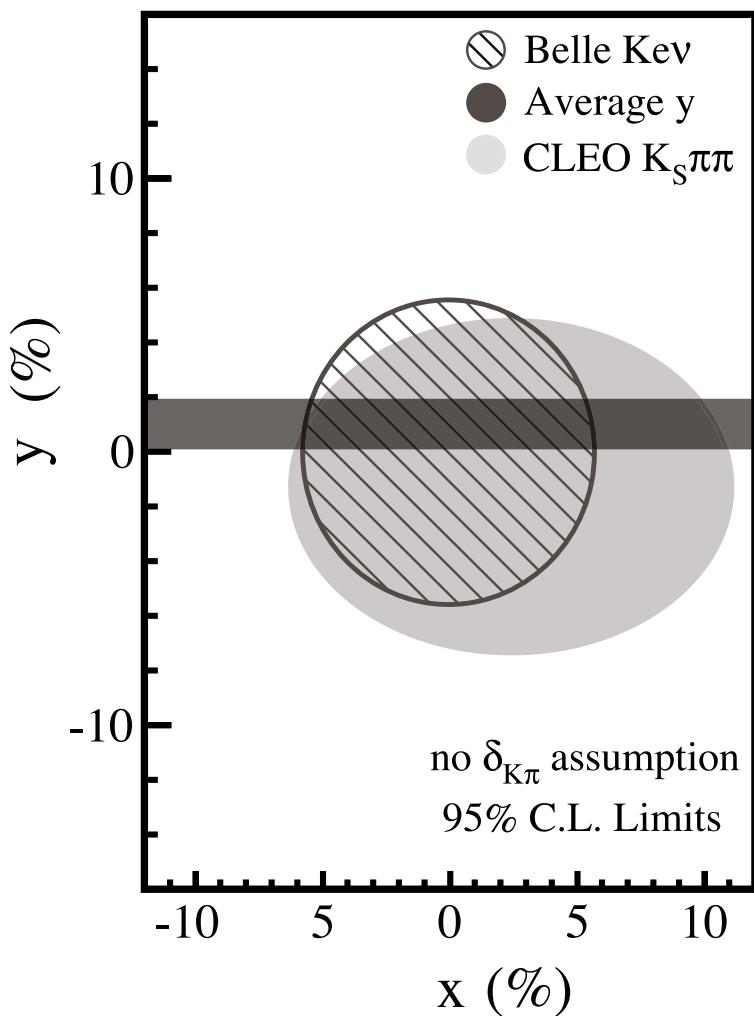
$$h_{\pm}(t) = g_+(t) \pm g_-(t) = e^{-iz_{\pm}t}, \quad (27)$$

and

$$\eta_{\pm} \equiv \frac{pA_{\pm} \mp q\bar{A}_{\pm}}{pA_{\pm} \pm q\bar{A}_{\pm}} = \frac{1 \mp \chi_{\pm}}{1 \pm \chi_{\pm}}. \quad (28)$$

The variable  $\eta_{\pm}$  describes  $CP$  violation; it can receive contributions from each of the three fundamental types of  $CP$  violation.

The quantity  $y$  may be measured by comparing the rate for decays to non- $CP$  eigenstates such as  $D^0 \rightarrow K^- \pi^+$  with decays to  $CP$  eigenstates such as  $D^0 \rightarrow K^+ K^-$  [11]. A positive



**Figure 2:** Allowed regions in the  $xy$  plane. No assumption is made regarding  $\delta$ . The allowed region for  $y$  is the average of the results from E791<sup>a</sup>, FOCUS<sup>b</sup>, CLEO<sup>c</sup>, BABAR<sup>d</sup>, and Belle<sup>e</sup>. Also shown is the limit from  $D^0 \rightarrow K^{(*)} \ell \nu$  from Belle<sup>f</sup>. The CLEO experiment has constrained  $x$  and  $y$  with the time-dependent Dalitz-plot analysis of  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ <sup>g</sup>. All results are consistent with the absence of mixing. See the end of the  $D^0$  listings for these references: <sup>a</sup>AITALA 99E, <sup>b</sup>LINK 00, <sup>c</sup>CSORNA 02, <sup>d</sup>AUBERT 03P, <sup>e</sup>ABE 02I, <sup>f</sup>BITENC 05, <sup>g</sup>ASNER 05.

$y$  would make  $K^+K^-$  decays appear to have a shorter lifetime than  $K^-\pi^+$  decays. The decay rate for a  $D^0$  into a  $CP$  eigenstate is not described by a single exponential in the presence of  $CP$  violation.

In the limit of weak mixing, where  $|ix + y| \ll 1$ , and small  $CP$  violation, where  $|A_M|$ ,  $|A_D|$ , and  $|\sin \phi| \ll 1$ , the time dependence of decays to  $CP$  eigenstates is proportional to a single exponential:

$$r_{\pm}(t) \propto \exp\left(-[1 \pm \left|\frac{p}{q}\right|(y \cos \phi - x \sin \phi)]t\right), \quad (29)$$

$$\bar{r}_{\pm}(t) \propto \exp\left(-[1 \pm \left|\frac{q}{p}\right|(y \cos \phi + x \sin \phi)]t\right), \quad (30)$$

$$r_{\pm}(t) + \bar{r}_{\pm}(t) \propto e^{-(1 \pm y_{CP})t}. \quad (31)$$

Here

$$\begin{aligned} y_{CP} = & y \cos \phi \left[ \frac{1}{2} \left( \left| \frac{p}{q} \right| + \left| \frac{q}{p} \right| \right) + \frac{A_{\text{prod}}}{2} \left( \left| \frac{p}{q} \right| - \left| \frac{q}{p} \right| \right) \right] \\ & - x \sin \phi \left[ \frac{1}{2} \left( \left| \frac{p}{q} \right| - \left| \frac{q}{p} \right| \right) + \frac{A_{\text{prod}}}{2} \left( \left| \frac{p}{q} \right| + \left| \frac{q}{p} \right| \right) \right], \end{aligned} \quad (32)$$

and

$$A_{\text{prod}} \equiv \frac{N(D^0) - N(\overline{D}^0)}{N(D^0) + N(\overline{D}^0)} \quad (33)$$

is defined as the production asymmetry of the  $D^0$  and  $\overline{D}^0$ .

The possibility of  $CP$  violation has been considered in the limit of weak mixing and small  $CP$  violation. In this limit there is no sensitivity to  $CP$  violation in direct decay. Belle [14] and BABAR [AUBERT 03P] have measure  $A_{\Gamma}$ , where

$$A_{\Gamma} \equiv \frac{r_{\pm}(t) - \bar{r}_{\pm}(t)}{r_{\pm}(t) + \bar{r}_{\pm}(t)} \approx A_M y \cos \phi - x \sin \phi,$$

allowing  $CP$  violation in interference and mixing.

In the limit of  $CP$  conservation,  $A_{\pm} = \pm \bar{A}_{\pm}$ ,  $\eta_{\pm} = 0$ ,  $y = y_{CP}$ , and

$$r_{\pm}(t) |\bar{A}_{\pm}|^2 = \bar{r}_{\pm}(t) |A_{\pm}|^2 \propto e^{-(1 \pm y_{CP})t}. \quad (34)$$

All measurements of  $y$  and  $A_{\Gamma}$  are relative to the  $D^0 \rightarrow K^-\pi^+$  decay rate. Table 5 summarizes the current status of measurements. The average of the six  $y_{CP}$  measurements is  $0.90 \pm 0.42\%$ .

**Table 5:** Results for  $y$  from  $D^0 \rightarrow K^+K^-$  and  $\pi^+\pi^-$ .

Year	Exper.	$D^0$ final state(s)	$y_{CP}(\%)$	$A_{\Gamma}(\times 10^{-3})$
2003	Belle [14]	$K^+K^-$	$1.15 \pm 0.69 \pm 0.38$	$-2.0 \pm 6.3 \pm 3.0$
2003	BABAR <sup>a</sup>	$K^+K^-, \pi^+\pi^-$	$0.8 \pm 0.4^{+0.5}_{-0.4}$	$-8 \pm 6 \pm 2$
2001	CLEO <sup>b</sup>	$K^+K^-, \pi^+\pi^-$	$-1.1 \pm 2.5 \pm 1.4$	—
2001	Belle <sup>c</sup>	$K^+K^-$	$-0.5 \pm 1.0^{+0.7}_{-0.8}$	—
2000	FOCUS <sup>d</sup>	$K^+K^-$	$3.4 \pm 1.4 \pm 0.7$	—
1999	E791 <sup>e</sup>	$K^+K^-$	$0.8 \pm 2.9 \pm 1.0$	—

See the end of the  $D^0$  listings for these references: <sup>a</sup>AUBERT 03P, <sup>b</sup>CSORNA 02, <sup>c</sup>ABE 02I, <sup>d</sup>LINK 00, <sup>e</sup>AITALA 99E.

Substantial work on the integrated  $CP$  asymmetries in decays to  $CP$  eigenstates indicates that  $A_{CP}$  is consistent with zero at the few percent level [15]. The expression for the integrated  $CP$  asymmetry that includes the possibility of  $CP$  violation in mixing is

$$A_{CP} = \frac{\Gamma(D^0 \rightarrow f_{\pm}) - \Gamma(\bar{D}^0 \rightarrow f_{\pm})}{\Gamma(D^0 \rightarrow f_{\pm}) + \Gamma(\bar{D}^0 \rightarrow f_{\pm})} \quad (35)$$

$$= \frac{|q|^2 - |p|^2}{|q|^2 + |p|^2} + 2\text{Re}(\eta_{\pm}). \quad (36)$$

**Coherent  $D^0\bar{D}^0$  Analyses:** Measurements of  $R_D$ ,  $\cos\delta$ ,  $x$ , and  $y$  can be made simultaneously in a combined fit to the single-tag (ST) and double-tag (DT) yields or individually by a series of “targeted” analyses [16, 17].

The “comprehensive” analysis simultaneously measures mixing and DCS parameters by examining various ST and DT rates. Due to quantum correlations in the  $C = -1$  and  $C = +1$   $D^0\bar{D}^0$  pairs produced in the reactions  $e^+e^- \rightarrow D^0\bar{D}^0(\pi^0)$  and  $e^+e^- \rightarrow D^0\bar{D}^0\gamma(\pi^0)$ , respectively, the time-integrated  $D^0\bar{D}^0$  decay rates are sensitive to interference between amplitudes for indistinguishable final states. The size of this interference is governed by the relevant amplitude ratios and can include contributions from  $D^0$ - $\bar{D}^0$  mixing.

**Table 6:** CLEO-c results from time-integrated yields at  $\psi(3770) \rightarrow D\bar{D}$ .

Parameter	CLEO-c fitted value	Other results
$y$ (Table 5)	$-0.058 \pm 0.066$	$(0.90 \pm 0.42)\%$
$\cos\delta_{K\pi}$	$1.09 \pm 0.66$	—
$R_M$ (Table 1)	$(1.7 \pm 1.5) \times 10^{-3}$	$< 0.1\%$ (95% C.L.)
$x^2/2$ (Table 3)	$< 0.44\%$ @ (95% C.L.)	$< 0.036\%$ (95% C.L.)

The following categories of final states are considered:

**$f$  or  $\bar{f}$ :** Hadronic states accessed from either  $D^0$  or  $\bar{D}^0$  decay but that are not  $CP$  eigenstates. An example is  $K^-\pi^+$ , which results from Cabibbo-favored  $D^0$  transitions or DCS  $\bar{D}^0$  transitions.

**$\ell^+$  or  $\ell^-$ :** Semileptonic or purely leptonic final states, which, in the absence of mixing, tag unambiguously the flavor of the parent  $D$ .

**$S_+$  or  $S_-$ :**  $CP$ -even and  $CP$ -odd eigenstates, respectively.

The decay rates for  $D^0\bar{D}^0$  pairs to all possible combinations of the above categories of final states are calculated in Ref. [4], for both  $C = -1$  and  $C = +1$ , reproducing the work of Refs. [5, 10]. Such  $D^0\bar{D}^0$  combinations, where both  $D$  final states are specified, are double tags. In addition, the rates for single tags, where either the  $D^0$  or  $\bar{D}^0$  is identified and the other neutral  $D$  decays generically are given in Ref. [4].

CLEO-c has reported results using 281 pb<sup>-1</sup> of  $e^+e^- \rightarrow \psi(3770)$  data [18], where the quantum coherent  $D^0\bar{D}^0$  pairs are in the  $C = -1$  state. The values of  $y$ ,  $R_M$ , and  $\cos\delta$  are determined from a combined fit to the ST (hadronic only) and DT yields. The hadronic final states included in the analysis are  $K^-\pi^+$  ( $f$ ),  $K^+\pi^-$  ( $\bar{f}$ ),  $K^-K^+$  ( $S_+$ ),  $\pi^+\pi^-$  ( $S_+$ ),  $K_S^0\pi^0\pi^0$  ( $S_+$ ), and  $K_S^0\pi^0$  ( $S_-$ ). Both of the two flavored final states,  $K^-\pi^+$  and  $K^+\pi^-$ , can be reached via CF or DCS transitions.

Semileptonic DT yields are also included, where one  $D$  is fully reconstructed in one of the hadronic modes listed above, and the other  $D$  is partially reconstructed, requiring that only the electron be found. When the electron is accompanied by a flavor tag ( $D \rightarrow K^-\pi^+$  or  $K^+\pi^-$ ), only the “right-sign” DT sample, where the electron and kaon charges are the same, is used. Extraction of the DCS “wrong-sign” semileptonic yield is not feasible with the current CLEO-c data sample, and the parameter  $r_{K\pi}$  is constrained to the world average. Table 6 shows the results of the fit to the CLEO-c data.

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$$|m_{D_1^0} - m_{D_2^0}|$$

The  $D_1^0$  and  $D_2^0$  are the mass eigenstates of the  $D^0$  meson, as described in the note on “ $D^0$ - $\bar{D}^0$  Mixing,” above.

<i>VALUE</i> ( $10^{10} \text{ } \text{\AA} \text{ s}^{-1}$ )	<i>CL%</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
< 7	95	5 ZHANG	06 BELL	$e^+ e^-$
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
–11 to +22		6 ASNER	05 CLEO	$e^+ e^- \approx 10 \text{ GeV}$
< 11	90	BITENC	05 BELL	
< 30	90	CAWLFIELD	05 CLEO	
< 7	95	5 LI	05A BELL	See ZHANG 06
< 22	95	7 LINK	05H FOCS	$\gamma$ nucleus
< 23	95	AUBERT	04Q BABR	
< 11	95	5 AUBERT	03Z BABR	$e^+ e^-, 10.6 \text{ GeV}$
< 7	95	8 GODANG	00 CLE2	$e^+ e^-$
< 32	90	9,10 AITALA	98 E791	$\pi^-$ nucleus, 500 GeV
< 24	90	11 AITALA	96C E791	$\pi^-$ nucleus, 500 GeV
< 21	90	10,12 ANJOS	88C E691	Photoproduction

<sup>5</sup> The AUBERT 03Z, LI 05A, and ZHANG 06 limits are inferred from the  $D^0$ - $\bar{D}^0$  mixing ratio  $\Gamma(K^+\pi^- \text{ (via } \bar{D}^0))/\Gamma(K^-\pi^+)$  given near the end of this  $D^0$  Listings. Decay-time information is used to distinguish DCS decays from  $D^0$ - $\bar{D}^0$  mixing. The limit allows interference between the DCS and mixing ratios, and also allows  $CP$  violation. AUBERT 03Z assumes the strong phase between  $D^0 \rightarrow K^+\pi^-$  and  $\bar{D}^0 \rightarrow K^+\pi^-$  amplitudes is small; if an arbitrary phase is allowed, the limit degrades by 20%. The LI 05A and ZHANG 06 limits are valid for an arbitrary strong phase.

- <sup>6</sup>This ASNER 05 limit is from the time-dependent Dalitz-plot analysis of  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ . Decay-time information and interference on the Dalitz plot are used to distinguish doubly Cabibbo-suppressed decays from mixing and to measure the relative phase between  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^+ \pi^-$ . This limit allows  $CP$  violation and is sensitive to the sign of  $\Delta m$ .
- <sup>7</sup>This LINK 05H limit is inferred from the  $D^0$ - $\bar{D}^0$  mixing ratio  $\Gamma(K^+ \pi^-)$  (via  $\bar{D}^0$ )/ $\Gamma(K^- \pi^+)$  given near the end of this  $D^0$  Listings. Decay-time information is used to distinguish DCS decays from  $D^0$ - $\bar{D}^0$  mixing. The limit allows interference between the DCS and mixing ratios, and also allows  $CP$  violation. The strong phase between  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^+ \pi^-$  is assumed to be small. If an arbitrary relative strong phase is allowed, the limit degrades by 25%.
- <sup>8</sup>This GODANG 00 limit is inferred from the  $D^0$ - $\bar{D}^0$  mixing ratio  $\Gamma(K^+ \pi^-)$  (via  $\bar{D}^0$ )/ $\Gamma(K^- \pi^+)$  given near the end of this  $D^0$  Listings. Decay-time information is used to distinguish DCS decays from  $D^0$ - $\bar{D}^0$  mixing. The limit allows interference between the DCS and mixing ratios, and also allows  $CP$  violation. The strong phase between  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^+ \pi^-$  is assumed to be small. If an arbitrary relative strong phase is allowed, the limit degrades by a factor of two.
- <sup>9</sup>AITALA 98 allows interference between the doubly Cabibbo-suppressed and mixing amplitudes, and also allows  $CP$  violation in this term, but assumes that  $A_D = A_R = 0$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above.
- <sup>10</sup>This limit is inferred from  $R_M$  for  $f = K^+ \pi^-$  and  $f = K^+ \pi^- \pi^+ \pi^-$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above. Decay-time information is used to distinguish doubly Cabibbo-suppressed decays from  $D^0$ - $\bar{D}^0$  mixing.
- <sup>11</sup>This limit is inferred from  $R_M$  for  $f = K^+ \ell^- \bar{\nu}_\ell$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above.
- <sup>12</sup>ANJOS 88C assumes that  $y = 0$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above. Without this assumption, the limit degrades by about a factor of two.

$$(\Gamma_{D_1^0} - \Gamma_{D_2^0})/\Gamma = 2y$$

The  $D_1^0$  and  $D_2^0$  are the mass eigenstates of the  $D^0$  meson, as described in the note on " $D^0$ - $\bar{D}^0$  Mixing," above.

VALUE (units $10^{-2}$ )	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>1.4 ± 1.0 OUR AVERAGE</b>					
-3.0 + 5.0 - 4.8 - 0.8	+1.6		13 ASNER	05 CLEO	$e^+ e^- \approx 10$ GeV
1.6 ± 0.8 - 0.8	+1.0	450k	14 AUBERT	03P BABR	$e^+ e^- \approx \gamma(4S)$
-1.0 ± 2.0 - 1.6	+1.4	18k	15 ABE	02I BELL	$e^+ e^- \approx \gamma(4S)$
-2.4 ± 5.0 6.84 ± 2.78 ± 1.48	±2.8	3393 10k	16 CSORNA 15 LINK	02 CLE2 00 FOCS	$e^+ e^- \approx \gamma(4S)$ $\gamma$ nucleus
+1.6 ± 5.8 - 2.1	±2.1		15 AITALA	99E E791	$K^- \pi^+, K^+ K^-$
• • • We do not use the following data for averages, fits, limits, etc. • • •					
-0.7 ± 4.9		4k ± 88	17,18 ZHANG	06 BELL	$e^+ e^-$
-0.3 ± 5.7			17,18 LI	05A BELL	See ZHANG 06
-5.2 +18.4 - 16.8			17,18 LINK	05H FOCS	$\gamma$ nucleus
1.6 + 6.2 - 12.8			17,18 AUBERT	03Z BABR	$e^+ e^-, 10.6$ GeV

$-5.0 \pm 2.8 \pm 0.6$	<sup>18</sup> GODANG	00	CLE2	$e^+ e^-$
$ \Delta\Gamma /\Gamma < 26$	90	<sup>19,20</sup> AITALA	98	E791 $\pi^-$ nucleus, 500 GeV
$ \Delta\Gamma /\Gamma < 20$	90	<sup>21</sup> AITALA	96C	E791 $\pi^-$ nucleus, 500 GeV
$ \Delta\Gamma /\Gamma < 17$	90	<sup>20,22</sup> ANJOS	88C	E691 Photoproduction

- <sup>13</sup>This ASNER 05 limit is from the time-dependent Dalitz-plot analysis of  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ . Decay-time information and interference on the Dalitz plot are used to distinguish doubly Cabibbo-suppressed decays from mixing and to measure the relative phase between  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^+ \pi^-$ . This limit allows  $CP$  violation.
- <sup>14</sup>AUBERT 03P measures  $Y \equiv 2 \tau^0 / (\tau^+ + \tau^-) - 1$ , where  $\tau^0$  is the  $D^0 \rightarrow K^- \pi^+$  (and  $\bar{D}^0 \rightarrow K^+ \pi^-$ ) lifetime, and  $\tau^+$  and  $\tau^-$  are the  $D^0$  and  $\bar{D}^0$  lifetimes to  $CP$ -even states (here  $K^- K^+$  and  $\pi^- \pi^+$ ). In the limit of  $CP$  conservation,  $Y = y \equiv \Delta\Gamma / 2\Gamma$  (we list  $2y = \Delta\Gamma/\Gamma$ ). AUBERT 03P also uses  $\tau^+ - \tau^-$  to get  $\Delta Y = -0.008 \pm 0.006 \pm 0.002$ .
- <sup>15</sup>LINK 00, AITALA 99E, and ABE 02I measure the lifetime difference between  $D^0 \rightarrow K^- K^+$  ( $CP$  even) decays and  $D^0 \rightarrow K^- \pi^+$  ( $CP$  mixed) decays, or  $y_{CP} = [\Gamma(CP+) - \Gamma(CP-)] / [\Gamma(CP+) + \Gamma(CP-)]$ . We list  $2y_{CP} = \Delta\Gamma/\Gamma$ .
- <sup>16</sup>CSORNA 02 measures the lifetime difference between  $D^0 \rightarrow K^- K^+$  and  $\pi^- \pi^+$  ( $CP$  even) decays and  $D^0 \rightarrow K^- \pi^+$  ( $CP$  mixed) decays, or  $y_{CP} = [\Gamma(CP+) - \Gamma(CP-)] / [\Gamma(CP+) + \Gamma(CP-)]$ . We list  $2y_{CP} = \Delta\Gamma/\Gamma$ .
- <sup>17</sup>The ranges of AUBERT 03Z, LINK 05H, LI 05A, and ZHANG 06 measurements are for 95% confidence level.
- <sup>18</sup>The GODANG 00, AUBERT 03Z, LINK 05H, LI 05A, and ZHANG 06 limits are inferred from the  $D^0$ - $\bar{D}^0$  mixing ratio  $\Gamma(K^+ \pi^-)$  (via  $\bar{D}^0$ )/ $\Gamma(K^- \pi^+)$  given near the end of this  $D^0$  Listings. Decay-time information is used to distinguish DCS decays from  $D^0$ - $\bar{D}^0$  mixing. The limit allows interference between the DCS and mixing ratios, and also allows  $CP$  violation. The phase between  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^+ \pi^-$  is assumed to be small. This is a measurement of  $y'$  and is not the same as the  $y_{CP}$  of our note above on " $D^0$ - $\bar{D}^0$  Mixing."
- <sup>19</sup>AITALA 98 allows interference between the doubly Cabibbo-suppressed and mixing amplitudes, and also allows  $CP$  violation in this term, but assumes that  $A_D = A_R = 0$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above.
- <sup>20</sup>This limit is inferred from  $R_M$  for  $f = K^+ \pi^-$  and  $f = K^+ \pi^- \pi^+ \pi^-$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above. Decay-time information is used to distinguish doubly Cabibbo-suppressed decays from  $D^0$ - $\bar{D}^0$  mixing.
- <sup>21</sup>This limit is inferred from  $R_M$  for  $f = K^+ \ell^- \bar{\nu}_\ell$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above.
- <sup>22</sup>ANJOS 88C assumes that  $y = 0$ . See the note on " $D^0$ - $\bar{D}^0$  Mixing," above. Without this assumption, the limit degrades by about a factor of two.

## $D^0$ DECAY MODES

Most decay modes (other than the semileptonic modes) that involve a neutral  $K$  meson are now given as  $K_S^0$  modes, not as  $\bar{K}^0$  modes. Nearly always it is a  $K_S^0$  that is measured, and interference between Cabibbo-allowed and doubly Cabibbo-suppressed modes can invalidate the assumption that  $2\Gamma(K_S^0) = \Gamma(\bar{K}^0)$ .

Mode	Fraction ( $\Gamma_i/\Gamma$ )	Scale factor/ Confidence level
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### Topological modes

$\Gamma_1$	0-prongs	[a]	(19 $\pm$ 6) %
$\Gamma_2$	2-prongs		(67 $\pm$ 6) %
$\Gamma_3$	4-prongs	[b]	(13.8 $\pm$ 0.5) %
$\Gamma_4$	6-prongs		( 1.2 $\pm$ 1.3 ) $\times 10^{-3}$

### Inclusive modes

$\Gamma_5$	$e^+$ anything	[c]	( 6.71 $\pm$ 0.29 ) %
$\Gamma_6$	$\mu^+$ anything		( 6.5 $\pm$ 0.7 ) %
$\Gamma_7$	$K^-$ anything		(53 $\pm$ 4) %
$\Gamma_8$	$\bar{K}^0$ anything + $K^0$ anything		(42 $\pm$ 5) %
$\Gamma_9$	$K^+$ anything		( 3.4 $\pm$ 0.6 ) %
$\Gamma_{10}$	$\bar{K}^*(892)^0$ anything		( 9 $\pm$ 4 ) %
$\Gamma_{11}$	$K^*(892)^0$ anything		( 2.8 $\pm$ 1.3 ) %
$\Gamma_{12}$	$\eta$ anything	[d]	< 13 %
$\Gamma_{13}$	$\phi$ anything		( 1.7 $\pm$ 0.8 ) %

### Semileptonic modes

$\Gamma_{14}$	$K^- \ell^+ \nu_\ell$		
$\Gamma_{15}$	$K^- e^+ \nu_e$		( 3.51 $\pm$ 0.11 ) %
$\Gamma_{16}$	$K^- \mu^+ \nu_\mu$		( 3.19 $\pm$ 0.16 ) %
$\Gamma_{17}$	$K^*(892)^- e^+ \nu_e$		( 2.17 $\pm$ 0.16 ) %
$\Gamma_{18}$	$K^*(892)^- \mu^+ \nu_\mu$		( 1.95 $\pm$ 0.25 ) %
$\Gamma_{19}$	$K^- \pi^0 e^+ \nu_e$		
$\Gamma_{20}$	$\bar{K}^0 \pi^- e^+ \nu_e$		
$\Gamma_{21}$	$\bar{K}^0 \pi^- \mu^+ \nu_\mu$		
$\Gamma_{22}$	$K^*(892)^- \ell^+ \nu_\ell$		
$\Gamma_{23}$	$\bar{K}^*(892)^0 \pi^- e^+ \nu_e$		
$\Gamma_{24}$	$K^- \pi^+ \pi^- \mu^+ \nu_\mu$	< 1.2	$\times 10^{-3}$ CL=90%
$\Gamma_{25}$	$(\bar{K}^*(892)\pi)^- \mu^+ \nu_\mu$	< 1.4	$\times 10^{-3}$ CL=90%
$\Gamma_{26}$	$\pi^- e^+ \nu_e$		( 2.81 $\pm$ 0.19 ) $\times 10^{-3}$
$\Gamma_{27}$	$\pi^- \mu^+ \nu_\mu$		( 2.4 $\pm$ 0.4 ) $\times 10^{-3}$
$\Gamma_{28}$	$\rho^- e^+ \nu_e$		( 1.9 $\pm$ 0.4 ) $\times 10^{-3}$

### Hadronic modes with one $\bar{K}$

$\Gamma_{29}$	$K^- \pi^+$		( 3.80 $\pm$ 0.07 ) %
$\Gamma_{30}$	$K_S^0 \pi^0$		( 1.14 $\pm$ 0.12 ) %
$\Gamma_{31}$	$K_S^0 \pi^+ \pi^-$	[e]	( 2.90 $\pm$ 0.19 ) %
$\Gamma_{32}$	$K_S^0 \rho^0$		( 7.5 $\pm$ 0.6 ) $\times 10^{-3}$
$\Gamma_{33}$	$K_S^0 \omega, \omega \rightarrow \pi^+ \pi^-$		( 2.1 $\pm$ 0.6 ) $\times 10^{-4}$
$\Gamma_{34}$	$K_S^0 f_0(980),$ $f_0(980) \rightarrow \pi^+ \pi^-$		( 1.36 $\pm$ 0.30 ) $\times 10^{-3}$

$\Gamma_{35}$	$K_S^0 f_2(1270),$ $f_2(1270) \rightarrow \pi^+ \pi^-$	$(1.3 \pm 0.7) \times 10^{-4}$
$\Gamma_{36}$	$K_S^0 f_0(1370),$ $f_0(1370) \rightarrow \pi^+ \pi^-$	$(2.5 \pm 0.6) \times 10^{-3}$
$\Gamma_{37}$	$K^*(892)^- \pi^+,$ $K^*(892)^- \rightarrow K_S^0 \pi^-$	$(1.91 \pm 0.14)\%$
$\Gamma_{38}$	$K^*(892)^+ \pi^-,$ $K^*(892)^+ \rightarrow K_S^0 \pi^+$	[f] $(10 \pm 12/4) \times 10^{-5}$
$\Gamma_{39}$	$K_0^*(1430)^- \pi^+,$ $K_0^*(1430)^- \rightarrow K_S^0 \pi^-$	$(2.8 \pm 0.6) \times 10^{-3}$
$\Gamma_{40}$	$K_2^*(1430)^- \pi^+,$ $K_2^*(1430)^- \rightarrow K_S^0 \pi^-$	$(3.2 \pm 2.1) \times 10^{-4}$
$\Gamma_{41}$	$K^*(1680)^- \pi^+,$ $K^*(1680)^- \rightarrow K_S^0 \pi^-$	$(6 \pm 5) \times 10^{-4}$
$\Gamma_{42}$	$K_S^0 \pi^+ \pi^-$ nonresonant	$(2.6 \pm 5.9) \times 10^{-4}$
$\Gamma_{43}$	$K^- \pi^+ \pi^0$	[e] $(14.1 \pm 0.5)\%$
$\Gamma_{44}$	$K^- \rho^+$	$(11.0 \pm 0.7)\%$
$\Gamma_{45}$	$K^- \rho(1700)^+,$ $\rho(1700)^+ \rightarrow \pi^+ \pi^0$	$(8.0 \pm 1.7) \times 10^{-3}$
$\Gamma_{46}$	$K^*(892)^- \pi^+,$ $K^*(892)^- \rightarrow K^- \pi^0$	$(2.25 \pm 0.36)\%$
$\Gamma_{47}$	$\bar{K}^*(892)^0 \pi^0,$ $\bar{K}^*(892)^0 \rightarrow K^- \pi^+$	$(1.91 \pm 0.24)\%$
$\Gamma_{48}$	$K_0^*(1430)^- \pi^+,$ $K_0^*(1430)^- \rightarrow K^- \pi^0$	$(4.6 \pm 2.2) \times 10^{-3}$
$\Gamma_{49}$	$\bar{K}_0^*(1430)^0 \pi^0,$ $\bar{K}_0^*(1430)^0 \rightarrow K^- \pi^+$	$(5.8 \pm 4.6) \times 10^{-3}$
$\Gamma_{50}$	$K^*(1680)^- \pi^+,$ $K^*(1680)^- \rightarrow K^- \pi^0$	$(1.8 \pm 0.7) \times 10^{-3}$
$\Gamma_{51}$	$K^- \pi^+ \pi^0$ nonresonant	$(1.13 \pm 0.54)\%$
$\Gamma_{52}$	$K_S^0 \pi^0 \pi^0$	—
$\Gamma_{53}$	$\bar{K}^*(892)^0 \pi^0,$ $\bar{K}^*(892)^0 \rightarrow K_S^0 \pi^0$	$(6.3 \pm 1.8) \times 10^{-3}$
$\Gamma_{54}$	$K_S^0 \pi^0 \pi^0$ nonresonant	$(4.2 \pm 1.1) \times 10^{-3}$
$\Gamma_{55}$	$K^- \pi^+ \pi^+ \pi^-$	[e] $(7.72 \pm 0.28)\%$
$\Gamma_{56}$	$K^- \pi^+ \rho^0$ total	$(6.4 \pm 0.4)\%$
$\Gamma_{57}$	$K^- \pi^+ \rho^0$ 3-body	$(4.9 \pm 2.2) \times 10^{-3}$
$\Gamma_{58}$	$\bar{K}^*(892)^0 \rho^0,$ $\bar{K}^*(892)^0 \rightarrow K^- \pi^+$	$(1.00 \pm 0.22)\%$

$\Gamma_{59}$	$K^- a_1(1260)^+$ , $a_1(1260)^+ \rightarrow \pi^+ \pi^+ \pi^-$	( 3.6 $\pm$ 0.6 ) %
$\Gamma_{60}$	$\bar{K}^*(892)^0 \pi^+ \pi^-$ total, $\bar{K}^*(892)^0 \rightarrow K^- \pi^+$	( 1.5 $\pm$ 0.4 ) %
$\Gamma_{61}$	$\bar{K}^*(892)^0 \pi^+ \pi^-$ 3-body, $\bar{K}^*(892)^0 \rightarrow K^- \pi^+$	( 9.7 $\pm$ 2.1 ) $\times 10^{-3}$
$\Gamma_{62}$	$K_1(1270)^- \pi^+$ , $K_1(1270)^- \rightarrow K^- \pi^+ \pi^-$	[g] ( 2.9 $\pm$ 0.3 ) $\times 10^{-3}$
$\Gamma_{63}$	$K^- \pi^+ \pi^+ \pi^-$ nonresonant	( 1.80 $\pm$ 0.25 ) %
$\Gamma_{64}$	$K_S^0 \pi^+ \pi^- \pi^0$	[e] ( 5.3 $\pm$ 0.6 ) %
$\Gamma_{65}$	$K_S^0 \eta, \eta \rightarrow \pi^+ \pi^- \pi^0$	( 8.6 $\pm$ 1.4 ) $\times 10^{-4}$
$\Gamma_{66}$	$K_S^0 \omega, \omega \rightarrow \pi^+ \pi^- \pi^0$	( 9.8 $\pm$ 1.8 ) $\times 10^{-3}$
$\Gamma_{67}$	$K^*(892)^- \rho^+$ , $K^*(892)^- \rightarrow K_S^0 \pi^-$	( 2.1 $\pm$ 0.8 ) %
$\Gamma_{68}$	$K_1(1270)^- \pi^+$ , $K_1(1270)^- \rightarrow K_S^0 \pi^- \pi^0$	[g] ( 2.2 $\pm$ 0.6 ) $\times 10^{-3}$
$\Gamma_{69}$	$\bar{K}^*(892)^0 \pi^+ \pi^-$ 3-body, $\bar{K}^*(892)^0 \rightarrow K_S^0 \pi^0$	( 2.4 $\pm$ 0.5 ) $\times 10^{-3}$
$\Gamma_{70}$	$K_S^0 \pi^+ \pi^- \pi^0$ nonresonant	( 1.1 $\pm$ 1.1 ) %
$\Gamma_{71}$	$K^- \pi^+ \pi^0 \pi^0$	
$\Gamma_{72}$	$K^- \pi^+ \pi^+ \pi^- \pi^0$	( 4.1 $\pm$ 0.4 ) %
$\Gamma_{73}$	$\bar{K}^*(892)^0 \pi^+ \pi^- \pi^0$ , $\bar{K}^*(892)^0 \rightarrow K^- \pi^+$	( 1.2 $\pm$ 0.6 ) %
$\Gamma_{74}$	$K^- \pi^+ \omega, \omega \rightarrow \pi^+ \pi^- \pi^0$	( 2.7 $\pm$ 0.5 ) %
$\Gamma_{75}$	$\bar{K}^*(892)^0 \omega$ , $\bar{K}^*(892)^0 \rightarrow K^- \pi^+$ , $\omega \rightarrow \pi^+ \pi^- \pi^0$	( 6.5 $\pm$ 2.4 ) $\times 10^{-3}$
$\Gamma_{76}$	$K_S^0 \eta \pi^0$	( 5.2 $\pm$ 1.2 ) $\times 10^{-3}$
$\Gamma_{77}$	$K_S^0 a_0(980)$ , $a_0(980) \rightarrow \eta \pi^0$	( 6.2 $\pm$ 2.0 ) $\times 10^{-3}$
$\Gamma_{78}$	$\bar{K}^*(892)^0 \eta, \bar{K}^*(892)^0 \rightarrow K_S^0 \pi^0$	( 1.5 $\pm$ 0.5 ) $\times 10^{-3}$
$\Gamma_{79}$	$K_S^0 2\pi^+ 2\pi^-$	( 2.75 $\pm$ 0.31 ) $\times 10^{-3}$
$\Gamma_{80}$	$K_S^0 \rho^0 \pi^+ \pi^-$ , no $K^*(892)^-$	( 1.1 $\pm$ 0.7 ) $\times 10^{-3}$
$\Gamma_{81}$	$K^*(892)^- \pi^+ \pi^+ \pi^-$ , $K^*(892)^- \rightarrow K_S^0 \pi^-, \text{no}$ $\rho^0$	( 5 $\pm$ 8 ) $\times 10^{-4}$
$\Gamma_{82}$	$K^*(892)^- \rho^0 \pi^+$ , $K^*(892)^- \rightarrow K_S^0 \pi^-$	( 1.7 $\pm$ 0.7 ) $\times 10^{-3}$
$\Gamma_{83}$	$K_S^0 2\pi^+ 2\pi^-$ nonresonant	< 1.3 $\times 10^{-3}$ CL=90%
$\Gamma_{84}$	$\bar{K}^0 \pi^+ \pi^- \pi^0 \pi^0 (\pi^0)$	
$\Gamma_{85}$	$K^- 3\pi^+ 2\pi^-$	( 2.1 $\pm$ 0.5 ) $\times 10^{-4}$

Fractions of many of the following modes with resonances have already appeared above as submodes of particular charged-particle modes. (Modes for which there are only upper limits and  $\bar{K}^*(892)\rho$  submodes only appear below.)

$\Gamma_{86}$	$K_S^0 \eta$	$(3.8 \pm 0.6) \times 10^{-3}$
$\Gamma_{87}$	$K_S^0 \omega$	$(1.10 \pm 0.20) \%$
$\Gamma_{88}$	$K_S^0 \eta'(958)$	$(9.1 \pm 1.4) \times 10^{-3}$
$\Gamma_{89}$	$K^- a_1(1260)^+$	$(7.5 \pm 1.1) \%$
$\Gamma_{90}$	$\bar{K}^0 a_1(1260)^0$	$< 1.9 \%$ CL=90%
$\Gamma_{91}$	$K^- a_2(1320)^+$	$< 2 \times 10^{-3}$ CL=90%
$\Gamma_{92}$	$\bar{K}^*(892)^0 \pi^+ \pi^-$ total	$(2.3 \pm 0.5) \%$
$\Gamma_{93}$	$\bar{K}^*(892)^0 \pi^+ \pi^-$ 3-body	$(1.46 \pm 0.32) \%$
$\Gamma_{94}$	$\bar{K}^*(892)^0 \rho^0$	$(1.50 \pm 0.33) \%$
$\Gamma_{95}$	$\bar{K}^*(892)^0 \rho^0$ transverse	$(1.6 \pm 0.5) \%$
$\Gamma_{96}$	$\bar{K}^*(892)^0 \rho^0$ S-wave	$(2.9 \pm 0.6) \%$
$\Gamma_{97}$	$\bar{K}^*(892)^0 \rho^0$ S-wave long.	$< 3 \times 10^{-3}$ CL=90%
$\Gamma_{98}$	$\bar{K}^*(892)^0 \rho^0$ P-wave	$< 3 \times 10^{-3}$ CL=90%
$\Gamma_{99}$	$\bar{K}^*(892)^0 \rho^0$ D-wave	$(2.0 \pm 0.6) \%$
$\Gamma_{100}$	$K^*(892)^- \rho^+$	$(6.4 \pm 2.5) \%$
$\Gamma_{101}$	$K^*(892)^- \rho^+$ longitudinal	$(3.1 \pm 1.2) \%$
$\Gamma_{102}$	$K^*(892)^- \rho^+$ transverse	$(3.4 \pm 2.0) \%$
$\Gamma_{103}$	$K^*(892)^- \rho^+$ P-wave	$< 1.5 \%$ CL=90%
$\Gamma_{104}$	$K^- \pi^+ f_0(980)$	
$\Gamma_{105}$	$\bar{K}^*(892)^0 f_0(980)$	
$\Gamma_{106}$	$K_1(1270)^- \pi^+$	[g] $(1.12 \pm 0.31) \%$
$\Gamma_{107}$	$K_1(1400)^- \pi^+$	$< 1.2 \%$ CL=90%
$\Gamma_{108}$	$\bar{K}_1(1400)^0 \pi^0$	$< 3.7 \%$ CL=90%
$\Gamma_{109}$	$K^*(1410)^- \pi^+$	
$\Gamma_{110}$	$\bar{K}^*(892)^0 \pi^+ \pi^- \pi^0$	$(1.8 \pm 0.9) \%$
$\Gamma_{111}$	$\bar{K}^*(892)^0 \eta$	
$\Gamma_{112}$	$K^- \pi^+ \omega$	$(3.0 \pm 0.6) \%$
$\Gamma_{113}$	$\bar{K}^*(892)^0 \omega$	$(1.1 \pm 0.4) \%$
$\Gamma_{114}$	$K^- \pi^+ \eta'(958)$	$(7.2 \pm 1.8) \times 10^{-3}$
$\Gamma_{115}$	$\bar{K}^*(892)^0 \eta'(958)$	$< 1.1 \times 10^{-3}$ CL=90%

### Hadronic modes with three $K$ 's

$\Gamma_{116}$	$K_S^0 K^+ K^-$	$(4.58 \pm 0.34) \times 10^{-3}$
$\Gamma_{117}$	$K_S^0 a_0(980)^0, a_0^0 \rightarrow K^+ K^-$	$(3.0 \pm 0.4) \times 10^{-3}$
$\Gamma_{118}$	$K^- a_0(980)^+, a_0^+ \rightarrow K^+ K_S^0$	$(6.1 \pm 1.8) \times 10^{-4}$
$\Gamma_{119}$	$K^+ a_0(980)^-, a_0^- \rightarrow K^- K_S^0$	$< 1.1 \times 10^{-4}$ CL=95%
$\Gamma_{120}$	$K_S^0 f_0(980), f_0 \rightarrow K^+ K^-$	$< 1.0 \times 10^{-4}$ CL=95%
$\Gamma_{121}$	$K_S^0 \phi, \phi \rightarrow K^+ K^-$	$(2.10 \pm 0.16) \times 10^{-3}$
$\Gamma_{122}$	$K_S^0 f_0(1400), f_0 \rightarrow K^+ K^-$	$(1.7 \pm 1.1) \times 10^{-4}$

$\Gamma_{123}$	$3K_S^0$	$(9.3 \pm 1.3) \times 10^{-4}$
$\Gamma_{124}$	$K^+ K^- K^- \pi^+$	$(2.11 \pm 0.31) \times 10^{-4}$
$\Gamma_{125}$	$K^+ K^- \bar{K}^*(892)^0,$ $\bar{K}^*(892)^0 \rightarrow K^- \pi^+$	$(4.2 \pm 1.7) \times 10^{-5}$
$\Gamma_{126}$	$K^- \pi^+ \phi, \phi \rightarrow K^+ K^-$	$(3.8 \pm 1.6) \times 10^{-5}$
$\Gamma_{127}$	$\phi \bar{K}^*(892)^0,$ $\phi \rightarrow K^+ K^-,$ $\bar{K}^*(892)^0 \rightarrow K^- \pi^+$	$(1.01 \pm 0.20) \times 10^{-4}$
$\Gamma_{128}$	$K^+ K^- K^- \pi^+$ nonresonant	$(3.2 \pm 1.4) \times 10^{-5}$
$\Gamma_{129}$	$K_S^0 K_S^0 K^\pm \pi^\mp$	$(6.1 \pm 1.3) \times 10^{-4}$

**Pionic modes**

$\Gamma_{130}$	$\pi^+ \pi^-$	$(1.364 \pm 0.032) \times 10^{-3}$
$\Gamma_{131}$	$\pi^0 \pi^0$	$(7.9 \pm 0.8) \times 10^{-4}$
$\Gamma_{132}$	$\pi^+ \pi^- \pi^0$	$(1.31 \pm 0.06) \%$
$\Gamma_{133}$	$\rho^+ \pi^-$	$(10.0 \pm 0.6) \times 10^{-3}$
$\Gamma_{134}$	$\rho^0 \pi^0$	$(3.2 \pm 0.4) \times 10^{-3}$
$\Gamma_{135}$	$\rho^- \pi^+$	$(4.5 \pm 0.4) \times 10^{-3}$
$\Gamma_{136}$	$f_0(980)\pi^0, f_0(980) \rightarrow \pi^+ \pi^-$	$< 3.4 \times 10^{-6}$ CL=95%
$\Gamma_{137}$	$f_0(600)\pi^0, f_0(600) \rightarrow \pi^+ \pi^-$	$< 2.7 \times 10^{-5}$ CL=95%
$\Gamma_{138}$	$(\pi^+ \pi^-)_{S\text{-wave}} \pi^0$	$< 2.5 \times 10^{-4}$ CL=95%
$\Gamma_{139}$	$3\pi^0$	$< 3.5 \times 10^{-4}$ CL=90%
$\Gamma_{140}$	$2\pi^+ 2\pi^-$	$(7.31 \pm 0.27) \times 10^{-3}$
$\Gamma_{141}$	$\pi^+ \pi^- 2\pi^0$	$(9.8 \pm 0.9) \times 10^{-3}$
$\Gamma_{142}$	$\eta \pi^0$	$[h] (5.6 \pm 1.4) \times 10^{-4}$
$\Gamma_{143}$	$\omega \pi^0$	$[h] < 2.6 \times 10^{-4}$ CL=90%
$\Gamma_{144}$	$2\pi^+ 2\pi^- \pi^0$	$(4.1 \pm 0.5) \times 10^{-3}$
$\Gamma_{145}$	$\eta \pi^+ \pi^-$	$[h] < 1.9 \times 10^{-3}$ CL=90%
$\Gamma_{146}$	$\omega \pi^+ \pi^-$	$[h] (1.6 \pm 0.5) \times 10^{-3}$
$\Gamma_{147}$	$3\pi^+ 3\pi^-$	$(4.0 \pm 1.1) \times 10^{-4}$

**Hadronic modes with a  $K\bar{K}$  pair**

$\Gamma_{148}$	$K^+ K^-$	$(3.84 \pm 0.10) \times 10^{-3}$
$\Gamma_{149}$	$2K_S^0$	$(3.7 \pm 0.7) \times 10^{-4}$
$\Gamma_{150}$	$K_S^0 K^- \pi^+$	$(3.4 \pm 0.5) \times 10^{-3}$ S=1.1
$\Gamma_{151}$	$\bar{K}^*(892)^0 K_S^0,$ $\bar{K}^*(892)^0 \rightarrow K^- \pi^+$	$< 6 \times 10^{-4}$ CL=90%
$\Gamma_{152}$	$K^*(892)^+ K^-, K^*(892)^+ \rightarrow K_S^0 \pi^+$	$(1.2 \pm 0.3) \times 10^{-3}$
$\Gamma_{153}$	$K_S^0 K^- \pi^+$ nonresonant	$(1.1 \pm 1.1) \times 10^{-3}$
$\Gamma_{154}$	$K_S^0 K^+ \pi^-$	$(2.6 \pm 0.5) \times 10^{-3}$
$\Gamma_{155}$	$K^*(892)^0 K_S^0, K^*(892)^0 \rightarrow K^+ \pi^-$	$< 3 \times 10^{-4}$ CL=90%
$\Gamma_{156}$	$K^*(892)^- K^+, K^*(892)^- \rightarrow K_S^0 \pi^-$	$(7 \pm 4) \times 10^{-4}$

$\Gamma_{157}$	$K_S^0 K^+ \pi^-$ nonresonant	$(1.9 \pm 1.1) \times 10^{-3}$
$\Gamma_{158}$	$K^+ K^- \pi^0$	$(1.3 \pm 0.4) \times 10^{-3}$
$\Gamma_{159}$	$K_S^0 K_S^0 \pi^0$	$< 5.9 \times 10^{-4}$
$\Gamma_{160}$	$K^+ K^- \pi^+ \pi^-$	$[i] (2.32 \pm 0.13) \times 10^{-3}$
$\Gamma_{161}$	$\phi \pi^+ \pi^-$ 3-body, $\phi \rightarrow K^+ K^-$	$(2.3 \pm 2.3) \times 10^{-5}$
$\Gamma_{162}$	$\phi \rho^0$ , $\phi \rightarrow K^+ K^-$	$(6.7 \pm 0.6) \times 10^{-4}$
$\Gamma_{163}$	$K^+ K^- \rho^0$ 3-body	$(5 \pm 7) \times 10^{-5}$
$\Gamma_{164}$	$f_0(980) \pi^+ \pi^-$ , $f_0 \rightarrow K^+ K^-$	$(3.5 \pm 0.9) \times 10^{-4}$
$\Gamma_{165}$	$K^*(892)^0 K^\mp \pi^\pm$ 3-body, $K^{*0} \rightarrow K^\pm \pi^\mp$	$[j] (2.5 \pm 0.5) \times 10^{-4}$
$\Gamma_{166}$	$K^*(892)^0 \bar{K}^*(892)^0$ , $K^{*0} \rightarrow K^\pm \pi^\mp$	$(7 \pm 5) \times 10^{-5}$
$\Gamma_{167}$	$K_1(1270)^\pm K^\mp$ , $K_1(1270)^\pm \rightarrow K^\pm \pi^+ \pi^-$	$(7.6 \pm 1.7) \times 10^{-4}$
$\Gamma_{168}$	$K_1(1400)^\pm K^\mp$ , $K_1(1400)^\pm \rightarrow K^\pm \pi^+ \pi^-$	$(5.1 \pm 1.2) \times 10^{-4}$
$\Gamma_{169}$	$K^+ K^- \pi^+ \pi^-$ non- $\phi$	
$\Gamma_{170}$	$K^+ K^- \pi^+ \pi^-$ nonresonant	
$\Gamma_{171}$	$K_S^0 K_S^0 \pi^+ \pi^-$	$(1.26 \pm 0.24) \times 10^{-3}$
$\Gamma_{172}$	$K_S^0 K^- \pi^+ \pi^+ \pi^-$	$< 1.5 \times 10^{-4}$ CL=90%
$\Gamma_{173}$	$K^+ K^- \pi^+ \pi^- \pi^0$	$(3.1 \pm 2.0) \times 10^{-3}$

Fractions of most of the following modes with resonances have already appeared above as submodes of particular charged-particle modes.

$\Gamma_{174}$	$\bar{K}^*(892)^0 K_S^0$	$< 8 \times 10^{-4}$ CL=90%
$\Gamma_{175}$	$K^*(892)^+ K^-$	$(3.7 \pm 0.8) \times 10^{-3}$
$\Gamma_{176}$	$K^*(892)^0 K_S^0$	$< 4 \times 10^{-4}$ CL=90%
$\Gamma_{177}$	$K^*(892)^- K^+$	$(2.0 \pm 1.1) \times 10^{-3}$
$\Gamma_{178}$	$\phi \pi^0$	$(7.4 \pm 0.5) \times 10^{-4}$
$\Gamma_{179}$	$\phi \eta$	$(1.4 \pm 0.4) \times 10^{-4}$
$\Gamma_{180}$	$\phi \omega$	$< 2.1 \times 10^{-3}$ CL=90%

### Radiative modes

$\Gamma_{181}$	$\rho^0 \gamma$	$< 2.4 \times 10^{-4}$ CL=90%
$\Gamma_{182}$	$\omega \gamma$	$< 2.4 \times 10^{-4}$ CL=90%
$\Gamma_{183}$	$\phi \gamma$	$(2.4 \pm 0.7) \times 10^{-5}$
$\Gamma_{184}$	$\bar{K}^*(892)^0 \gamma$	$< 7.6 \times 10^{-4}$ CL=90%

### Doubly Cabibbo suppressed (DC) modes or $\Delta C = 2$ forbidden via mixing (C2M) modes

$\Gamma_{185}$	$K^+ \ell^- \bar{\nu}_\ell$ (via $\bar{D}^0$ )	$C2M < 1.8 \times 10^{-4}$ CL=90%
$\Gamma_{186}$	$K^+$ or $K^*(892)^+ e^- \bar{\nu}_e$ (via $\bar{D}^0$ )	$C2M < 6 \times 10^{-5}$ CL=90%
$\Gamma_{187}$	$K^+ \pi^-$	$DC (1.43 \pm 0.04) \times 10^{-4}$

$\Gamma_{188}$	$K^+ \pi^-$ (via $\bar{D}^0$ )	$C2M$	< 1.5	$\times 10^{-5}$	CL=95%
$\Gamma_{189}$	$K_S^0 \pi^+ \pi^-$ (in $D^0 \rightarrow \bar{D}^0$ )	$C2M$	< 1.8	$\times 10^{-4}$	CL=95%
$\Gamma_{190}$	$K^*(892)^+ \pi^-$ , $K^*(892)^+ \rightarrow K_S^0 \pi^+$	$DC$	(10 $^{+12}_{-4}$ )	$\times 10^{-5}$	
$\Gamma_{191}$	$K^+ \pi^- \pi^0$	$DC$	( 3.29 $^{+0.30}_{-0.27}$ )	$\times 10^{-4}$	
$\Gamma_{192}$	$K^+ \pi^- \pi^+ \pi^-$	$DC$	( 2.49 $^{+0.21}_{-0.19}$ )	$\times 10^{-4}$	
$\Gamma_{193}$	$K^+ \pi^- \pi^+ \pi^-$ (via $\bar{D}^0$ )	$C2M$	< 4	$\times 10^{-4}$	CL=90%
$\Gamma_{194}$	$K^+ \pi^-$ or $K^+ \pi^- \pi^+ \pi^-$ (via $\bar{D}^0$ )				
$\Gamma_{195}$	$\mu^-$ anything (via $\bar{D}^0$ )	$C2M$	< 4	$\times 10^{-4}$	CL=90%

**$\Delta C = 1$  weak neutral current (*C1*) modes,  
Lepton Family number (*LF*) violating modes, or  
Lepton number (*L*) violating modes**

$\Gamma_{196}$	$\gamma \gamma$	$C1$	< 2.6	$\times 10^{-5}$	CL=90%
$\Gamma_{197}$	$e^+ e^-$	$C1$	< 1.2	$\times 10^{-6}$	CL=90%
$\Gamma_{198}$	$\mu^+ \mu^-$	$C1$	< 1.3	$\times 10^{-6}$	CL=90%
$\Gamma_{199}$	$\pi^0 e^+ e^-$	$C1$	< 4.5	$\times 10^{-5}$	CL=90%
$\Gamma_{200}$	$\pi^0 \mu^+ \mu^-$	$C1$	< 1.8	$\times 10^{-4}$	CL=90%
$\Gamma_{201}$	$\eta e^+ e^-$	$C1$	< 1.1	$\times 10^{-4}$	CL=90%
$\Gamma_{202}$	$\eta \mu^+ \mu^-$	$C1$	< 5.3	$\times 10^{-4}$	CL=90%
$\Gamma_{203}$	$\pi^+ \pi^- e^+ e^-$	$C1$	< 3.73	$\times 10^{-4}$	CL=90%
$\Gamma_{204}$	$\rho^0 e^+ e^-$	$C1$	< 1.0	$\times 10^{-4}$	CL=90%
$\Gamma_{205}$	$\pi^+ \pi^- \mu^+ \mu^-$	$C1$	< 3.0	$\times 10^{-5}$	CL=90%
$\Gamma_{206}$	$\rho^0 \mu^+ \mu^-$	$C1$	< 2.2	$\times 10^{-5}$	CL=90%
$\Gamma_{207}$	$\omega e^+ e^-$	$C1$	< 1.8	$\times 10^{-4}$	CL=90%
$\Gamma_{208}$	$\omega \mu^+ \mu^-$	$C1$	< 8.3	$\times 10^{-4}$	CL=90%
$\Gamma_{209}$	$K^- K^+ e^+ e^-$	$C1$	< 3.15	$\times 10^{-4}$	CL=90%
$\Gamma_{210}$	$\phi e^+ e^-$	$C1$	< 5.2	$\times 10^{-5}$	CL=90%
$\Gamma_{211}$	$K^- K^+ \mu^+ \mu^-$	$C1$	< 3.3	$\times 10^{-5}$	CL=90%
$\Gamma_{212}$	$\phi \mu^+ \mu^-$	$C1$	< 3.1	$\times 10^{-5}$	CL=90%
$\Gamma_{213}$	$\bar{K}^0 e^+ e^-$	[ $k$ ]	< 1.1	$\times 10^{-4}$	CL=90%
$\Gamma_{214}$	$\bar{K}^0 \mu^+ \mu^-$	[ $k$ ]	< 2.6	$\times 10^{-4}$	CL=90%
$\Gamma_{215}$	$K^- \pi^+ e^+ e^-$	$C1$	< 3.85	$\times 10^{-4}$	CL=90%
$\Gamma_{216}$	$\bar{K}^*(892)^0 e^+ e^-$	[ $k$ ]	< 4.7	$\times 10^{-5}$	CL=90%
$\Gamma_{217}$	$K^- \pi^+ \mu^+ \mu^-$	$C1$	< 3.59	$\times 10^{-4}$	CL=90%
$\Gamma_{218}$	$\bar{K}^*(892)^0 \mu^+ \mu^-$	[ $k$ ]	< 2.4	$\times 10^{-5}$	CL=90%
$\Gamma_{219}$	$\pi^+ \pi^- \pi^0 \mu^+ \mu^-$	$C1$	< 8.1	$\times 10^{-4}$	CL=90%
$\Gamma_{220}$	$\mu^\pm e^\mp$	$LF$	[ $ I $ ] < 8.1	$\times 10^{-7}$	CL=90%
$\Gamma_{221}$	$\pi^0 e^\pm \mu^\mp$	$LF$	[ $ I $ ] < 8.6	$\times 10^{-5}$	CL=90%
$\Gamma_{222}$	$\eta e^\pm \mu^\mp$	$LF$	[ $ I $ ] < 1.0	$\times 10^{-4}$	CL=90%
$\Gamma_{223}$	$\pi^+ \pi^- e^\pm \mu^\mp$	$LF$	[ $ I $ ] < 1.5	$\times 10^{-5}$	CL=90%
$\Gamma_{224}$	$\rho^0 e^\pm \mu^\mp$	$LF$	[ $ I $ ] < 4.9	$\times 10^{-5}$	CL=90%

$\Gamma_{225}$	$\omega e^\pm \mu^\mp$	$LF$	$ I  < 1.2$	$\times 10^{-4}$	CL=90%
$\Gamma_{226}$	$K^- K^+ e^\pm \mu^\mp$	$LF$	$ I  < 1.8$	$\times 10^{-4}$	CL=90%
$\Gamma_{227}$	$\phi e^\pm \mu^\mp$	$LF$	$ I  < 3.4$	$\times 10^{-5}$	CL=90%
$\Gamma_{228}$	$\bar{K}^0 e^\pm \mu^\mp$	$LF$	$ I  < 1.0$	$\times 10^{-4}$	CL=90%
$\Gamma_{229}$	$K^- \pi^+ e^\pm \mu^\mp$	$LF$	$ I  < 5.53$	$\times 10^{-4}$	CL=90%
$\Gamma_{230}$	$\bar{K}^*(892)^0 e^\pm \mu^\mp$	$LF$	$ I  < 8.3$	$\times 10^{-5}$	CL=90%
$\Gamma_{231}$	$\pi^- \pi^- e^+ e^+ + \text{c.c.}$	$L$	$< 1.12$	$\times 10^{-4}$	CL=90%
$\Gamma_{232}$	$\pi^- \pi^- \mu^+ \mu^+ + \text{c.c.}$	$L$	$< 2.9$	$\times 10^{-5}$	CL=90%
$\Gamma_{233}$	$K^- \pi^- e^+ e^+ + \text{c.c.}$	$L$	$< 2.06$	$\times 10^{-4}$	CL=90%
$\Gamma_{234}$	$K^- \pi^- \mu^+ \mu^+ + \text{c.c.}$	$L$	$< 3.9$	$\times 10^{-4}$	CL=90%
$\Gamma_{235}$	$K^- K^- e^+ e^+ + \text{c.c.}$	$L$	$< 1.52$	$\times 10^{-4}$	CL=90%
$\Gamma_{236}$	$K^- K^- \mu^+ \mu^+ + \text{c.c.}$	$L$	$< 9.4$	$\times 10^{-5}$	CL=90%
$\Gamma_{237}$	$\pi^- \pi^- e^+ \mu^+ + \text{c.c.}$	$L$	$< 7.9$	$\times 10^{-5}$	CL=90%
$\Gamma_{238}$	$K^- \pi^- e^+ \mu^+ + \text{c.c.}$	$L$	$< 2.18$	$\times 10^{-4}$	CL=90%
$\Gamma_{239}$	$K^- K^- e^+ \mu^+ + \text{c.c.}$	$L$	$< 5.7$	$\times 10^{-5}$	CL=90%

$\Gamma_{240}$  A dummy mode used by the fit.  $(38.0 \pm 1.9) \%$

- [a] This value is obtained by subtracting the branching fractions for 2-, 4- and 6-prongs from unity.
- [b] This is the sum of our  $K^- \pi^+ \pi^+ \pi^-$ ,  $K^- \pi^+ \pi^+ \pi^- \pi^0$ ,  $\bar{K}^0 2\pi^+ 2\pi^-$ ,  $2\pi^+ 2\pi^-$ ,  $2\pi^+ 2\pi^- \pi^0$ ,  $K^+ K^- \pi^+ \pi^-$ , and  $K^+ K^- \pi^+ \pi^- \pi^0$ , branching fractions.
- [c] The branching fractions for the  $K^- e^+ \nu_e$ ,  $K^*(892)^- e^+ \nu_e$ ,  $\pi^- e^+ \nu_e$ , and  $\rho^- e^+ \nu_e$  modes add up to  $6.14 \pm 0.20 \%$ .
- [d] This is a weighted average of  $D^\pm$  (44%) and  $D^0$  (56%) branching fractions. See “ $D^+$  and  $D^0 \rightarrow (\eta \text{ anything}) / (\text{total } D^+ \text{ and } D^0)$ ” under “ $D^+$  Branching Ratios” in these Particle Listings.
- [e] The branching fraction for this mode may differ from the sum of the submodes that contribute to it, due to interference effects. See the relevant papers.
- [f] This is a doubly Cabibbo-suppressed mode.
- [g] The two experiments measuring this fraction are in serious disagreement. See the Particle Listings.
- [h] This branching fraction includes all the decay modes of the resonance in the final state.
- [i] The experiments on the division of this charge mode amongst its submodes disagree, and the submode branching fractions here add up to considerably more than the charged-mode fraction.
- [j] However, these upper limits are in serious disagreement with values obtained in another experiment.

- [k] This mode is not a useful test for a  $\Delta C=1$  weak neutral current because both quarks must change flavor in this decay.
- [l] The value is for the sum of the charge states or particle/antiparticle states indicated.

## CONSTRAINED FIT INFORMATION

An overall fit to 41 branching ratios uses 78 measurements and one constraint to determine 22 parameters. The overall fit has a  $\chi^2 = 42.6$  for 57 degrees of freedom.

The following *off-diagonal* array elements are the correlation coefficients  $\langle \delta x_i \delta x_j \rangle / (\delta x_i \cdot \delta x_j)$ , in percent, from the fit to the branching fractions,  $x_i \equiv \Gamma_i / \Gamma_{\text{total}}$ . The fit constrains the  $x_i$  whose labels appear in this array to sum to one.

$x_{15}$	2									
$x_{16}$	26	9								
$x_{17}$	0	0	1							
$x_{26}$	1	23	2	0						
$x_{29}$	9	26	35	2	6					
$x_{30}$	1	2	2	10	0	6				
$x_{31}$	1	3	3	16	1	10	63			
$x_{43}$	3	7	10	0	2	28	2	3		
$x_{55}$	4	10	14	1	2	38	3	4	57	
$x_{64}$	0	1	2	7	0	5	29	46	2	3
$x_{72}$	2	4	6	0	1	16	1	2	15	25
$x_{86}$	0	1	2	7	0	4	58	45	1	2
$x_{87}$	0	1	1	5	0	3	21	33	1	2
$x_{93}$	1	2	2	0	0	6	0	1	9	16
$x_{95}$	0	1	1	1	0	4	2	4	5	9
$x_{106}$	0	1	1	3	0	3	11	18	2	3
$x_{140}$	4	12	16	1	3	46	3	5	26	41
$x_{150}$	1	1	2	6	0	6	23	37	2	2
$x_{154}$	0	1	2	4	0	5	16	26	1	2
$x_{175}$	0	1	1	4	0	3	18	28	1	1
$x_{240}$	-39	-15	-28	-15	-4	-36	-33	-46	-48	-49
	$x_6$	$x_{15}$	$x_{16}$	$x_{17}$	$x_{26}$	$x_{29}$	$x_{30}$	$x_{31}$	$x_{43}$	$x_{55}$

$x_{72}$	1								
$x_{86}$	21	1							
$x_{87}$	40	1	15						
$x_{93}$	0	4	0	0					
$x_{95}$	8	2	2	3	1				
$x_{106}$	38	1	8	15	0	3			
$x_{140}$	3	13	2	2	7	4	2		
$x_{150}$	17	1	16	12	0	2	6	3	
$x_{154}$	12	1	12	9	0	1	5	2	9
$x_{175}$	13	1	13	9	0	1	5	2	10
$x_{240}$	-57	-33	-25	-35	-24	-36	-37	-27	-20
									-15
	$x_{64}$	$x_{72}$	$x_{86}$	$x_{87}$	$x_{93}$	$x_{95}$	$x_{106}$	$x_{140}$	$x_{150}$
$x_{240}$									-17
									$x_{175}$

## CONSTRAINED FIT INFORMATION

An overall fit to 3 branching ratios uses 3 measurements and one constraint to determine 4 parameters. The overall fit has a  $\chi^2 = 0.0$  for 0 degrees of freedom.

The following *off-diagonal* array elements are the correlation coefficients  $\langle \delta x_i \delta x_j \rangle / (\delta x_i \cdot \delta x_j)$ , in percent, from the fit to the branching fractions,  $x_i \equiv \Gamma_i / \Gamma_{\text{total}}$ . The fit constrains the  $x_i$  whose labels appear in this array to sum to one.

$x_2$	-100			
$x_3$	-48	41		
$x_4$	-2	0	0	
	$x_1$	$x_2$	$x_3$	

## $D^0$ BRANCHING RATIOS

Some older now obsolete results have been omitted from these Listings.

### Topological modes

#### $\Gamma(0\text{-prongs})/\Gamma_{\text{total}}$

This value is obtained by subtracting the branching fractions for 2-, 4-, and 6-prongs from unity.

VALUE  
 **$0.19 \pm 0.06$  OUR FIT**

DOCUMENT ID

#### $\Gamma_1/\Gamma$

### $\Gamma(4\text{-prongs})/\Gamma_{\text{total}}$

$\Gamma_3/\Gamma$

This is the sum of our  $K^-\pi^+\pi^+\pi^-$ ,  $K^-\pi^+\pi^+\pi^-\pi^0$ ,  $\bar{K}^02\pi^+2\pi^-$ ,  $2\pi^+2\pi^-$ ,  $2\pi^+2\pi^-\pi^0$ ,  $K^+K^-\pi^+\pi^-$ , and  $K^+K^-\pi^+\pi^-\pi^0$  branching fractions.

VALUE	DOCUMENT ID
<b>0.138±0.005 OUR FIT</b>	
<b>0.138±0.005</b>	PDG 06



### $\Gamma(4\text{-prongs})/\Gamma(2\text{-prongs})$

$\Gamma_3/\Gamma_2$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.207±0.016 OUR FIT</b>				
<b>0.207±0.016±0.004</b>	226	ONENGUT	05 CHRS	$\nu_\mu$ emulsion, $\bar{E}_\nu \approx 27$ GeV



### $\Gamma(6\text{-prongs})/\Gamma_{\text{total}}$

$\Gamma_4/\Gamma$

VALUE (units $10^{-3}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>1.2<sup>+1.3</sup><sub>-0.7</sub> OUR FIT</b>				
<b>1.2<sup>+1.3</sup><sub>-0.9</sub>±0.2</b>	3	ONENGUT	05 CHRS	$\nu_\mu$ emulsion, $\bar{E}_\nu \approx 27$ GeV



### Inclusive modes

### $\Gamma(e^+\text{ anything})/\Gamma_{\text{total}}$

$\Gamma_5/\Gamma$

The branching fractions for the  $K^-e^+\nu_e$ ,  $K^*(892)^-e^+\nu_e$ ,  $\pi^-e^+\nu_e$ , and  $\rho^-e^+\nu_e$  modes add up to  $6.14 \pm 0.20$  %.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.0671±0.0029 OUR AVERAGE</b>				
0.069 ± 0.003 ± 0.005	1670	ALBRECHT	96C ARG	$e^+e^- \approx 10$ GeV
0.0664±0.0018±0.0029	4609	23 KUBOTA	96B CLE2	$e^+e^- \approx \gamma(4S)$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.15 ± 0.05		AGUILAR-...	87E HYBR	$\pi p, pp$ 360, 400 GeV
0.075 ± 0.011 ± 0.004	137	BALTRUSAIT..85B	MRK3	$e^+e^-$ 3.77 GeV
0.055 ± 0.037	12	SCHINDLER	81 MRK2	$e^+e^-$ 3.771 GeV
23 KUBOTA 96B uses $D^{*+} \rightarrow D^0\pi^+$ (and charge conjugate) events in which the $D^0$ subsequently decays to $X e^+\nu_e$ .				



### $\Gamma(\mu^+\text{ anything})/\Gamma_{\text{total}}$

$\Gamma_6/\Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.065±0.007 OUR FIT</b>				
<b>0.063±0.009 OUR AVERAGE</b>				
0.065±0.012±0.003	36	KAYIS-TOPAK.05	CHRS	$\nu_\mu$ emulsion
0.060±0.007±0.012	310	ALBRECHT	96C ARG	$e^+e^- \approx 10$ GeV



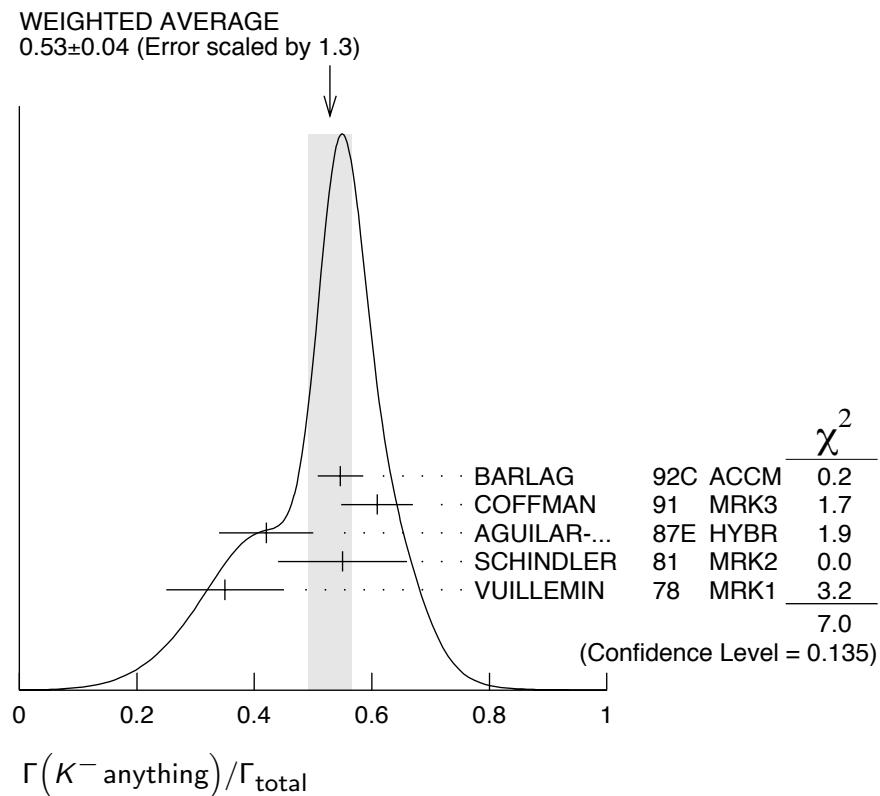
### $\Gamma(K^-\text{ anything})/\Gamma_{\text{total}}$

$\Gamma_7/\Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.53 ± 0.04 OUR AVERAGE</b>				Error includes scale factor of 1.3. See the ideogram below.
0.546 <sup>+0.039</sup> <sub>-0.038</sub>		24 BARLAG	92C ACCM	$\pi^-$ Cu 230 GeV
0.609±0.032±0.052		COFFMAN	91 MRK3	$e^+e^-$ 3.77 GeV
0.42 ± 0.08		AGUILAR-...	87E HYBR	$\pi p, pp$ 360, 400 GeV
0.55 ± 0.11	121	SCHINDLER	81 MRK2	$e^+e^-$ 3.771 GeV
0.35 ± 0.10	19	VUILLEMIN	78 MRK1	$e^+e^-$ 3.772 GeV



<sup>24</sup> BARLAG 92C computes the branching fraction using topological normalization.



$$[\Gamma(\bar{K}^0 \text{anything}) + \Gamma(K^0 \text{anything})]/\Gamma_{\text{total}}$$

$$\Gamma_8/\Gamma$$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.42 ±0.05 OUR AVERAGE</b>				
0.455±0.050±0.032		COFFMAN 91	MRK3	$e^+ e^-$ 3.77 GeV
0.29 ±0.11	13	SCHINDLER 81	MRK2	$e^+ e^-$ 3.771 GeV
0.57 ±0.26	6	VUILLEMIN 78	MRK1	$e^+ e^-$ 3.772 GeV

$$\Gamma(K^+ \text{anything})/\Gamma_{\text{total}}$$

$$\Gamma_9/\Gamma$$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.034<sup>+0.006</sup><sub>-0.004</sub> OUR AVERAGE</b>				
0.034 <sup>+0.007</sup> <sub>-0.005</sub>		25 BARLAG	92C ACCM	$\pi^-$ Cu 230 GeV
0.028±0.009±0.004		COFFMAN 91	MRK3	$e^+ e^-$ 3.77 GeV
0.03 <sup>+0.05</sup> <sub>-0.02</sub>		AGUILAR-...	87E HYBR	$\pi p, pp$ 360, 400 GeV
0.08 ±0.03	25	SCHINDLER 81	MRK2	$e^+ e^-$ 3.771 GeV

<sup>25</sup> BARLAG 92C computes the branching fraction using topological normalization.

$$\Gamma(\bar{K}^*(892)^0 \text{anything})/\Gamma_{\text{total}}$$

$$\Gamma_{10}/\Gamma$$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.087±0.040±0.012</b>	96 ± 44	ABLIKIM 05P	BES	$e^+ e^-$ ≈ 3773 MeV

$\Gamma(K^*(892)^0 \text{anything})/\Gamma_{\text{total}}$ 

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	$\Gamma_{11}/\Gamma$
<b>0.028±0.012±0.004</b>	$31 \pm 12$	ABLIKIM	05P BES	$e^+ e^- \approx 3773 \text{ MeV}$	

 $\Gamma(\phi \text{ anything})/\Gamma_{\text{total}}$ 

<u>VALUE (units <math>10^{-2}</math>)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	$\Gamma_{13}/\Gamma$
<b>1.71<math>^{+0.76}_{-0.71}</math>±0.17</b>	9	26 BAI	00C BES	$e^+ e^- \rightarrow D\bar{D}^*, D^*\bar{D}^*$	

<sup>26</sup> BAI 00C finds the average ( $\phi$  anything) branching fraction for the 4.03-GeV mix of  $D^+$  and  $D^0$  mesons to be  $(1.34 \pm 0.52 \pm 0.12)\%$ .

**Semileptonic modes** $\Gamma(K^- e^+ \nu_e)/\Gamma_{\text{total}}$ 

<u>VALUE (units <math>10^{-2}</math>)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	$\Gamma_{15}/\Gamma$
<b>3.51±0.11 OUR FIT</b>					
<b>3.47±0.13 OUR AVERAGE</b>					
3.44±0.10±0.10	$1311 \pm 37$	COAN	05 CLEO	$e^+ e^- \text{ at } \psi(3770)$	
3.82±0.40±0.27	$104 \pm 11$	ABLIKIM	04C BES	$e^+ e^-, 3.773 \text{ GeV}$	
3.4 ± 0.5 ± 0.4	55	ADLER	89 MRK3	$e^+ e^- 3.77 \text{ GeV}$	

 $\Gamma(K^- e^+ \nu_e)/\Gamma(K^- \pi^+)$ 

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	$\Gamma_{15}/\Gamma_{29}$
<b>0.923±0.029 OUR FIT</b>					
<b>0.95 ± 0.04 OUR AVERAGE</b>					
0.978±0.027±0.044	2510	27 BEAN	93C CLE2	$e^+ e^- \approx \gamma(4S)$	
0.90 ± 0.06 ± 0.06	584	28 CRAWFORD	91B CLEO	$e^+ e^- \approx 10.5 \text{ GeV}$	
0.91 ± 0.07 ± 0.11	250	29 ANJOS	89F E691	Photoproduction	

<sup>27</sup> BEAN 93C uses  $K^- \mu^+ \nu_\mu$  as well as  $K^- e^+ \nu_e$  events and makes a small phase-space adjustment to the number of the  $\mu^+$  events to use them as  $e^+$  events. A pole mass of  $2.00 \pm 0.12 \pm 0.18 \text{ GeV}/c^2$  is obtained from the  $q^2$  dependence of the decay rate.

<sup>28</sup> CRAWFORD 91B uses  $K^- e^+ \nu_e$  and  $K^- \mu^+ \nu_\mu$  candidates to measure a pole mass of  $2.1^{+0.4+0.3}_{-0.2-0.2} \text{ GeV}/c^2$  from the  $q^2$  dependence of the decay rate.

<sup>29</sup> ANJOS 89F measures a pole mass of  $2.1^{+0.4}_{-0.2} \pm 0.2 \text{ GeV}/c^2$  from the  $q^2$  dependence of the decay rate.

 $\Gamma(K^- \mu^+ \nu_\mu)/\Gamma(K^- \pi^+)$ 

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	$\Gamma_{16}/\Gamma_{29}$
<b>0.84 ± 0.04 OUR FIT</b>					
<b>0.84 ± 0.04 OUR AVERAGE</b>					

0.852±0.034±0.028	1897	30 FRABETTI	95G E687	$\gamma \text{Be } \bar{E}_\gamma = 220 \text{ GeV}$	
0.82 ± 0.13 ± 0.13	338	31 FRABETTI	93I E687	$\gamma \text{Be } \bar{E}_\gamma = 221 \text{ GeV}$	
0.79 ± 0.08 ± 0.09	231	32 CRAWFORD	91B CLEO	$e^+ e^- \approx 10.5 \text{ GeV}$	

<sup>30</sup> FRABETTI 95G extracts the ratio of form factors  $f_-(0)/f_+(0) = -1.3^{+3.6}_{-3.4} \pm 0.6$ , and measures a pole mass of  $1.87^{+0.11+0.07}_{-0.08-0.06} \text{ GeV}/c^2$  from the  $q^2$  dependence of the decay rate.

<sup>31</sup> FRABETTI 93I measures a pole mass of  $2.1^{+0.7+0.7}_{-0.3-0.3} \text{ GeV}/c^2$  from the  $q^2$  dependence of the decay rate.

<sup>32</sup> CRAWFORD 91B measures a pole mass of  $2.00 \pm 0.12 \pm 0.18 \text{ GeV}/c^2$  from the  $q^2$  dependence of the decay rate.

### $\Gamma(K^-\mu^+\nu_\mu)/\Gamma(\mu^+ \text{anything})$

$\Gamma_{16}/\Gamma_6$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.49 ± 0.05 OUR FIT</b>				
<b>0.472 ± 0.051 ± 0.040</b>	232	KODAMA	94 E653	$\pi^-$ emulsion 600 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.32 ± 0.05 ± 0.05	124	KODAMA	91 EMUL	pA 800 GeV

### $\Gamma(K^-\pi^0e^+\nu_e)/\Gamma_{\text{total}}$

$\Gamma_{19}/\Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				

$0.016^{+0.013}_{-0.005} \pm 0.002$       4      33 BAI      91 MRK3       $e^+ e^- \approx 3.77$  GeV

33 BAI 91 finds that a fraction  $0.79^{+0.15}_{-0.17}{}^{+0.09}_{-0.03}$  of combined  $D^+$  and  $D^0$  decays to  $\bar{K}\pi e^+\nu_e$  (24 events) are  $\bar{K}^*(892)e^+\nu_e$ . BAI 91 uses 56  $K^-e^+\nu_e$  events to measure a pole mass of  $1.8 \pm 0.3 \pm 0.2$  GeV/ $c^2$  from the  $q^2$  dependence of the decay rate.

### $\Gamma(\bar{K}^0\pi^-e^+\nu_e)/\Gamma_{\text{total}}$

$\Gamma_{20}/\Gamma$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				

$0.028^{+0.017}_{-0.008} \pm 0.003$       6      34 BAI      91 MRK3       $e^+ e^- \approx 3.77$  GeV

34 BAI 91 finds that a fraction  $0.79^{+0.15}_{-0.17}{}^{+0.09}_{-0.03}$  of combined  $D^+$  and  $D^0$  decays to  $\bar{K}\pi e^+\nu_e$  (24 events) are  $\bar{K}^*(892)e^+\nu_e$ .

### $\Gamma(K^*(892)^-e^+\nu_e)/\Gamma_{\text{total}}$

$\Gamma_{17}/\Gamma$

Both decay modes of the  $K^*(892)^-$  are included.

VALUE (units $10^{-2}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>2.17 ± 0.16 OUR FIT</b>				

**2.16 ± 0.15 ± 0.08**       $219 \pm 16$       35 COAN      05 CLEO       $e^+ e^-$  at  $\psi(3770)$

35 COAN 05 uses both  $K^-\pi^0$  and  $K_S^0\pi^-$  events.

### $\Gamma(K^*(892)^-e^+\nu_e)/\Gamma(K_S^0\pi^+\pi^-)$

$\Gamma_{17}/\Gamma_{31}$

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.75 ± 0.07 OUR FIT</b>				

**0.76 ± 0.12 ± 0.06**      152      36 BEAN      93C CLE2       $e^+ e^- \approx \Upsilon(4S)$

36 BEAN 93C uses  $K^-\mu^+\nu_\mu$  as well as  $K^-\pi^+\nu_e$  events and makes a small phase-space adjustment to the number of the  $\mu^+$  events to use them as  $e^+$  events.

### $\Gamma(K^*(892)^-e^+\nu_e)/\Gamma(K^-\pi^+\nu_e)$

$\Gamma_{17}/\Gamma_{15}$

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				

$0.51 \pm 0.18 \pm 0.06$       CRAWFORD      91B CLEO       $e^+ e^- \approx 10.5$  GeV

$\Gamma(K^*(892)^-\mu^+\nu_\mu)/\Gamma(K_S^0\pi^+\pi^-)$   $\Gamma_{18}/\Gamma_{31}$ 

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.674±0.068±0.026</b>	$175 \pm 17$	37 LINK	05B FOCS	$\gamma A, \bar{E}_\gamma \approx 180$ GeV

37 LINK 05B finds that in  $D^0 \rightarrow \bar{K}^0\pi^-\mu^+\nu_\mu$  the  $\bar{K}^0\pi^-$  system is 6% in  $S$ -wave.

 $\Gamma(K^*(892)^-\ell^+\nu_\ell)/\Gamma(K_S^0\pi^+\pi^-)$   $\Gamma_{22}/\Gamma_{31}$ 

This an average of the  $K^*(892)^-e^+\nu_e$  and  $K^*(892)^-\mu^+\nu_\mu$  ratios. Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
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• • • We do not use the following data for averages, fits, limits, etc. • • •

$0.48 \pm 0.14 \pm 0.12$  137 38 ALEXANDER 90B CLEO  $e^+e^-$  10.5–11 GeV

38 ALEXANDER 90B cannot exclude extra  $\pi^0$ 's in the final state.

 $\Gamma(K^-\pi^+\pi^-\mu^+\nu_\mu)/\Gamma(K^-\mu^+\nu_\mu)$   $\Gamma_{24}/\Gamma_{16}$ 

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.037</b>	90	KODAMA	93B E653	$\pi^-$ emulsion 600 GeV

 $\Gamma((\bar{K}^*(892)\pi)^-\mu^+\nu_\mu)/\Gamma(K^-\mu^+\nu_\mu)$   $\Gamma_{25}/\Gamma_{16}$ 

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.043</b>	90	39 KODAMA	93B E653	$\pi^-$ emulsion 600 GeV

39 KODAMA 93B searched in  $K^-\pi^+\pi^-\mu^+\nu_\mu$ , but the limit includes other  $(\bar{K}^*(892)\pi)^-$  charge states.

 $\Gamma(\pi^-e^+\nu_e)/\Gamma_{\text{total}}$   $\Gamma_{26}/\Gamma$ 

VALUE (units $10^{-2}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.281±0.019 OUR FIT</b>				

**0.262±0.025±0.008**  $117 \pm 11$  COAN 05 CLEO  $e^+e^-$  at  $\psi(3770)$

• • • We do not use the following data for averages, fits, limits, etc. • • •

$0.33 \pm 0.13 \pm 0.03$   $9 \pm 4$  40 ABLIKIM 04C BES  $e^+e^-$ , 3.773 GeV

$0.39^{+0.23}_{-0.11} \pm 0.04$   $7$  41 ADLER 89 MRK3  $e^+e^-$  3.77 GeV

40 ABLIKIM 04C measures  $|\frac{f_+^\pi(0)}{f_+^K(0)}|$  to be  $0.93 \pm 0.19 \pm 0.07$ .

41 This result of ADLER 89 gives  $|\frac{V_{cd}}{V_{cs}} \cdot \frac{f_+^\pi(0)}{f_+^K(0)}|^2 = 0.057^{+0.038}_{-0.015} \pm 0.005$ .

 $\Gamma(\pi^-e^+\nu_e)/\Gamma(K^-\ell^+\nu_\ell)$   $\Gamma_{26}/\Gamma_{15}$ 

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.080±0.005 OUR FIT</b>				

**0.085±0.007 OUR AVERAGE**

$0.082 \pm 0.006 \pm 0.005$  42 HUANG 05 CLEO  $e^+e^- \approx \gamma(4S)$

$0.101 \pm 0.020 \pm 0.003$  91 43 FRABETTI 96B E687  $\gamma Be, \bar{E}_\gamma \approx 200$  GeV

$0.103 \pm 0.039 \pm 0.013$  87 44 BUTLER 95 CLE2  $< 0.156$  (90% CL)

<sup>42</sup>HUANG 05 uses both  $e$  and  $\mu$  events, and makes a small correction to the  $\mu$  events to make them effectively  $e$  events. This result gives  $|\frac{V_{cd}}{V_{cs}} \cdot \frac{f_+^\pi(0)}{f_+^K(0)}|^2 = 0.038^{+0.006+0.005}_{-0.007-0.003}$ .

<sup>43</sup>FRABETTI 96B uses both  $e$  and  $\mu$  events, and makes a small correction to the  $\mu$  events to make them effectively  $e$  events. This result gives  $|\frac{V_{cd}}{V_{cs}} \cdot \frac{f_+^\pi(0)}{f_+^K(0)}|^2 = 0.050 \pm 0.011 \pm 0.002$ .

<sup>44</sup>BUTLER 95 has  $87 \pm 33 \pi^- e^+ \nu_e$  events. The result gives  $|\frac{V_{cd}}{V_{cs}} \cdot \frac{f_+^\pi(0)}{f_+^K(0)}|^2 = 0.052 \pm 0.020 \pm 0.007$ .

### $\Gamma(\pi^- \mu^+ \nu_\mu)/\Gamma(K^- \mu^+ \nu_\mu)$

$\Gamma_{27}/\Gamma_{16}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.074±0.008±0.007</b>	$288 \pm 29$	45 LINK	05 FOCS	$\gamma A, \bar{E}_\gamma \approx 180 \text{ GeV}$

<sup>45</sup>LINK 05 finds the form-factor ratio  $|f_0^\pi(0)/f_0^K(0)|$  to be  $0.85 \pm 0.04 \pm 0.04 \pm 0.01$ .

### $\Gamma(\rho^- e^+ \nu_e)/\Gamma_{\text{total}}$

$\Gamma_{28}/\Gamma$

VALUE (units $10^{-2}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.194±0.039±0.013</b>	$31 \pm 6$	COAN	05 CLEO	$e^+ e^-$ at $\psi(3770)$

## Hadronic modes with a single $\bar{K}$

### $\Gamma(K^- \pi^+)/\Gamma_{\text{total}}$

$\Gamma_{29}/\Gamma$

VALUE (units $10^{-2}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>3.80±0.07 OUR FIT</b>	Error includes scale factor of 1.1.			
<b>3.85±0.07 OUR AVERAGE</b>				
3.91±0.08±0.09	$10.3k \pm 100$	46 HE	05 CLEO	$e^+ e^-$ at $\psi(3770)$
3.82±0.07±0.12		47 ARTUSO	98 CLE2	CLEO average
3.90±0.09±0.12	5392	48 BARATE	97C ALEP	From $Z$ decays
3.41±0.12±0.28	$1173 \pm 37$	48 ALBRECHT	94F ARG	$e^+ e^- \approx \gamma(4S)$
3.62±0.34±0.44		48 DECOMP	91J ALEP	From $Z$ decays

• • • We do not use the following data for averages, fits, limits, etc. • • •

3.81±0.15±0.16	1165	49 ARTUSO	98 CLE2	$e^+ e^-$ at $\gamma(4S)$
3.69±0.11±0.16		50 COAN	98 CLE2	See ARTUSO 98
4.5 ± 0.6 ± 0.4		51 ALBRECHT	94 ARG	$e^+ e^- \approx \gamma(4S)$
3.95±0.08±0.17	4208	48,52 AKERIB	93 CLE2	See ARTUSO 98
4.5 ± 0.8 ± 0.5	56	48 ABACHI	88 HRS	$e^+ e^-$ 29 GeV
4.2 ± 0.4 ± 0.4	930	ADLER	88C MRK3	$e^+ e^-$ 3.77 GeV
4.1 ± 0.6	$263 \pm 17$	53 SCHINDLER	81 MRK2	$e^+ e^-$ 3.771 GeV
4.3 ± 1.0	130	54 PERUZZI	77 MRK1	$e^+ e^-$ 3.77 GeV

<sup>46</sup>HE 05 uses single- and double-tagged events in an overall fit. The fraction here includes (unobserved) final-state photons.

<sup>47</sup>This combines the CLEO results of ARTUSO 98, COAN 98, and AKERIB 93.

<sup>48</sup>ABACHI 88, DECOMP 91J, AKERIB 93, ALBRECHT 94F, and BARATE 97C use  $D^*(2010)^+ \rightarrow D^0 \pi^+$  decays. The  $\pi^+$  is both slow and of low  $p_T$  with respect to the event thrust axis or nearest jet ( $\approx D^{*+}$  direction). The excess number of such  $\pi^+$ 's over background gives the number of  $D^*(2010)^+ \rightarrow D^0 \pi^+$  events, and the fraction with  $D^0 \rightarrow K^- \pi^+$  gives the  $D^0 \rightarrow K^- \pi^+$  branching fraction.

<sup>49</sup> ARTUSO 98, following ALBRECHT 94, uses  $D^0$  mesons from  $\bar{B}^0 \rightarrow D^*(2010)^+ X \ell^- \bar{\nu}_\ell$  decays. Our average uses the CLEO average of this value with the values of COAN 98 and AKERIB 93.

<sup>50</sup> COAN 98 assumes that  $\Gamma(B \rightarrow \bar{D} X \ell^+ \nu)/\Gamma(B \rightarrow X \ell^+ \nu) = 1.0 - 3|V_{ub}/V_{cb}|^2 - 0.010 \pm 0.005$ , the last term accounting for  $\bar{B} \rightarrow D_s^+ K X \ell^- \bar{\nu}$ . COAN 98 is included in the CLEO average in ARTUSO 98.

<sup>51</sup> ALBRECHT 94 uses  $D^0$  mesons from  $\bar{B}^0 \rightarrow D^{*+} \ell^- \bar{\nu}_\ell$  decays. This is a different set of events than used by ALBRECHT 94F.

<sup>52</sup> This AKERIB 93 value includes radiative corrections; without them, the value is  $0.0391 \pm 0.0008 \pm 0.0017$ . AKERIB 93 is included in the CLEO average in ARTUSO 98.

<sup>53</sup> SCHINDLER 81 (MARK-2) measures  $\sigma(e^+ e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.24 \pm 0.02$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

<sup>54</sup> PERUZZI 77 (MARK-1) measures  $\sigma(e^+ e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.25 \pm 0.05$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

$\Gamma(K_S^0 \pi^0)/\Gamma(K^- \pi^+)$		$\Gamma_{30}/\Gamma_{29}$		
VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.300±0.031 OUR FIT</b>				
<b>0.68 ± 0.12 ± 0.11</b>	119	ANJOS	92B E691	$\gamma$ Be 80–240 GeV

$\Gamma(K_S^0 \pi^0)/\Gamma(K_S^0 \pi^+ \pi^-)$		$\Gamma_{30}/\Gamma_{31}$		
VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.393±0.033 OUR FIT</b>		Error includes scale factor of 1.1.		
<b>0.378±0.033 OUR AVERAGE</b>				
0.44 ± 0.02 ± 0.05	1942 ± 64	PROCARIO	93B CLE2	$e^+ e^-$ 10.36–10.7 GeV
0.34 ± 0.04 ± 0.02	92	<sup>55</sup> ALBRECHT	92P ARG	$e^+ e^- \approx$ 10 GeV
0.36 ± 0.04 ± 0.08	104	KINOSHITA	91 CLEO	$e^+ e^- \sim$ 10.7 GeV

<sup>55</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(K_S^0 \pi^+ \pi^-)/\Gamma_{\text{total}}$		$\Gamma_{31}/\Gamma$		
VALUE (units $10^{-2}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>2.90±0.19 OUR FIT</b>				
<b>2.68±0.29 OUR AVERAGE</b>				
2.52 ± 0.20 ± 0.25	284 ± 22	<sup>56</sup> ALBRECHT	94F ARG	$e^+ e^- \approx \gamma(4S)$
3.2 ± 0.3 ± 0.5		ADLER	87 MRK3	$e^+ e^-$ 3.77 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
2.6 ± 0.8	32 ± 8	<sup>57</sup> SCHINDLER	81 MRK2	$e^+ e^-$ 3.771 GeV
4.0 ± 1.2	28	<sup>58</sup> PERUZZI	77 MRK1	$e^+ e^-$ 3.77 GeV

<sup>56</sup> See the footnote on the ALBRECHT 94F measurement of  $\Gamma(K^- \pi^+)/\Gamma_{\text{total}}$  for the method used.

<sup>57</sup> SCHINDLER 81 (MARK-2) measures  $\sigma(e^+ e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.30 \pm 0.08$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

<sup>58</sup> PERUZZI 77 (MARK-1) measures  $\sigma(e^+ e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.46 \pm 0.12$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

$\Gamma(K_S^0 \pi^+ \pi^-)/\Gamma(K^- \pi^+)$		$\Gamma_{31}/\Gamma_{29}$		
VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.76±0.05 OUR FIT</b>				
<b>0.81±0.05±0.08</b>	856 ± 35	FRAEBETTI	94J E687	$\gamma$ Be $\bar{E}_\gamma=220$ GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.85 ± 0.40	35	AVERY	80 SPEC	$\gamma N \rightarrow D^{*+}$
1.4 ± 0.5	116	PICCOLO	77 MRK1	$e^+ e^-$ 4.03, 4.41 GeV

$\Gamma(K_S^0 \rho^0)/\Gamma(K_S^0 \pi^+ \pi^-)$  $\Gamma_{32}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.259<sup>+0.014</sup><sub>-0.023</sub> OUR AVERAGE</b>	Error includes scale factor of 1.1.		
0.264 $\pm 0.009^{+0.010}_{-0.026}$	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
0.350 $\pm 0.028^{+0.067}_{-0.067}$	FRABETTI 94G	E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV
0.227 $\pm 0.032^{+0.009}_{-0.009}$	ALBRECHT 93D	ARG	$e^+ e^- \approx 10$ GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$			
0.267 $\pm 0.011^{+0.009}_{-0.028}$	ASNER 04A	CLEO	See MURAMATSU 02
0.215 $\pm 0.051^{+0.037}_{-0.037}$	ANJOS 93	E691	$\gamma$ Be 90–260 GeV
0.20 $\pm 0.06 \pm 0.03$	FRABETTI 92B	E687	$\gamma$ Be, $\bar{E}_\gamma = 221$ GeV
0.12 $\pm 0.01 \pm 0.07$	ADLER 87	MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K_S^0 \omega, \omega \rightarrow \pi^+ \pi^-)/\Gamma(K_S^0 \pi^+ \pi^-)$  $\Gamma_{33}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.0072<sup>+0.0018</sup><sub>-0.0009</sub> OUR AVERAGE</b>	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$			
0.0081 $\pm 0.0019^{+0.0018}_{-0.0010}$	ASNER 04A	CLEO	See MURAMATSU 02

 $\Gamma(K_S^0 f_0(980), f_0(980) \rightarrow \pi^+ \pi^-)/\Gamma(K_S^0 \pi^+ \pi^-)$  $\Gamma_{34}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.047<sup>+0.010</sup><sub>-0.007</sub> OUR AVERAGE</b>			
0.043 $\pm 0.005^{+0.012}_{-0.006}$	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
0.068 $\pm 0.016^{+0.018}_{-0.018}$	FRABETTI 94G	E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV
0.046 $\pm 0.018^{+0.006}_{-0.006}$	ALBRECHT 93D	ARG	$e^+ e^- \approx 10$ GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$			
0.042 $\pm 0.005^{+0.011}_{-0.005}$	ASNER 04A	CLEO	See MURAMATSU 02

 $\Gamma(K_S^0 f_2(1270), f_2(1270) \rightarrow \pi^+ \pi^-)/\Gamma(K_S^0 \pi^+ \pi^-)$  $\Gamma_{35}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis. Note the large difference between the CLEO results and earlier measurements.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.0045<sup>+0.0039</sup><sub>-0.0022</sub> OUR AVERAGE</b>			
0.0027 $\pm 0.0015^{+0.0037}_{-0.0017}$	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
0.037 $\pm 0.014 \pm 0.017$	FRABETTI 94G	E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV
0.050 $\pm 0.021 \pm 0.008$	ALBRECHT 93D	ARG	$e^+ e^- \approx 10$ GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$			
0.0036 $\pm 0.0022^{+0.0032}_{-0.0019}$	ASNER 04A	CLEO	See MURAMATSU 02

$\Gamma(K_S^0 f_0(1370), f_0(1370) \rightarrow \pi^+ \pi^-)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{36}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.085<sup>+0.019</sup><sub>-0.021</sub> OUR AVERAGE</b>			
0.099 $\pm 0.011^{+0.028}_{-0.044}$	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
0.077 $\pm 0.022^{+0.031}_{-0.031}$	FRABETTI 94G	E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV
0.082 $\pm 0.028^{+0.013}_{-0.013}$	ALBRECHT 93D	ARG	$e^+ e^- \approx 10$ GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$			
0.098 $\pm 0.014^{+0.026}_{-0.036}$	ASNER 04A	CLEO	See MURAMATSU 02

 $\Gamma(K^*(892)^- \pi^+, K^*(892)^- \rightarrow K_S^0 \pi^-)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{37}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.660<sup>+0.019</sup><sub>-0.026</sub> OUR AVERAGE</b>			
0.657 $\pm 0.013^{+0.018}_{-0.040}$	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
0.625 $\pm 0.036^{+0.026}_{-0.026}$	FRABETTI 94G	E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV
0.718 $\pm 0.042^{+0.030}_{-0.030}$	ALBRECHT 93D	ARG	$e^+ e^- \approx 10$ GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$			
0.663 $\pm 0.013^{+0.024}_{-0.043}$	ASNER 04A	CLEO	See MURAMATSU 02
0.480 $\pm 0.097$	ANJOS 93	E691	$\gamma$ Be 90–260 GeV
0.56 $\pm 0.04$ $\pm 0.05$	ADLER 87	MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K^*(892)^+ \pi^-, K^*(892)^+ \rightarrow K_S^0 \pi^+)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{190}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis. This is a doubly Cabibbo-suppressed mode.

VALUE (units $10^{-3}$ )	DOCUMENT ID	TECN	COMMENT
<b>3.4<math>\pm 1.3^{+4.1}_{-0.4}</math></b>	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$			
3.4 $\pm 1.3^{+3.6}_{-0.5}$	ASNER 04A	CLEO	See MURAMATSU 02

 $\Gamma(K_0^*(1430)^- \pi^+, K_0^*(1430)^- \rightarrow K_S^0 \pi^-)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{39}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.096<sup>+0.021</sup><sub>-0.012</sub> OUR AVERAGE</b>			
0.073 $\pm 0.007^{+0.031}_{-0.011}$	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
0.109 $\pm 0.027^{+0.029}_{-0.029}$	FRABETTI 94G	E687	$\gamma$ Be, $\bar{E}_\gamma \approx 220$ GeV
0.129 $\pm 0.034^{+0.021}_{-0.021}$	ALBRECHT 93D	ARG	$e^+ e^- \approx 10$ GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$			
0.072 $\pm 0.007^{+0.014}_{-0.013}$	ASNER 04A	CLEO	See MURAMATSU 02

$\Gamma(K_2^*(1430)^-\pi^+, K_2^*(1430)^-\rightarrow K_S^0\pi^-)/\Gamma(K_S^0\pi^+\pi^-)$   $\Gamma_{40}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.011±0.002<sup>+0.007</sup><sub>-0.003</sub></b>	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>			
0.011±0.002 <sup>+0.005</sup> <sub>-0.003</sub>	ASNER	04A CLEO	See MURAMATSU 02

 $\Gamma(K^*(1680)^-\pi^+, K^*(1680)^-\rightarrow K_S^0\pi^-)/\Gamma(K_S^0\pi^+\pi^-)$   $\Gamma_{41}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.022±0.004<sup>+0.018</sup><sub>-0.015</sub></b>	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>			
0.023±0.005 <sup>+0.007</sup> <sub>-0.014</sub>	ASNER	04A CLEO	See MURAMATSU 02

 $\Gamma(K_S^0\pi^+\pi^- \text{ nonresonant})/\Gamma(K_S^0\pi^+\pi^-)$   $\Gamma_{42}/\Gamma_{31}$ 

This is the "fit fraction" from the Dalitz-plot analysis. Neither FRABETTI 94G nor ALBRECHT 93D (quoted in many of the earlier submodes of  $K_S^0\pi^+\pi^-$ ) sees evidence for a nonresonant component.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.009±0.004<sup>+0.020</sup><sub>-0.004</sub></b>	MURAMATSU 02	CLE2	Dalitz fit, 5299 evts
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>			
0.007±0.007 <sup>+0.021</sup> <sub>-0.006</sub>	ASNER	04A CLEO	See MURAMATSU 02
0.263±0.024±0.041	ANJOS	93 E691	$\gamma$ Be 90–260 GeV
0.26 ± 0.08 ± 0.05	FRABETTI	92B E687	$\gamma$ Be, $\bar{E}_\gamma = 221$ GeV
0.33 ± 0.05 ± 0.10	ADLER	87 MRK3	$e^+e^-$ 3.77 GeV

 $\Gamma(K^-\pi^+\pi^0)/\Gamma_{\text{total}}$   $\Gamma_{43}/\Gamma$ 

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.141±0.005 OUR FIT</b>				Error includes scale factor of 1.2.
<b>0.149±0.003±0.005</b>	19k ± 150	59 HE	05 CLEO	$e^+e^-$ at $\psi(3770)$
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>				
0.133±0.012±0.013	931	ADLER	88C MRK3	$e^+e^-$ 3.77 GeV
0.117±0.043	37	60 SCHINDLER	81 MRK2	$e^+e^-$ 3.771 GeV

<sup>59</sup> HE 05 uses single- and double-tagged events in an overall fit. The fraction here includes (unobserved) final-state photons.

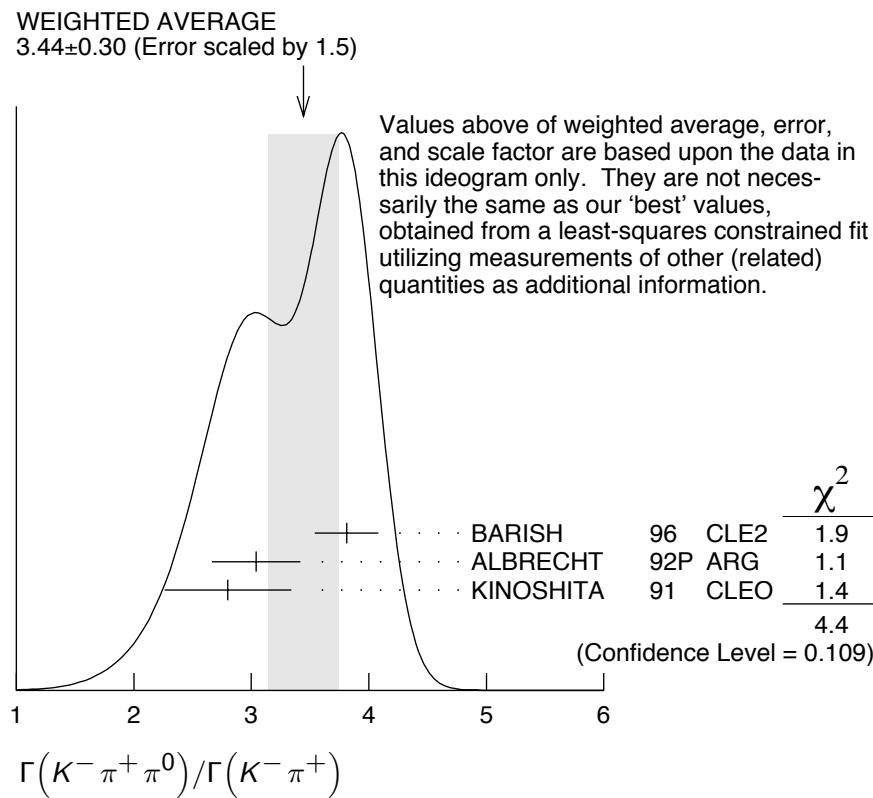
<sup>60</sup> SCHINDLER 81 (MARK-2) measures  $\sigma(e^+e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.68 \pm 0.23$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

## $\Gamma(K^-\pi^+\pi^0)/\Gamma(K^-\pi^+)$

$\Gamma_{43}/\Gamma_{29}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>3.71 \pm 0.14</math> OUR FIT</b>	Error includes scale factor of 1.4.			
<b><math>3.44 \pm 0.30</math> OUR AVERAGE</b>	Error includes scale factor of 1.5. See the ideogram below.			
$3.81 \pm 0.07 \pm 0.26$	10k	BARISH	96 CLE2	$e^+e^- \approx \gamma(4S)$
$3.04 \pm 0.16 \pm 0.34$	931	ALBRECHT	92P ARG	$e^+e^- \approx 10 \text{ GeV}$
$2.8 \pm 0.14 \pm 0.52$	1050	KINOSHITA	91 CLEO	$e^+e^- \sim 10.7 \text{ GeV}$
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
$4.0 \pm 0.9 \pm 1.0$	69	ALVAREZ	91B NA14	Photoproduction
$4.2 \pm 1.4$	41	SUMMERS	84 E691	Photoproduction

<sup>61</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.



## $\Gamma(K^-\rho^+)/\Gamma(K^-\pi^+\pi^0)$

$\Gamma_{44}/\Gamma_{43}$

This is the "fit fraction" from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b><math>0.78 \pm 0.04</math> OUR AVERAGE</b>			
$0.788 \pm 0.019 \pm 0.048$	KOPP	01 CLE2	$e^+e^- \approx 10.6 \text{ GeV}$
$0.765 \pm 0.041 \pm 0.054$	FRABETTI	94G E687	$\gamma\text{Be}, \bar{E}_\gamma \approx 220 \text{ GeV}$
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$			
$0.647 \pm 0.039 \pm 0.150$	ANJOS	93 E691	$\gamma\text{Be} 90\text{--}260 \text{ GeV}$
$0.81 \pm 0.03 \pm 0.06$	ADLER	87 MRK3	$e^+e^- 3.77 \text{ GeV}$

## $\Gamma(K^-\rho(1700)^+, \rho(1700)^+ \rightarrow \pi^+\pi^0)/\Gamma(K^-\pi^+\pi^0)$

$\Gamma_{45}/\Gamma_{43}$

This is the "fit fraction" from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b><math>0.057 \pm 0.008 \pm 0.009</math></b>	KOPP	01 CLE2	$e^+e^- \approx 10.6 \text{ GeV}$

$\Gamma(K^*(892)^-\pi^+, K^*(892)^-\rightarrow K^-\pi^0)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{46}/\Gamma_{43}$

This is the “fit fraction” from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.160<sup>+0.025</sup><sub>-0.013</sub> OUR AVERAGE</b>			
0.161 $\pm$ 0.007 <sup>+0.027</sup> <sub>-0.011</sub>	KOPP 01	CLE2	$e^+e^- \approx 10.6$ GeV
0.148 $\pm$ 0.028 $\pm$ 0.049	FRABETTI 94G E687	$\gamma$ Be, $\bar{E}_\gamma$	$\approx 220$ GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •			
0.084 $\pm$ 0.011 $\pm$ 0.012	ANJOS 93	E691	$\gamma$ Be 90–260 GeV
0.12 $\pm$ 0.02 $\pm$ 0.03	ADLER 87	MRK3	$e^+e^-$ 3.77 GeV

$\Gamma(\bar{K}^*(892)^0\pi^0, \bar{K}^*(892)^0\rightarrow K^-\pi^+)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{47}/\Gamma_{43}$

This is the “fit fraction” from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.135<math>\pm</math>0.016 OUR AVERAGE</b>			
0.127 $\pm$ 0.009 $\pm$ 0.016	KOPP 01	CLE2	$e^+e^- \approx 10.6$ GeV
0.165 $\pm$ 0.031 $\pm$ 0.015	FRABETTI 94G E687	$\gamma$ Be, $\bar{E}_\gamma$	$\approx 220$ GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •			
0.142 $\pm$ 0.018 $\pm$ 0.024	ANJOS 93	E691	$\gamma$ Be 90–260 GeV
0.13 $\pm$ 0.02 $\pm$ 0.03	ADLER 87	MRK3	$e^+e^-$ 3.77 GeV

$\Gamma(K_0^*(1430)^-\pi^+, K_0^*(1430)^-\rightarrow K^-\pi^0)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{48}/\Gamma_{43}$

This is the “fit fraction” from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.033<math>\pm</math>0.006<math>\pm</math>0.014</b>			

$\Gamma(\bar{K}_0^*(1430)^0\pi^0, \bar{K}_0^*(1430)^0\rightarrow K^-\pi^+)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{49}/\Gamma_{43}$

This is the “fit fraction” from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.041<math>\pm</math>0.006<sup>+0.032</sup><sub>-0.009</sub></b>			

$\Gamma(K^*(1680)^-\pi^+, K^*(1680)^-\rightarrow K^-\pi^0)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{50}/\Gamma_{43}$

This is the “fit fraction” from the Dalitz-plot analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.013<math>\pm</math>0.003<math>\pm</math>0.004</b>			

$\Gamma(K^-\pi^+\pi^0 \text{ nonresonant})/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{51}/\Gamma_{43}$

This is the “fit fraction” from the Dalitz-plot analysis.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.080<sup>+0.038</sup><sub>-0.014</sub> OUR AVERAGE</b>				
0.075 $\pm$ 0.009 <sup>+0.056</sup> <sub>-0.011</sub>	KOPP 01	CLE2	$e^+e^- \approx 10.6$ GeV	
0.101 $\pm$ 0.033 $\pm$ 0.040	FRABETTI 94G E687	$\gamma$ Be, $\bar{E}_\gamma$	$\approx 220$ GeV	
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.036 $\pm$ 0.004 $\pm$ 0.018	ANJOS 93	E691	$\gamma$ Be 90–260 GeV	
0.09 $\pm$ 0.02 $\pm$ 0.04	ADLER 87	MRK3	$e^+e^-$ 3.77 GeV	
0.51 $\pm$ 0.22	21	SUMMERS 84	E691 Photoproduction	

$\Gamma(\bar{K}^*(892)^0\pi^0, \bar{K}^*(892)^0 \rightarrow K_S^0\pi^0)/\Gamma(K_S^0\pi^0)$				$\Gamma_{53}/\Gamma_{30}$
<i>VALUE</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>	
$0.55^{+0.13}_{-0.10} \pm 0.07$	PROCARIO	93B CLE2	Dalitz plot fit, 122 evts	

$\Gamma(K_S^0\pi^0\pi^0 \text{ nonresonant})/\Gamma(K_S^0\pi^0)$				$\Gamma_{54}/\Gamma_{30}$
<i>VALUE</i>	<i>EVTS</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
$0.37 \pm 0.08 \pm 0.04$	76	PROCARIO	93B CLE2	Dalitz plot fit, 122 evts

$\Gamma(K^-\pi^+\pi^+\pi^-)/\Gamma_{\text{total}}$				$\Gamma_{55}/\Gamma$
<i>VALUE (units <math>10^{-2}</math>)</i>	<i>EVTS</i>	<i>DOCUMENT ID</i>	<i>TECN</i>	<i>COMMENT</i>
$7.72 \pm 0.28$ OUR FIT	Error includes scale factor of 1.3.			

**8.0 ± 0.4 OUR AVERAGE** Error includes scale factor of 1.3. See the ideogram below.

8.3 ± 0.2 ± 0.3	15k ± 130	<sup>62</sup> HE	05 CLEO	$e^+e^-$ at $\psi(3770)$
7.9 ± 1.5 ± 0.9		<sup>63</sup> ALBRECHT	94 ARG	$e^+e^- \approx \gamma(4S)$
$6.80 \pm 0.27 \pm 0.57$	$1430 \pm 52$	<sup>64</sup> ALBRECHT	94F ARG	$e^+e^- \approx \gamma(4S)$
9.1 ± 0.8 ± 0.8	992	ADLER	88C MRK3	$e^+e^-$ 3.77 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

11.7 ± 2.5	185	<sup>65</sup> SCHINDLER	81 MRK2	$e^+e^-$ 3.771 GeV
6.2 ± 1.9	44	<sup>66</sup> PERUZZI	77 MRK1	$e^+e^-$ 3.77 GeV

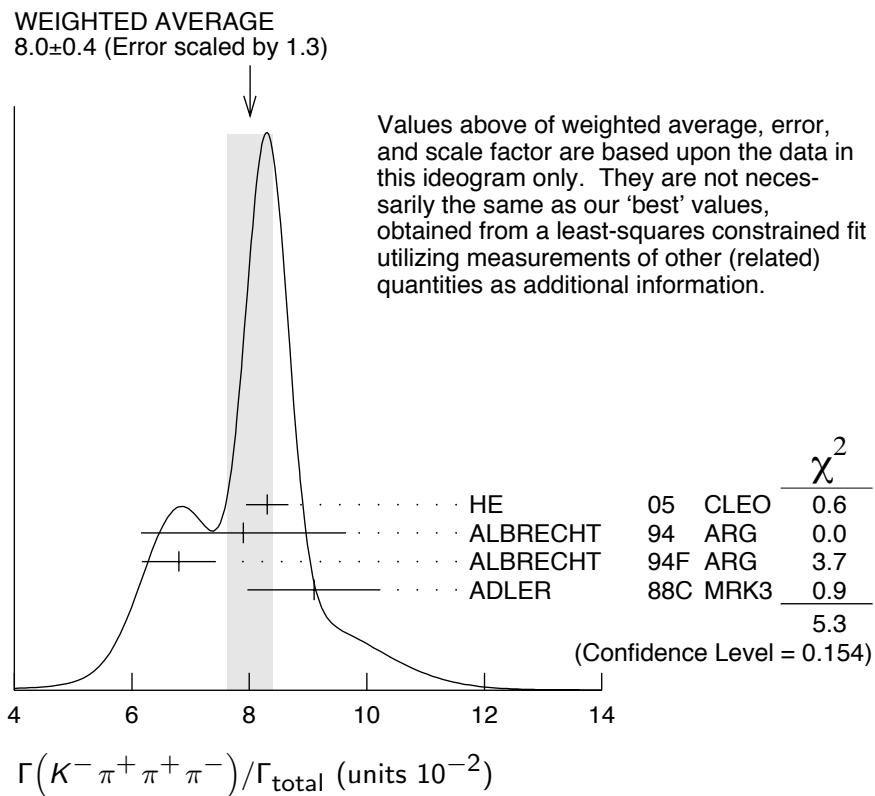
<sup>62</sup> HE 05 uses single- and double-tagged events in an overall fit. The fraction here includes (unobserved) final-state photons.

<sup>63</sup> ALBRECHT 94 uses  $D^0$  mesons from  $\bar{B}^0 \rightarrow D^*+\ell^-\bar{\nu}_\ell$  decays. This is a different set of events than used by ALBRECHT 94F.

<sup>64</sup> See the footnote on the ALBRECHT 94F measurement of  $\Gamma(K^-\pi^+)/\Gamma_{\text{total}}$  for the method used.

<sup>65</sup> SCHINDLER 81 (MARK-2) measures  $\sigma(e^+e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.68 \pm 0.11$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.

<sup>66</sup> PERUZZI 77 (MARK-1) measures  $\sigma(e^+e^- \rightarrow \psi(3770)) \times$  branching fraction to be  $0.36 \pm 0.10$  nb. We use the MARK-3 (ADLER 88C) value of  $\sigma = 5.8 \pm 0.5 \pm 0.6$  nb.



### $\Gamma(K^- \pi^+ \pi^+ \pi^-)/\Gamma(K^- \pi^+)$

### $\Gamma_{55}/\Gamma_{29}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>2.03±0.07 OUR FIT</b>				Error includes scale factor of 1.5.
<b>1.97±0.09 OUR AVERAGE</b>				
1.94±0.07 <sup>+0.09</sup> <sub>-0.11</sub>		JUN	00 SELX	$\Sigma^-$ nucleus, 600 GeV
1.7 ± 0.2 ± 0.2	1745	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
1.90±0.25±0.20	337	ALVAREZ	91B NA14	Photoproduction
2.12±0.16±0.09		BORTOLETTI088	CLEO	$e^+ e^-$ 10.55 GeV
2.17±0.28±0.23		ALBRECHT	85F ARG	$e^+ e^-$ 10 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
2.0 ± 0.9	48	BAILEY	86 ACCM	$\pi^-$ Be fixed target
2.0 ± 1.0	10	BAILEY	83B SPEC	$\pi^-$ Be → $D^0$
2.2 ± 0.8	214	PICCOLO	77 MRK1	$e^+ e^-$ 4.03, 4.41 GeV

### $\Gamma(K^- \pi^+ \rho^0 \text{total})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$

### $\Gamma_{56}/\Gamma_{55}$

This includes  $K^- a_1(1260)^+$ ,  $\bar{K}^*(892)^0 \rho^0$ , etc. The next entry gives the specifically 3-body fraction. We rely on the MARK III and E691 full amplitude analyses of the  $K^- \pi^+ \pi^+ \pi^-$  channel for values of the resonant substructure.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.835±0.035 OUR AVERAGE</b>			
0.80 ± 0.03 ± 0.05	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
0.855±0.032±0.030	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •			
0.98 ± 0.12 ± 0.10	ALVAREZ	91B NA14	Photoproduction

$\Gamma(K^-\pi^+\rho^0\text{3-body})/\Gamma(K^-\pi^+\pi^+\pi^-)$  $\Gamma_{57}/\Gamma_{55}$ 

We rely on the MARK III and E691 full amplitude analyses of the  $K^-\pi^+\pi^+\pi^-$  channel for values of the resonant substructure.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.063±0.028 OUR AVERAGE</b>				
0.05 ± 0.03 ± 0.02		ANJOS	92C E691	$\gamma$ Be 90–260 GeV
0.084±0.022±0.04		COFFMAN	92B MRK3	$e^+e^-$ 3.77 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.77 ± 0.06 ± 0.06	67	ALVAREZ	91B NA14	Photoproduction
0.85 $^{+0.11}_{-0.22}$	180	PICCOLO	77 MRK1	$e^+e^-$ 4.03, 4.41 GeV

<sup>67</sup> This value is for  $\rho^0$  ( $K^-\pi^+$ )-nonresonant. ALVAREZ 91B cannot determine what fraction of this is  $K^-\pi_1(1260)^+$ .

 $\Gamma(\bar{K}^*(892)^0\rho^0)/\Gamma(K^-\pi^+\pi^+\pi^-)$  $\Gamma_{94}/\Gamma_{55}$ 

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included. We rely on the MARK III and E691 full amplitude analyses of the  $K^-\pi^+\pi^+\pi^-$  channel for values of the resonant substructure.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.195±0.03±0.03</b>				
		ANJOS	92C E691	$\gamma$ Be 90–260 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.34 ± 0.09 ± 0.09		ALVAREZ	91B NA14	Photoproduction
0.75 ± 0.3	5	BAILEY	83B SPEC	$\pi$ Be → $D^0$
0.15 $^{+0.16}_{-0.15}$	20	PICCOLO	77 MRK1	$e^+e^-$ 4.03, 4.41 GeV

 $\Gamma(\bar{K}^*(892)^0\rho^0\text{transverse})/\Gamma(K^-\pi^+\pi^+\pi^-)$  $\Gamma_{95}/\Gamma_{55}$ 

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.20 ± 0.07 OUR FIT</b>			
<b>0.213±0.024±0.075</b>	COFFMAN	92B MRK3	$e^+e^-$ 3.77 GeV

 $\Gamma(\bar{K}^*(892)^0\rho^0S\text{-wave})/\Gamma(K^-\pi^+\pi^+\pi^-)$  $\Gamma_{96}/\Gamma_{55}$ 

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.375±0.045±0.06</b>			

 $\Gamma(\bar{K}^*(892)^0\rho^0S\text{-wave long.})/\Gamma_{\text{total}}$  $\Gamma_{97}/\Gamma$ 

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<0.003	90	COFFMAN	92B MRK3	$e^+e^-$ 3.77 GeV

 $\Gamma(\bar{K}^*(892)^0\rho^0P\text{-wave})/\Gamma_{\text{total}}$  $\Gamma_{98}/\Gamma$ 

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<0.003	90	COFFMAN	92B MRK3	$e^+e^-$ 3.77 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.009	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
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$\Gamma(\bar{K}^*(892)^0 \rho^0 D\text{-wave})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{99}/\Gamma_{55}$ Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b><math>0.255 \pm 0.045 \pm 0.06</math></b>	ANJOS	92C E691	$\gamma$ Be 90–260 GeV

 $\Gamma(K^- \pi^+ f_0(980))/\Gamma_{\text{total}}$   $\Gamma_{104}/\Gamma$ 

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
<0.011	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV

 $\Gamma(\bar{K}^*(892)^0 f_0(980))/\Gamma_{\text{total}}$   $\Gamma_{105}/\Gamma$ 

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
<0.007	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV

 $\Gamma(K^- a_1(1260)^+)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{89}/\Gamma_{55}$ Unseen decay modes of the  $a_1(1260)^+$  are included, assuming that the  $a_1(1260)^+$  decays entirely to  $\rho\pi$  [or at least to  $(\pi\pi)_{I=1}\pi$ ].

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.97 <math>\pm 0.14</math> OUR AVERAGE</b>			
0.94 $\pm 0.13 \pm 0.20$	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
$0.984 \pm 0.048 \pm 0.16$	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K^- a_2(1320)^+)/\Gamma_{\text{total}}$   $\Gamma_{91}/\Gamma$ Unseen decay modes of the  $a_2(1320)^+$  are included.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.002</b>	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
<0.006	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K_1(1270)^- \pi^+)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{106}/\Gamma_{55}$ Unseen decay modes of the  $K_1(1270)^-$  are included. The MARK3 and E691 experiments disagree considerably here.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.15 <math>\pm 0.04</math> OUR FIT</b>				
<b><math>0.194 \pm 0.056 \pm 0.088</math></b>		COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
<0.013	90	ANJOS	92C E691	$\gamma$ Be 90–260 GeV

 $\Gamma(K_1(1400)^- \pi^+)/\Gamma_{\text{total}}$   $\Gamma_{107}/\Gamma$ 

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.012</b>	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K^*(1410)^- \pi^+)/\Gamma_{\text{total}}$   $\Gamma_{109}/\Gamma$ 

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$				
<0.012	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(\bar{K}^*(892)^0 \pi^+ \pi^- \text{total})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{92}/\Gamma_{55}$

This includes  $\bar{K}^*(892)^0 \rho^0$ , etc. The next entry gives the specifically 3-body fraction.  
Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.30 ± 0.06 ± 0.03</b>	ANJOS	92C E691	$\gamma$ Be 90–260 GeV

$\Gamma(\bar{K}^*(892)^0 \pi^+ \pi^- \text{3-body})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{93}/\Gamma_{55}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.19 ± 0.04 OUR FIT</b>			
<b>0.18 ± 0.04 OUR AVERAGE</b>			

0.165 ± 0.03 ± 0.045 ANJOS 92C E691  $\gamma$ Be 90–260 GeV  
0.210 ± 0.027 ± 0.06 COFFMAN 92B MRK3  $e^+ e^-$  3.77 GeV

$\Gamma(K^- \pi^+ \pi^+ \pi^- \text{nonresonant})/\Gamma(K^- \pi^+ \pi^+ \pi^-)$   $\Gamma_{63}/\Gamma_{55}$

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.233 ± 0.032 OUR AVERAGE</b>			
0.23 ± 0.02 ± 0.03	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
0.242 ± 0.025 ± 0.06	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

$\Gamma(K_S^0 \pi^+ \pi^- \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{64}/\Gamma$

VALUE (units $10^{-2}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>5.3 ± 0.6 OUR FIT</b>				

**5.2 ± 1.1 ± 1.2** 140 COFFMAN 92B MRK3  $e^+ e^-$  3.77 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$6.7^{+1.6}_{-1.7}$  68 BARLAG 92C ACCM  $\pi^-$  Cu 230 GeV

68 BARLAG 92C computes the branching fraction using topological normalization.

$\Gamma(K_S^0 \pi^+ \pi^- \pi^0)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{64}/\Gamma_{31}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>1.83 ± 0.20 OUR FIT</b>				
<b>1.86 ± 0.23 OUR AVERAGE</b>				
1.80 ± 0.20 ± 0.21	190	69 ALBRECHT	92P ARG	$e^+ e^- \approx 10$ GeV
2.8 ± 0.8 ± 0.8	46	ANJOS	92C E691	$\gamma$ Be 90–260 GeV
1.85 ± 0.26 ± 0.30	158	KINOSHITA	91 CLEO	$e^+ e^- \sim 10.7$ GeV

69 This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(K_S^0 \eta)/\Gamma(K_S^0 \pi^0)$   $\Gamma_{86}/\Gamma_{30}$

Unseen decay modes of the  $\eta$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.33 ± 0.04 OUR FIT</b>				
<b>0.32 ± 0.04 ± 0.03</b>	225 ± 30	PROCARIO	93B CLE2	$\eta \rightarrow \gamma\gamma$

$\Gamma(K_S^0 \eta)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{86}/\Gamma_{31}$

Unseen decay modes of the  $\eta$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.131 ± 0.018 OUR FIT</b>				
<b>0.14 ± 0.02 ± 0.02</b>	80 ± 12	PROCARIO	93B CLE2	$\eta \rightarrow \pi^+ \pi^- \pi^0$

$\Gamma(K_S^0 \omega)/\Gamma(K^- \pi^+)$  $\Gamma_{87}/\Gamma_{29}$ Unseen decay modes of the  $\omega$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.29±0.05 OUR FIT</b>			
<b>0.50±0.18±0.10</b>	ALBRECHT	89D ARG	$e^+ e^-$ 10 GeV

 $\Gamma(K_S^0 \omega)/\Gamma(K_S^0 \pi^+ \pi^-)$  $\Gamma_{87}/\Gamma_{31}$ Unseen decay modes of the  $\omega$  are included.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.38±0.07 OUR FIT</b>				
<b>0.33±0.09 OUR AVERAGE</b>				Error includes scale factor of 1.1.

0.29±0.08±0.05	16	70 ALBRECHT	92P ARG	$e^+ e^- \approx$ 10 GeV
0.54±0.14±0.16	40	KINOSHITA	91 CLEO	$e^+ e^- \sim$ 10.7 GeV

70 This value is calculated from numbers in Table 1 of ALBRECHT 92P.

 $\Gamma(K_S^0 \omega)/\Gamma(K_S^0 \pi^+ \pi^- \pi^0)$  $\Gamma_{87}/\Gamma_{64}$ Unseen decay modes of the  $\omega$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.21 ±0.04 OUR FIT</b>			
<b>0.220±0.048±0.0116</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K_S^0 \eta'(958))/\Gamma(K_S^0 \pi^+ \pi^-)$  $\Gamma_{88}/\Gamma_{31}$ Unseen decay modes of the  $\eta'(958)$  are included.

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.32±0.04 OUR AVERAGE</b>				
0.31±0.02±0.04	594	PROCARIO	93B CLE2	$\eta' \rightarrow \eta \pi^+ \pi^-$ , $\rho^0 \gamma$
0.37±0.13±0.06	18	71 ALBRECHT	92P ARG	$e^+ e^- \approx$ 10 GeV

71 This value is calculated from numbers in Table 1 of ALBRECHT 92P.

 $\Gamma(K^*(892)^- \rho^+)/\Gamma(K_S^0 \pi^+ \pi^- \pi^0)$  $\Gamma_{100}/\Gamma_{64}$ Unseen decay modes of the  $K^*(892)^-$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>1.212±0.376±0.252</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K^*(892)^- \rho^+ \text{ longitudinal})/\Gamma(K_S^0 \pi^+ \pi^- \pi^0)$  $\Gamma_{101}/\Gamma_{64}$ Unseen decay modes of the  $K^*(892)^-$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.580±0.222</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K^*(892)^- \rho^+ \text{ transverse})/\Gamma(K_S^0 \pi^+ \pi^- \pi^0)$  $\Gamma_{102}/\Gamma_{64}$ Unseen decay modes of the  $K^*(892)^-$  are included.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.634±0.360</b>	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K^*(892)^- \rho^+ P\text{-wave})/\Gamma_{\text{total}}$  $\Gamma_{103}/\Gamma$ Unseen decay modes of the  $K^*(892)^-$  are included.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.015</b>	90	72 COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

72 Obtained using other  $\bar{K}^*(892)^- \rho$  P-wave limits and isospin relations.

$\Gamma(\bar{K}^*(892)^0 \rho^0 \text{transverse})/\Gamma(K_S^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{95}/\Gamma_{64}$ 
Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>		<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.30 ± 0.11 OUR FIT</b>				
<b>0.252 ± 0.222</b>		COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(\bar{K}^0 a_1(1260)^0)/\Gamma_{\text{total}}$   $\Gamma_{90}/\Gamma$ 
Unseen decay modes of the  $a_1(1260)^+$  are included, assuming that the  $a_1(1260)^+$  decays entirely to  $\rho\pi$  [or at least to  $(\pi\pi)_{I=1}\pi$ ].

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<0.019	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K_1(1270)^- \pi^+)/\Gamma(K_S^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{106}/\Gamma_{64}$ 
Unseen decay modes of the  $K_1(1270)^-$  are included.

<u>VALUE</u>		<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.21 ± 0.06 OUR FIT</b>				
<b>0.20 ± 0.06</b>		COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(\bar{K}_1(1400)^0 \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{108}/\Gamma$ 

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<0.037	90	COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(\bar{K}^*(892)^0 \pi^+ \pi^- \text{3-body})/\Gamma(K_S^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{93}/\Gamma_{64}$ 
Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

<u>VALUE</u>		<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.28 ± 0.07 OUR FIT</b>	Error includes scale factor of 1.1.			
<b>0.382 ± 0.210</b>		COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K_S^0 \pi^+ \pi^- \pi^0 \text{nonresonant})/\Gamma(K_S^0 \pi^+ \pi^- \pi^0)$   $\Gamma_{70}/\Gamma_{64}$ 

<u>VALUE</u>		<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.210 ± 0.147 ± 0.150</b>		COFFMAN	92B MRK3	$e^+ e^-$ 3.77 GeV

 $\Gamma(K^- \pi^+ \pi^0 \pi^0)/\Gamma_{\text{total}}$   $\Gamma_{71}/\Gamma$ 

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
0.177 ± 0.029		73 BARLAG	92C ACCM	$\pi^-$ Cu 230 GeV
0.149 ± 0.037 ± 0.030	24	74 ADLER	88C MRK3	$e^+ e^-$ 3.77 GeV
0.209 ± 0.074 ± 0.012	9	73 AGUILAR-...	87F HYBR	$\pi p, pp$ 360, 400 GeV

<sup>73</sup> AGUILAR-BENITEZ 87F and BARLAG 92C compute the branching fraction using topological normalization. They do not distinguish the presence of a third  $\pi^0$ , and thus are not included in the average.

<sup>74</sup> ADLER 88C uses an absolute normalization method finding this decay channel opposite a detected  $\bar{D}^0 \rightarrow K^+ \pi^-$  in pure  $D\bar{D}$  events.

 $\Gamma(K^- \pi^+ \pi^+ \pi^- \pi^0)/\Gamma(K^- \pi^+)$   $\Gamma_{72}/\Gamma_{29}$ 

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>1.08 ± 0.10 OUR FIT</b>				
<b>0.98 ± 0.11 ± 0.11</b>	225	75 ALBRECHT	92P ARG	$e^+ e^-$ ≈ 10 GeV

<sup>75</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(K^-\pi^+\pi^+\pi^-\pi^0)/\Gamma(K^-\pi^+\pi^+\pi^-)$   $\Gamma_{72}/\Gamma_{55}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.53±0.05 OUR FIT</b>				
<b>0.56±0.07 OUR AVERAGE</b>				
$0.55 \pm 0.07^{+0.12}_{-0.09}$	167	KINOSHITA 91	CLEO	$e^+e^- \sim 10.7 \text{ GeV}$
$0.57 \pm 0.06 \pm 0.05$	180	ANJOS 90D	E691	Photoproduction

$\Gamma(\bar{K}^*(892)^0\pi^+\pi^-\pi^0)/\Gamma(K^-\pi^+\pi^+\pi^-\pi^0)$   $\Gamma_{110}/\Gamma_{72}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.45±0.15±0.15</b>		ANJOS 90D	E691	Photoproduction

$\Gamma(\bar{K}^*(892)^0\eta)/\Gamma(K^-\pi^+)$   $\Gamma_{111}/\Gamma_{29}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  and  $\eta$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				

$0.58 \pm 0.19^{+0.24}_{-0.28}$	46	KINOSHITA 91	CLEO	$e^+e^- \sim 10.7 \text{ GeV}$
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$\Gamma(\bar{K}^*(892)^0\eta)/\Gamma(K^-\pi^+\pi^0)$   $\Gamma_{111}/\Gamma_{43}$

Unseen decay modes of the  $\bar{K}^*(892)^0$  and  $\eta$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
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• • • We do not use the following data for averages, fits, limits, etc. • • •

$0.13 \pm 0.02 \pm 0.03$	214	PROCARIO 93B	CLE2	$\bar{K}^*{}^0\eta \rightarrow K^-\pi^+/\gamma\gamma$
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$\Gamma(K_S^0\eta\pi^0)/\Gamma(K_S^0\pi^0)$   $\Gamma_{76}/\Gamma_{30}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.46±0.07±0.06</b>	$155 \pm 22$	76 RUBIN	04 CLEO	$e^+e^- \approx 10 \text{ GeV}$

<sup>76</sup> The  $\eta$  here is detected in its  $\gamma\gamma$  mode, but other  $\eta$  modes are included in the value given.

$\Gamma(K_S^0a_0(980), a_0(980) \rightarrow \eta\pi^0)/\Gamma(K_S^0\eta\pi^0)$   $\Gamma_{77}/\Gamma_{76}$

This is the “fit fraction” from the Dalitz-plot analysis, with interference.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>1.19±0.09±0.26</b>	77 RUBIN	04 CLEO	Dalitz fit, 155 evts	

<sup>77</sup> In addition to  $K_S^0a_0(980)$  and  $\bar{K}^*(892)^0\eta$  modes, RUBIN 04 finds a fit fraction of  $0.246 \pm 0.092 \pm 0.091$  for other, undetermined modes.

$\Gamma(\bar{K}^*(892)^0\eta, \bar{K}^*(892)^0 \rightarrow K_S^0\pi^0)/\Gamma(K_S^0\eta\pi^0)$   $\Gamma_{78}/\Gamma_{76}$

This is the “fit fraction” from the Dalitz-plot analysis, with interference.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.293±0.062±0.035</b>	78 RUBIN	04 CLEO	Dalitz fit, 155 evts	

<sup>78</sup> See the note on RUBIN 04 in the preceding data block.

$\Gamma(K^-\pi^+\omega)/\Gamma(K^-\pi^+)$   $\Gamma_{112}/\Gamma_{29}$

Unseen decay modes of the  $\omega$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.78±0.12±0.10</b>	99	79 ALBRECHT	92P ARG	$e^+e^- \approx 10 \text{ GeV}$

<sup>79</sup> This value is calculated from numbers in Table 1 of ALBRECHT 92P.

$\Gamma(\bar{K}^*(892)^0 \omega)/\Gamma(K^- \pi^+)$  $\Gamma_{113}/\Gamma_{29}$ Unseen decay modes of the  $\bar{K}^*(892)^0$  and  $\omega$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.28±0.11±0.04</b>	17	80 ALBRECHT	92P ARG	$e^+ e^- \approx 10 \text{ GeV}$

80 This value is calculated from numbers in Table 1 of ALBRECHT 92P.

 $\Gamma(K^- \pi^+ \eta'(958))/\Gamma(K^- \pi^+ \pi^+ \pi^-)$  $\Gamma_{114}/\Gamma_{55}$ Unseen decay modes of the  $\eta'(958)$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.093±0.014±0.019</b>	286	PROCARIO	93B CLE2	$\eta' \rightarrow \eta \pi^+ \pi^-, \rho^0 \gamma$

 $\Gamma(\bar{K}^*(892)^0 \eta'(958))/\Gamma(K^- \pi^+ \eta'(958))$  $\Gamma_{115}/\Gamma_{114}$ Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	CL%	DOCUMENT ID	TECN
<b>&lt;0.15</b>	90	PROCARIO	93B CLE2

 $\Gamma(K_S^0 2\pi^+ 2\pi^-)/\Gamma(K_S^0 \pi^+ \pi^-)$  $\Gamma_{79}/\Gamma_{31}$ 

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.095±0.005±0.007</b>	1283 ± 57	LINK	04D FOCS	$\gamma A, \bar{E}_\gamma \approx 180 \text{ GeV}$

• • • We do not use the following data for averages, fits, limits, etc. • • •

0.07 ± 0.02 ± 0.01	11	81 ALBRECHT	92P ARG	$e^+ e^- \approx 10 \text{ GeV}$
0.149 ± 0.026	56	AMMAR	91 CLEO	$e^+ e^- \approx 10.5 \text{ GeV}$
0.18 ± 0.07 ± 0.04	6	ANJOS	90D E691	Photoproduction

81 This value is calculated from numbers in Table 1 of ALBRECHT 92P.

 $\Gamma(K_S^0 \rho^0 \pi^+ \pi^-, \text{no } K^*(892)^-)/\Gamma(K_S^0 2\pi^+ 2\pi^-)$  $\Gamma_{80}/\Gamma_{79}$ 

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.40±0.24±0.07</b>	LINK	04D FOCS	$\gamma A, \bar{E}_\gamma \approx 180 \text{ GeV}$

 $\Gamma(K^*(892)^- \pi^+ \pi^+ \pi^-, K^*(892)^- \rightarrow K_S^0 \pi^-, \text{no } \rho^0)/\Gamma(K_S^0 2\pi^+ 2\pi^-)$  $\Gamma_{81}/\Gamma_{79}$ 

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.17±0.28±0.02</b>	LINK	04D FOCS	$\gamma A, \bar{E}_\gamma \approx 180 \text{ GeV}$

 $\Gamma(K^*(892)^- \rho^0 \pi^+, K^*(892)^- \rightarrow K_S^0 \pi^-)/\Gamma(K_S^0 2\pi^+ 2\pi^-)$  $\Gamma_{82}/\Gamma_{79}$ 

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.60±0.21±0.09</b>	LINK	04D FOCS	$\gamma A, \bar{E}_\gamma \approx 180 \text{ GeV}$

 $\Gamma(K_S^0 2\pi^+ 2\pi^- \text{ nonresonant})/\Gamma(K_S^0 2\pi^+ 2\pi^-)$  $\Gamma_{83}/\Gamma_{79}$ 

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.46</b>	90	LINK	04D FOCS	$\gamma A, \bar{E}_\gamma \approx 180 \text{ GeV}$

 $\Gamma(K^- 3\pi^+ 2\pi^-)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$  $\Gamma_{85}/\Gamma_{55}$ 

VALUE (units $10^{-3}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>2.70±0.58±0.38</b>	48 ± 10	LINK	04B FOCS	$\gamma A, \bar{E}_\gamma \approx 180 \text{ GeV}$

**Hadronic modes with three  $K$ 's**

$$\Gamma(K_S^0 K^+ K^-)/\Gamma(K_S^0 \pi^+ \pi^-)$$

$$\Gamma_{116}/\Gamma_{31}$$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.158±0.001±0.005</b>	14k±116	AUBERT,B	05J BABR	$e^+ e^- \approx \Upsilon(4S)$
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>				
0.20 ± 0.05 ± 0.04	47	FRABETTI	92B E687	$\gamma$ Be, $\bar{E}_\gamma = 221$ GeV
0.170±0.022	136	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
0.24 ± 0.08		BEBEK	86 CLEO	$e^+ e^-$ near $\Upsilon(4S)$
0.185±0.055	52	ALBRECHT	85B ARG	$e^+ e^-$ 10 GeV

$$\Gamma(K_S^0 a_0(980)^0, a_0^0 \rightarrow K^+ K^-)/\Gamma(K_S^0 K^+ K^-)$$

$$\Gamma_{117}/\Gamma_{116}$$

This is the "fit fraction" from the Dalitz-plot analysis, with interference.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.664±0.016±0.070</b>	AUBERT,B	05J BABR	Dalitz fit, 12540 ± 112 evts

$$\Gamma(K^- a_0(980)^+, a_0^+ \rightarrow K^+ K_S^0)/\Gamma(K_S^0 K^+ K^-)$$

$$\Gamma_{118}/\Gamma_{116}$$

This is the "fit fraction" from the Dalitz-plot analysis, with interference.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.134±0.011±0.037</b>	AUBERT,B	05J BABR	Dalitz fit, 12540 ± 112 evts

$$\Gamma(K^+ a_0(980)^-, a_0^- \rightarrow K^- K_S^0)/\Gamma(K_S^0 K^+ K^-)$$

$$\Gamma_{119}/\Gamma_{116}$$

This is a doubly Cabibbo-suppressed mode.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<0.025	95	AUBERT,B	05J BABR	Dalitz fit, 12540 ± 112 evts

$$\Gamma(K_S^0 f_0(980), f_0 \rightarrow K^+ K^-)/\Gamma(K_S^0 K^+ K^-)$$

$$\Gamma_{120}/\Gamma_{116}$$

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<0.021	95	AUBERT,B	05J BABR	Dalitz fit, 12540 ± 112 evts

$$\Gamma(K_S^0 \phi, \phi \rightarrow K^+ K^-)/\Gamma(K_S^0 K^+ K^-)$$

$$\Gamma_{121}/\Gamma_{116}$$

This is the "fit fraction" from the Dalitz-plot analysis, with interference.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.459±0.007±0.007</b>	AUBERT,B	05J BABR	Dalitz fit, 12540 ± 112 evts

$$\Gamma(K_S^0 f_0(1400), f_0 \rightarrow K^+ K^-)/\Gamma(K_S^0 K^+ K^-)$$

$$\Gamma_{122}/\Gamma_{116}$$

This is the "fit fraction" from the Dalitz-plot analysis, with interference.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.038±0.007±0.023</b>	AUBERT,B	05J BABR	Dalitz fit, 12540 ± 112 evts

$$\Gamma(3K_S^0)/\Gamma(K_S^0 \pi^+ \pi^-)$$

$$\Gamma_{123}/\Gamma_{31}$$

VALUE (units $10^{-2}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>3.2 ± 0.4 OUR AVERAGE</b>				
3.58±0.54±0.52	170 ± 26	LINK	05A FOCS	$\gamma$ Be, $\bar{E}_\gamma \approx 180$ GeV
2.78±0.38±0.48	61	ASNER	96B CLE2	$e^+ e^- \approx \Upsilon(4S)$
7.0 ± 2.4 ± 1.2	10 ± 3	FRABETTI	94J E687	$\gamma$ Be, $\bar{E}_\gamma = 220$ GeV
3.2 ± 1.0	22	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
3.4 ± 1.4 ± 1.0	5	ALBRECHT	90C ARG	$e^+ e^- \approx 10$ GeV

$\Gamma(K^+ K^- \bar{K}^*(892)^0)/\Gamma(K^+ \pi^+ \pi^-)$   $\Gamma_{124}/\Gamma_{55}$ 

<u>VALUE</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	
<b>0.0027 ± 0.0004 OUR AVERAGE</b>		Error includes scale factor of 1.1.			
0.00257 ± 0.00034 ± 0.00024	143	LINK	03G FOCS	$\gamma$ A, $\bar{E}_\gamma \approx 180$ GeV	
0.0054 ± 0.0016 ± 0.0008	18	AITALA	01D E791	$\pi^-$ A, 500 GeV	
0.0028 ± 0.0007 ± 0.0001	20	FRABETTI	95C E687	$\gamma$ Be, $\bar{E}_\gamma \approx 200$ GeV	

 $\Gamma(\phi \bar{K}^*(892)^0, \phi \rightarrow K^+ K^-, \bar{K}^*(892)^0 \rightarrow K^- \pi^+)/\Gamma(K^+ K^- K^- \pi^+)$   $\Gamma_{127}/\Gamma_{124}$ 

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.48 ± 0.06 ± 0.01</b>	LINK	03G FOCS	$\gamma$ A, $\bar{E}_\gamma \approx 180$ GeV

 $\Gamma(K^- \pi^+ \phi, \phi \rightarrow K^+ K^-)/\Gamma(K^+ K^- K^- \pi^+)$   $\Gamma_{126}/\Gamma_{124}$ 

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.18 ± 0.06 ± 0.04</b>	LINK	03G FOCS	$\gamma$ A, $\bar{E}_\gamma \approx 180$ GeV

 $\Gamma(K^+ K^- \bar{K}^*(892)^0, \bar{K}^*(892)^0 \rightarrow K^- \pi^+)/\Gamma(K^+ K^- K^- \pi^+)$   $\Gamma_{125}/\Gamma_{124}$ 

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.20 ± 0.07 ± 0.02</b>	LINK	03G FOCS	$\gamma$ A, $\bar{E}_\gamma \approx 180$ GeV

 $\Gamma(K^+ K^- K^- \pi^+ \text{nonresonant})/\Gamma(K^+ K^- K^- \pi^+)$   $\Gamma_{128}/\Gamma_{124}$ 

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>0.15 ± 0.06 ± 0.02</b>	LINK	03G FOCS	$\gamma$ A, $\bar{E}_\gamma \approx 180$ GeV

 $\Gamma(K_S^0 K_S^0 K^\pm \pi^\mp)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{129}/\Gamma_{31}$ 

<u>VALUE (units <math>10^{-2}</math>)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>2.12 ± 0.38 ± 0.20</b>	57 ± 10	LINK	05A FOCS	$\gamma$ Be, $\bar{E}_\gamma \approx 180$ GeV

**Pionic modes**

 $\Gamma(\pi^+ \pi^-)/\Gamma(K^- \pi^+)$   $\Gamma_{130}/\Gamma_{29}$ 

<u>VALUE (units <math>10^{-2}</math>)</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>3.59 ± 0.05 OUR AVERAGE</b>				
3.62 ± 0.10 ± 0.08	2085 ± 54	RUBIN	06 CLEO	$e^+ e^-$ at $\psi(3770)$
3.594 ± 0.054 ± 0.040	7334 ± 97	ACOSTA	05C CDF	$p\bar{p}$ , $\sqrt{s} = 1.96$ TeV
3.53 ± 0.12 ± 0.06	3453	LINK	03 FOCS	$\gamma$ A, $\bar{E}_\gamma \approx 180$ GeV
3.51 ± 0.16 ± 0.17	710	CSORNA	02 CLE2	$e^+ e^- \approx \gamma(4S)$
4.0 ± 0.2 ± 0.3	2043	AITALA	98C E791	$\pi^-$ A, 500 GeV
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
3.4 ± 0.7 ± 0.1	76 ± 15	ABLIKIM	05F BES	$e^+ e^- \approx \psi(3770)$
4.3 ± 0.7 ± 0.3	177	FRABETTI	94C E687	$\gamma$ Be, $\bar{E}_\gamma = 220$ GeV
3.48 ± 0.30 ± 0.23	227	SELEN	93 CLE2	$e^+ e^- \approx \gamma(4S)$
5.5 ± 0.8 ± 0.5	120	ANJOS	91D E691	Photoproduction
5.0 ± 0.7 ± 0.5	110	ALEXANDER	90 CLEO	$e^+ e^-$ 10.5–11 GeV

$\Gamma(\pi^0\pi^0)/\Gamma(K^-\pi^+)$  $\Gamma_{131}/\Gamma_{29}$ 

<u>VALUE</u> (units $10^{-2}$ )	<u>EVTS</u>
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**2.07±0.19 OUR AVERAGE**

$2.05 \pm 0.13 \pm 0.16$	$499 \pm 32$
$2.2 \pm 0.4 \pm 0.4$	40

<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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RUBIN	06	CLEO $e^+e^-$ at $\psi(3770)$
SELEN	93	CLE2 $e^+e^- \approx \gamma(4S)$

 $\Gamma(\pi^+\pi^-\pi^0)/\Gamma(K^-\pi^+)$  $\Gamma_{132}/\Gamma_{29}$ 

<u>VALUE</u> (units $10^{-2}$ )	<u>EVTS</u>
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**34.4±0.5±1.2** $11k \pm 164$ 

<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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RUBIN	06	CLEO $e^+e^-$ at $\psi(3770)$
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 $\Gamma(\rho^+\pi^-)/\Gamma(\pi^+\pi^-\pi^0)$  $\Gamma_{133}/\Gamma_{132}$ 

This is the "fit fraction" from the Dalitz-plot analysis, with interference.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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**0.763±0.019±0.025**

CRONIN-HEN..05	CLEO	$e^+e^- \approx 10 \text{ GeV}$
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 $\Gamma(\rho^0\pi^0)/\Gamma(\pi^+\pi^-\pi^0)$  $\Gamma_{134}/\Gamma_{132}$ 

This is the "fit fraction" from the Dalitz-plot analysis, with interference.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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**0.244±0.020±0.021**

CRONIN-HEN..05	CLEO	$e^+e^- \approx 10 \text{ GeV}$
----------------	------	---------------------------------

 $\Gamma(\rho^-\pi^+)/\Gamma(\pi^+\pi^-\pi^0)$  $\Gamma_{135}/\Gamma_{132}$ 

This is the "fit fraction" from the Dalitz-plot analysis, with interference.

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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**0.345±0.024±0.013**

CRONIN-HEN..05	CLEO	$e^+e^- \approx 10 \text{ GeV}$
----------------	------	---------------------------------

 $\Gamma(f_0(980)\pi^0, f_0(980) \rightarrow \pi^+\pi^-)/\Gamma(\pi^+\pi^-\pi^0)$  $\Gamma_{136}/\Gamma_{132}$ 

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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**<2.6 × 10<sup>-4</sup>**

95

82

CRONIN-HEN..05 CLEO     $e^+e^- \approx 10 \text{ GeV}$ 

<sup>82</sup> The CRONIN-HENNESSY 05 fit here includes, in addition to the three  $\rho\pi$  charged states, only the  $f_0(980)\pi^0$  mode. See also the next entries for limits obtained in the same way for the  $f_0(600)\pi^0$  mode and for an  $S$ -wave  $\pi^+\pi^-$  parametrized using a  $K$ -matrix. Our  $\rho\pi$  branching ratios, given above, use the fit with the  $K$ -matrix  $S$  wave.

 $\Gamma(f_0(600)\pi^0, f_0(600) \rightarrow \pi^+\pi^-)/\Gamma(\pi^+\pi^-\pi^0)$  $\Gamma_{137}/\Gamma_{132}$ The  $f_0(600)$  is the  $\sigma$ .

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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**<2.1 × 10<sup>-3</sup>**

95

83

CRONIN-HEN..05 CLEO     $e^+e^- \approx 10 \text{ GeV}$ 

<sup>83</sup> See the note on CRONIN-HENNESSY 05 in the proceeding data block.

 $\Gamma((\pi^+\pi^-)S\text{-wave}\pi^0)/\Gamma(\pi^+\pi^-\pi^0)$  $\Gamma_{138}/\Gamma_{132}$ 

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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**<0.019**

95

84

CRONIN-HEN..05 CLEO     $e^+e^- \approx 10 \text{ GeV}$ 

<sup>84</sup> See the note on CRONIN-HENNESSY 05 two data blocks up.

 $\Gamma(3\pi^0)/\Gamma_{\text{total}}$  $\Gamma_{139}/\Gamma$ 

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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**<3.5 × 10<sup>-4</sup>**

90

RUBIN	06	CLEO $e^+e^-$ at $\psi(3770)$
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$\Gamma(2\pi^+ 2\pi^-)/\Gamma(K^- \pi^+)$  $\Gamma_{140}/\Gamma_{29}$ VALUE (units  $10^{-2}$ ) EVTS **$19.2 \pm 0.6$  OUR FIT** **$19.1 \pm 0.4 \pm 0.6$**  $7331 \pm 130$ DOCUMENT IDTECNCOMMENT

RUBIN

06

CLEO

 $e^+ e^-$  at  $\psi(3770)$ 

|

 $\Gamma(2\pi^+ 2\pi^-)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$  $\Gamma_{140}/\Gamma_{55}$ VALUE EVTS **$0.095 \pm 0.004$  OUR FIT** Error includes scale factor of 1.2. **$0.096 \pm 0.005$  OUR AVERAGE** $0.079 \pm 0.018 \pm 0.005$ 

162

ABLIKIM

05F

BES

 $e^+ e^- \approx \psi(3770)$  $0.095 \pm 0.007 \pm 0.002$ 

814

FRABETTI

95C

E687

 $\gamma$ Be, $\bar{E}_\gamma \approx 200$  GeV $0.115 \pm 0.023 \pm 0.016$ 

64

ADAMOVICH

92

OMEG

 $\pi^-$  340 GeV $0.108 \pm 0.024 \pm 0.008$ 

79

FRABETTI

92

E687

 $\gamma$ Be $0.102 \pm 0.013$ 

345

85 AMMAR

91

CLEO

 $e^+ e^- \approx 10.5$  GeV $0.096 \pm 0.018 \pm 0.007$ 

66

ANJOS

91

E691

 $\gamma$ Be 80–240 GeV

85 AMMAR 91 finds  $1.25 \pm 0.25 \pm 0.25$   $\rho^0$ 's per  $\pi^+ \pi^+ \pi^- \pi^-$  decay, but can't untangle the resonant substructure ( $\rho^0 \rho^0$ ,  $a_1^\pm \pi^\mp$ ,  $\rho^0 \pi^+ \pi^-$ ).

 $\Gamma(\pi^+ \pi^- 2\pi^0)/\Gamma(K^- \pi^+)$  $\Gamma_{141}/\Gamma_{29}$ VALUE (units  $10^{-2}$ ) EVTS **$25.8 \pm 1.5 \pm 1.8$**  $2724 \pm 166$ DOCUMENT ID

RUBIN

06

CLEO

 $e^+ e^-$  at  $\psi(3770)$ 

|

 $\Gamma(\eta \pi^0)/\Gamma(K^- \pi^+)$  $\Gamma_{142}/\Gamma_{29}$ Unseen decay modes of the  $\eta$  are included.VALUE (units  $10^{-2}$ ) EVTS **$1.47 \pm 0.34 \pm 0.11$**  $62 \pm 14$ DOCUMENT ID

RUBIN

06

CLEO

 $e^+ e^-$  at  $\psi(3770)$ 

|

 $\Gamma(\omega \pi^0)/\Gamma_{\text{total}}$  $\Gamma_{143}/\Gamma$ Unseen decay modes of the  $\omega$  are included.VALUE CL% **$<2.6 \times 10^{-4}$** 

90

DOCUMENT ID

RUBIN

06

CLEO

 $e^+ e^-$  at  $\psi(3770)$ 

|

 $\Gamma(2\pi^+ 2\pi^- \pi^0)/\Gamma(K^- \pi^+)$  $\Gamma_{144}/\Gamma_{29}$ VALUE (units  $10^{-2}$ ) EVTS **$10.7 \pm 1.2 \pm 0.5$**  $1614 \pm 171$ DOCUMENT ID

RUBIN

06

CLEO

 $e^+ e^-$  at  $\psi(3770)$ 

|

 $\Gamma(\eta \pi^+ \pi^-)/\Gamma_{\text{total}}$  $\Gamma_{145}/\Gamma$ Unseen decay modes of the  $\eta$  are included.VALUE CL% **$<1.9 \times 10^{-3}$** 

90

DOCUMENT ID

RUBIN

06

CLEO

 $e^+ e^-$  at  $\psi(3770)$ 

|

 $\Gamma(\omega \pi^+ \pi^-)/\Gamma(K^- \pi^+)$  $\Gamma_{146}/\Gamma_{29}$ Unseen decay modes of the  $\omega$  are included.VALUE (units  $10^{-2}$ ) EVTS **$4.1 \pm 1.2 \pm 0.4$**  $472 \pm 132$ DOCUMENT ID

RUBIN

06

CLEO

 $e^+ e^-$  at  $\psi(3770)$ 

|

 $\Gamma(3\pi^+ 3\pi^-)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$  $\Gamma_{147}/\Gamma_{55}$ VALUE (units  $10^{-3}$ ) EVTS **$5.23 \pm 0.59 \pm 1.35$**  $149 \pm 17$ DOCUMENT ID

LINK

04B

FOCS

 $\gamma$ A,  $\bar{E}_\gamma \approx 180$  GeV

|

$\Gamma(3\pi^+ 3\pi^-)/\Gamma(K^- 3\pi^+ 2\pi^-)$

$\Gamma_{147}/\Gamma_{85}$

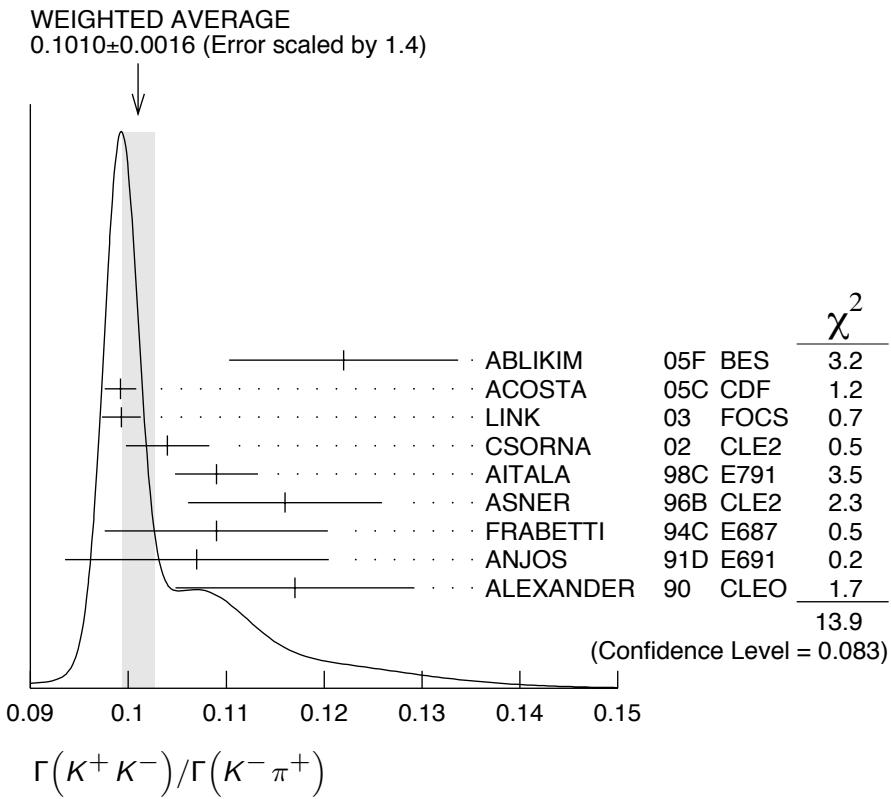
VALUE	DOCUMENT ID	TECN	COMMENT
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>			
1.93 $\pm$ 0.047 $\pm$ 0.48	86 LINK	04B FOCS	$\gamma A, \bar{E}_\gamma \approx 180$ GeV
86 This LINK 04B result is not independent of other results in these Listings.			

———— Hadronic modes with a  $K\bar{K}$  pair ———

$\Gamma(K^+ K^-)/\Gamma(K^- \pi^+)$

$\Gamma_{148}/\Gamma_{29}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.1010 <math>\pm</math> 0.0016 OUR AVERAGE</b>				Error includes scale factor of 1.4. See the ideogram below.
0.122 $\pm$ 0.011 $\pm$ 0.004	242 $\pm$ 20	ABLIKIM	05F BES	$e^+ e^- \approx \psi(3770)$
0.0992 $\pm$ 0.0011 $\pm$ 0.0012	16k $\pm$ 200	ACOSTA	05C CDF	$p\bar{p}, \sqrt{s}=1.96$ TeV
0.0993 $\pm$ 0.0014 $\pm$ 0.0014	11k	LINK	03 FOCS	$\gamma$ nucleus, $\bar{E}_\gamma \approx 180$ GeV
0.1040 $\pm$ 0.0033 $\pm$ 0.0027	1900	CSORNA	02 CLE2	$e^+ e^- \approx \Upsilon(4S)$
0.109 $\pm$ 0.003 $\pm$ 0.003	3317	AITALA	98C E791	$\pi^-$ nucleus, 500 GeV
0.116 $\pm$ 0.007 $\pm$ 0.007	1102	ASNER	96B CLE2	$e^+ e^- \approx \Upsilon(4S)$
0.109 $\pm$ 0.007 $\pm$ 0.009	581	FRABETTI	94C E687	$\gamma Be, \bar{E}_\gamma = 220$ GeV
0.107 $\pm$ 0.010 $\pm$ 0.009	193	ANJOS	91D E691	Photoproduction
0.117 $\pm$ 0.010 $\pm$ 0.007	249	ALEXANDER	90 CLEO	$e^+ e^-$ 10.5–11 GeV
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
0.107 $\pm$ 0.029 $\pm$ 0.015	103	ADAMOVICH	92 OMEG	$\pi^-$ 340 GeV
0.138 $\pm$ 0.027 $\pm$ 0.010	155	FRABETTI	92 E687	$\gamma Be$
0.16 $\pm$ 0.05	34	ALVAREZ	91B NA14	Photoproduction
0.10 $\pm$ 0.02 $\pm$ 0.01	131	ALBRECHT	90C ARG	$e^+ e^- \approx 10$ GeV
0.122 $\pm$ 0.018 $\pm$ 0.012	118	BALTRUSAIT..85E	MRK3	$e^+ e^-$ 3.77 GeV
0.113 $\pm$ 0.030		ABRAMS	79D MRK2	$e^+ e^-$ 3.77 GeV



### $\Gamma(K^+ K^-)/\Gamma(\pi^+ \pi^-)$

### $\Gamma_{148}/\Gamma_{130}$

The unused results here are redundant with  $\Gamma(K^+ K^-)/\Gamma(K^- \pi^+)$  and  $\Gamma(\pi^+ \pi^-)/\Gamma(K^- \pi^+)$  measurements by the same experiments.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
2.760±0.040±0.034	7334	ACOSTA	05C CDF	$p\bar{p}$ , $\sqrt{s}=1.96$ TeV
2.81 ± 0.10 ± 0.06		LINK	03 FOCS	$\gamma$ nucleus, $\bar{E}_\gamma \approx 180$ GeV
2.96 ± 0.16 ± 0.15	710	CSORNA	02 CLE2	$e^+ e^- \approx \Upsilon(4S)$
2.75 ± 0.15 ± 0.16		ITALA	98C E791	$\pi^-$ nucleus, 500 GeV
2.53 ± 0.46 ± 0.19		FRAZETTI	94C E687	$\gamma$ Be $\bar{E}_\gamma = 220$ GeV
2.23 ± 0.81 ± 0.46		ADAMOVICH	92 OMEG	$\pi^-$ 340 GeV
1.95 ± 0.34 ± 0.22		ANJOS	91D E691	Photoproduction
2.5 ± 0.7		ALBRECHT	90C ARG	$e^+ e^- \approx 10$ GeV
2.35 ± 0.37 ± 0.28		ALEXANDER	90 CLEO	$e^+ e^-$ 10.5–11 GeV

### $\Gamma(2K_S^0)/\Gamma(K_S^0 \pi^+ \pi^-)$

### $\Gamma_{149}/\Gamma_{31}$

This is the same as  $\Gamma(K^0 \bar{K}^0)/\Gamma(\bar{K}^0 \pi^+ \pi^-)$  because  $D^0 \rightarrow K_S^0 K_L^0$  is forbidden by  $CP$  conservation.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.0126±0.0022 OUR AVERAGE</b>				
0.0144±0.0032±0.0016	79 ± 17	LINK	05A FOCS	$\gamma$ Be, $\bar{E}_\gamma \approx 180$ GeV
0.0101±0.0022±0.0016	26	ASNER	96B CLE2	$e^+ e^- \approx \Upsilon(4S)$
0.039 ± 0.013 ± 0.013	20 ± 7	FRAZETTI	94J E687	$\gamma$ Be $\bar{E}_\gamma = 220$ GeV
0.021 $^{+0.011}_{-0.008}$ ± 0.002	5	ALEXANDER	90 CLEO	$e^+ e^-$ 10.5–11 GeV

$\Gamma(K_S^0 K^- \pi^+)/\Gamma(K^- \pi^+)$   $\Gamma_{150}/\Gamma_{29}$ 

VALUE	DOCUMENT ID	TECN	COMMENT
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**0.089±0.014 OUR FIT** Error includes scale factor of 1.1.**0.08 ±0.03** <sup>87</sup> ANJOS 91 E691  $\gamma$ Be 80–240 GeV

87 The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

 $\Gamma(K_S^0 K^- \pi^+)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{150}/\Gamma_{31}$ 

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.117±0.017 OUR FIT** Error includes scale factor of 1.1.**0.119±0.021 OUR AVERAGE** Error includes scale factor of 1.3.

0.108±0.019	61	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
0.16 ±0.03 ±0.02	39	ALBRECHT	90C ARG	$e^+ e^- \approx 10$ GeV

 $\Gamma(\bar{K}^*(892)^0 K_S^0)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{174}/\Gamma_{31}$ 
Unseen decay modes of the  $\bar{K}^*(892)^0$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
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**<0.029** 90 AMMAR 91 CLEO  $e^+ e^- \approx 10.5$  GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.03 90 ALBRECHT 90C ARG  $e^+ e^- \approx 10$  GeV
 $\Gamma(K^*(892)^+ K^-)/\Gamma(K^- \pi^+)$   $\Gamma_{175}/\Gamma_{29}$ 
Unseen decay modes of the  $K^*(892)^+$  are included.

VALUE	DOCUMENT ID	TECN	COMMENT
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**0.097±0.021 OUR FIT****0.16 +0.08  
-0.06** <sup>88</sup> ANJOS 91 E691  $\gamma$ Be 80–240 GeV

88 The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

 $\Gamma(K^*(892)^+ K^-)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{175}/\Gamma_{31}$ 
Unseen decay modes of the  $K^*(892)^+$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
-------	------	-------------	------	---------

**0.127±0.027 OUR FIT****0.117±0.028 OUR AVERAGE**

0.128±0.036	23	AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
0.10 ±0.04 ±0.02	15	ALBRECHT	90C ARG	$e^+ e^- \approx 10$ GeV

 $\Gamma(K_S^0 K^- \pi^+ \text{nonresonant})/\Gamma(K^- \pi^+)$   $\Gamma_{153}/\Gamma_{29}$ 

VALUE	DOCUMENT ID	TECN	COMMENT
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**0.03±0.03** 89 ANJOS 91 E691  $\gamma$ Be 80–240 GeV

89 The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

 $\Gamma(K_S^0 K^+ \pi^-)/\Gamma(K^- \pi^+)$   $\Gamma_{154}/\Gamma_{29}$ 

VALUE	DOCUMENT ID	TECN	COMMENT
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**0.068±0.013 OUR FIT****0.05 ±0.025** <sup>90</sup> ANJOS 91 E691  $\gamma$ Be 80–240 GeV

90 The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

 $\Gamma(K_S^0 K^+ \pi^-)/\Gamma(K_S^0 \pi^+ \pi^-)$   $\Gamma_{154}/\Gamma_{31}$ 

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
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**0.089±0.017 OUR FIT****0.098±0.020** 55 AMMAR 91 CLEO  $e^+ e^- \approx 10.5$  GeV

$\Gamma(K^*(892)^0 K_S^0)/\Gamma(K_S^0 \pi^+ \pi^-)$

$\Gamma_{176}/\Gamma_{31}$

Unseen decay modes of the  $K^*(892)^0$  are included.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.015</b>	90	AMMAR	91	CLEO $e^+ e^- \approx 10.5$ GeV

$\Gamma(K^*(892)^- K^+)/\Gamma(K_S^0 \pi^+ \pi^-)$

$\Gamma_{177}/\Gamma_{31}$

Unseen decay modes of the  $K^*(892)^-$  are included.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.068±0.038</b>	12	AMMAR	91	CLEO $e^+ e^- \approx 10.5$ GeV

$\Gamma(K_S^0 K^+ \pi^- \text{nonresonant})/\Gamma(K^- \pi^+)$

$\Gamma_{157}/\Gamma_{29}$

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.05<sup>+0.03</sup><sub>-0.02</sub></b>	91	ANJOS	91    E691 $\gamma$ Be 80–240 GeV

<sup>91</sup> The factor 100 at the top of column 2 of Table I of ANJOS 91 should be omitted.

$\Gamma(K^+ K^- \pi^0)/\Gamma(K^- \pi^+ \pi^0)$

$\Gamma_{158}/\Gamma_{43}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.0095±0.0026</b>	151	ASNER	96B	CLE2 $e^+ e^- \approx \gamma(4S)$

$\Gamma(K_S^0 K_S^0 \pi^0)/\Gamma_{\text{total}}$

$\Gamma_{159}/\Gamma$

VALUE	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.00059</b>	ASNER	96B	CLE2 $e^+ e^- \approx \gamma(4S)$

$\Gamma(\phi \pi^0)/\Gamma_{\text{total}}$

$\Gamma_{178}/\Gamma$

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
• • • We do not use the following data for averages, fits, limits, etc. • • •				

<0.0014

90

ALBRECHT    94I    ARG     $e^+ e^- \approx 10$  GeV

$\Gamma(\phi \pi^0)/\Gamma(K^+ K^-)$

$\Gamma_{178}/\Gamma_{148}$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.194±0.006±0.009</b>	1254	TAJIMA	04	BELL $e^+ e^-$ at $\gamma(4S)$

$\Gamma(\phi \eta)/\Gamma(K^+ K^-)$

$\Gamma_{179}/\Gamma_{148}$

VALUE (units $10^{-2}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>3.59±1.14±0.18</b>	31	TAJIMA	04	BELL $e^+ e^-$ at $\gamma(4S)$

$\Gamma(\phi \omega)/\Gamma_{\text{total}}$

$\Gamma_{180}/\Gamma$

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.0021</b>	90	ALBRECHT	94I	ARG $e^+ e^- \approx 10$ GeV

$\Gamma(K^+ K^- \pi^+ \pi^-)/\Gamma(K^- \pi^+ \pi^+ \pi^-)$

$\Gamma_{160}/\Gamma_{55}$

VALUE (units $10^{-2}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>3.00±0.13 OUR AVERAGE</b>				
2.95±0.11±0.08	2669 ± 101	92	LINK	05G FOCS $\gamma$ Be, $\bar{E}_\gamma \approx 180$ GeV
3.13±0.37±0.36	136 ± 15		AITALA	98D E791 $\pi^-$ nucleus, 500 GeV
3.5 ± 0.4 ± 0.2	244 ± 26		FRAZETTI	95C E687 $\gamma$ Be, $\bar{E}_\gamma \approx 200$ GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$4.4 \pm 1.8 \pm 0.5$	$19 \pm 8$	ABLIKIM	05F BES	$e^+ e^- \approx \psi(3770)$
$4.1 \pm 0.7 \pm 0.5$	$114 \pm 20$	ALBRECHT	94I ARG	$e^+ e^- \approx 10 \text{ GeV}$
$3.14 \pm 1.0$	$89 \pm 29$	AMMAR	91 CLEO	$e^+ e^- \approx 10.5 \text{ GeV}$
$2.8^{+0.8}_{-0.7}$		ANJOS	91 E691	$\gamma \text{Be } 80\text{--}240 \text{ GeV}$

92 LINK 05G uses a smaller, cleaner subset of  $1279 \pm 48$  events for the amplitude analysis that gives the results in the next data blocks.

$$\Gamma(\phi \pi^+ \pi^- 3\text{-body}, \phi \rightarrow K^+ K^-)/\Gamma(K^+ K^- \pi^+ \pi^-) \quad \Gamma_{161}/\Gamma_{160}$$

This is the fraction from a coherent amplitude analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.01 <math>\pm 0.01</math></b>	LINK	05G FOCS	$1279 \pm 48 K^+ K^- \pi^+ \pi^-$ evts.

$$\Gamma(\phi \rho^0, \phi \rightarrow K^+ K^-)/\Gamma(K^+ K^- \pi^+ \pi^-) \quad \Gamma_{162}/\Gamma_{160}$$

This is the fraction from a coherent amplitude analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.29 <math>\pm 0.02 \pm 0.01</math></b>	LINK	05G FOCS	$1279 \pm 48 K^+ K^- \pi^+ \pi^-$ evts.

$$\Gamma(K^+ K^- \rho^0 3\text{-body})/\Gamma(K^+ K^- \pi^+ \pi^-) \quad \Gamma_{163}/\Gamma_{160}$$

This is the fraction from a coherent amplitude analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.02 <math>\pm 0.02 \pm 0.02</math></b>	LINK	05G FOCS	$1279 \pm 48 K^+ K^- \pi^+ \pi^-$ evts.

$$\Gamma(f_0(980)\pi^+ \pi^-, f_0 \rightarrow K^+ K^-)/\Gamma(K^+ K^- \pi^+ \pi^-) \quad \Gamma_{164}/\Gamma_{160}$$

This is the fraction from a coherent amplitude analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.15 <math>\pm 0.03 \pm 0.02</math></b>	LINK	05G FOCS	$1279 \pm 48 K^+ K^- \pi^+ \pi^-$ evts.

$$\Gamma(K^*(892)^0 K^\mp \pi^\pm 3\text{-body}, K^{*0} \rightarrow K^\pm \pi^\mp)/\Gamma(K^+ K^- \pi^+ \pi^-) \quad \Gamma_{165}/\Gamma_{160}$$

This is the fraction from a coherent amplitude analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.11 <math>\pm 0.02 \pm 0.01</math></b>	LINK	05G FOCS	$1279 \pm 48 K^+ K^- \pi^+ \pi^-$ evts.

$$\Gamma(K^*(892)^0 \bar{K}^*(892)^0, K^{*0} \rightarrow K^\pm \pi^\mp)/\Gamma(K^+ K^- \pi^+ \pi^-) \quad \Gamma_{166}/\Gamma_{160}$$

This is the fraction from a coherent amplitude analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.03 <math>\pm 0.02 \pm 0.01</math></b>	LINK	05G FOCS	$1279 \pm 48 K^+ K^- \pi^+ \pi^-$ evts.

$$\Gamma(K_1(1270)^\pm K^\mp, K_1(1270)^\pm \rightarrow K^\pm \pi^+ \pi^-)/\Gamma(K^+ K^- \pi^+ \pi^-) \quad \Gamma_{167}/\Gamma_{160}$$

This is the fraction from a coherent amplitude analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.33 <math>\pm 0.06 \pm 0.04</math></b>	93 LINK	05G FOCS	$1279 \pm 48 K^+ K^- \pi^+ \pi^-$ evts.

93 This LINK 05G value includes  $K_1(1270)^\pm \rightarrow \rho^0 K^\pm$ ,  $\rightarrow K_0^*(1430)^0 \pi^\pm$ , and  $K^*(892)^0 \pi^\pm$ .

$$\Gamma(K_1(1400)^\pm K^\mp, K_1(1400)^\pm \rightarrow K^\pm \pi^+ \pi^-)/\Gamma(K^+ K^- \pi^+ \pi^-) \quad \Gamma_{168}/\Gamma_{160}$$

This is the fraction from a coherent amplitude analysis.

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.22 <math>\pm 0.03 \pm 0.04</math></b>	LINK	05G FOCS	$1279 \pm 48 K^+ K^- \pi^+ \pi^-$ evts.

$\Gamma(K_S^0 K_S^0 \pi^+ \pi^-)/\Gamma(K_S^0 \pi^+ \pi^-)$

$\Gamma_{171}/\Gamma_{31}$

<u>VALUE</u> (units $10^{-2}$ )	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	
<b>4.3 ± 0.8 OUR AVERAGE</b>					
$4.16 \pm 0.70 \pm 0.42$	$113 \pm 21$	LINK	05A FOCS	$\gamma$ Be, $\bar{E}_\gamma \approx 180$ GeV	
$6.2 \pm 2.0 \pm 1.6$	25	ALBRECHT	94I ARG	$e^+ e^- \approx 10$ GeV	

$\Gamma(K_S^0 K^- \pi^+ \pi^+ \pi^-)/\Gamma(K_S^0 2\pi^+ 2\pi^-)$

$\Gamma_{172}/\Gamma_{79}$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	
<0.054	90	LINK	04D FOCS	$\gamma$ A, $\bar{E}_\gamma \approx 180$ GeV	

$\Gamma(K^+ K^- \pi^+ \pi^- \pi^0)/\Gamma_{\text{total}}$

$\Gamma_{173}/\Gamma$

<u>VALUE</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	
<b><math>0.0031 \pm 0.0020</math></b>	94 BARLAG	92C ACCM	$\pi^-$ Cu	230 GeV

<sup>94</sup> BARLAG 92C computes the branching fraction using topological normalization.

———— Radiative modes ————

$\Gamma(\rho^0 \gamma)/\Gamma_{\text{total}}$

$\Gamma_{181}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	
< $2.4 \times 10^{-4}$	90	ASNER	98	CLE2

$\Gamma(\omega \gamma)/\Gamma_{\text{total}}$

$\Gamma_{182}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	
< $2.4 \times 10^{-4}$	90	ASNER	98	CLE2

$\Gamma(\phi \gamma)/\Gamma_{\text{total}}$

$\Gamma_{183}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	
• • • We do not use the following data for averages, fits, limits, etc. • • •				
< $1.9 \times 10^{-4}$	90	ASNER	98	CLE2

$\Gamma(\phi \gamma)/\Gamma(K^+ K^-)$

$\Gamma_{183}/\Gamma_{148}$

<u>VALUE</u> (units $10^{-3}$ )	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	
$6.31^{+1.70+0.30}_{-1.48-0.36}$	28	TAJIMA	04	BELL	$e^+ e^-$ at $\gamma(4S)$

$\Gamma(\bar{K}^*(892)^0 \gamma)/\Gamma_{\text{total}}$

$\Gamma_{184}/\Gamma$

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	
< $7.6 \times 10^{-4}$	90	ASNER	98	CLE2

———— Doubly Cabibbo-suppressed / Mixing modes ————

$\Gamma(K^+ \ell^- \bar{\nu}_\ell (\text{via } \bar{D}^0))/\Gamma(K^- \ell^+ \nu_\ell)$

$\Gamma_{185}/\Gamma_{14}$

This is a limit on  $R_M$  without the complications of possible doubly-Cabibbo-suppressed decays that occur when using hadronic modes. For the limits on  $|m_1 - m_2|$  and  $(\Gamma_1 - \Gamma_2)/\Gamma$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>	
<0.005	90	95 AITALA	96C E791	$\pi^-$ nucleus, 500 GeV	

<sup>95</sup> AITALA 96C uses  $D^{*+} \rightarrow D^0 \pi^+$  (and charge conjugate) decays to identify the charm at production and  $D^0 \rightarrow K^- \ell^+ \nu_\ell$  (and charge conjugate) decays to identify the charm at decay.

$$\Gamma(K^+ \text{ or } K^*(892)^+ e^- \bar{\nu}_e (\text{via } \bar{D}^0)) / [\Gamma(K^- e^+ \nu_e) + \Gamma(K^*(892)^- e^+ \nu_e)] = \Gamma_{186}/(\Gamma_{15} + \Gamma_{17})$$

This is a limit on  $R_M$  without the complications of possible doubly-Cabibbo-suppressed decays that occur when using hadronic modes. For the limits on  $|m_1 - m_2|$  and  $(\Gamma_1 - \Gamma_2)/\Gamma$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<0.001	90	96 BITENC	05 BELL	$e^+ e^- \approx 10.6 \text{ GeV}$
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>				
<0.0078	90	96 CAWLFIELD	05 CLEO	$e^+ e^- \approx 10.6 \text{ GeV}$
<0.0042	90	96 AUBERT,B	04Q BABR	$e^+ e^- \approx \gamma(4S)$

96 AUBERT,B 04Q, CAWLFIELD 05, and BITENC 05 use  $D^{*+} \rightarrow D^0 \pi^+$  (and charge conjugate) decays to identify the charm at production and the charge of the  $e$  to identify the charm at decay. These limits do not allow  $CP$  violation.

$$\Gamma(K^+ \pi^-) / \Gamma(K^- \pi^+) = \Gamma_{187}/\Gamma_{29}$$

This is  $R_D$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing,” near the start of the  $D^0$  Listings. The experiments here use the charge of the pion in  $D^*(2010)^\pm \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^\pm$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born. The  $D^0 \rightarrow K^+ \pi^-$  decay can occur directly by doubly Cabibbo-suppressed (DCS) decay, or indirectly by  $D^0 \rightarrow \bar{D}^0$  mixing followed by  $\bar{D}^0 \rightarrow K^+ \pi^-$  decay. Some of the experiments can use the decay-time information to disentangle the two mechanisms. Here, we list the experimental branching ratio, which if there is no mixing is the DCS ratio. See the next data block for limits on the mixing ratio  $R_M$ , see the section on  $CP$ -violating asymmetries near the end of this  $D^0$  Listing for values of  $A_D$ , and see the note on “ $D^0$ - $\bar{D}^0$  Mixing” for limits on  $x'$  and  $y'$ .

Some early limits have been omitted from this Listing; see our 1998 edition (EPJ **C3** 1).

VALUE (units $10^{-3}$ )	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>3.76 ± 0.09 OUR AVERAGE</b>					
3.77 ± 0.08 ± 0.05	4024 ± 88	97 ZHANG	06 BELL	$e^+ e^-$	
4.29 $^{+0.63}_{-0.61}$ ± 0.27	234	98 LINK	05H FOCS	$\gamma$ nucleus	
3.59 ± 0.20 ± 0.27		99 AUBERT	03Z BABR	$e^+ e^-$ , 10.6 GeV	
3.32 $^{+0.63}_{-0.65}$ ± 0.40	45	100 GODANG	00 CLE2	$e^+ e^-$	
6.8 $^{+3.4}_{-3.3}$ ± 0.7	34	101 AITALA	98 E791	$\pi^-$ nucl., 500 GeV	
<b>• • • We do not use the following data for averages, fits, limits, etc. • • •</b>					
3.81 ± 0.17 $^{+0.08}_{-0.16}$	845 ± 40	102 LI	05A BELL	See ZHANG 06	
4.04 ± 0.85 ± 0.25	149	103 LINK	01 FOCS	$\gamma$ nucleus	
18.4 ± 5.9 ± 3.4	19	104 BARATE	98W ALEP	$e^+ e^-$ at $Z^0$	
7.7 ± 2.5 ± 2.5	19	105 CINABRO	94 CLE2	$e^+ e^- \approx \gamma(4S)$	
< 11	90	105 AMMAR	91 CLEO	$e^+ e^- \approx 10.5 \text{ GeV}$	
< 15	90	106 ANJOS	88C E691	Photoproduction	
< 14	90	105 ALBRECHT	87K ARG	$e^+ e^-$ 10 GeV	

- 97 This ZHANG 06 result assumes no mixing. If mixing but no  $CP$  violation is allowed,  $R_D = (3.64 \pm 0.17) \times 10^{-3}$ .
- 98 This LINK 05H result assumes no mixing or  $CP$  violation. Allowing  $CP$  violation but no mixing,  $R_D = (4.29 \pm 0.63 \pm 0.28) \times 10^{-3}$  — negligibly different. Allowing mixing but no  $CP$  violation,  $R_D = (3.81^{+1.67}_{-1.63} \pm 0.92) \times 10^{-3}$ . Allowing mixing and  $CP$  violation,  $R_D = (5.17^{+1.47}_{-1.58} \pm 0.76) \times 10^{-3}$ .
- 99 This AUBERT 03Z result is for no mixing or  $CP$  violation. If  $CP$  violation but no mixing is allowed,  $R_D = 0.00357 \pm 0.00022 \pm 0.00027$ . If only mixing is allowed, the 95% confidence-level interval is  $(2.4 < R_D < 4.9) \times 10^{-3}$ . If both mixing and  $CP$  violation are allowed, this interval becomes  $(2.3 < R_D < 5.2) \times 10^{-3}$ .
- 100 This GODANG 00 result assumes no  $D^0$ - $\bar{D}^0$  mixing ( $R_M = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings) but allows  $CP$  violation. The DCS ratio becomes  $0.0048 \pm 0.0012 \pm 0.0004$  when mixing is allowed.
- 101 This AITALA 98 result assumes no  $CP$  violation or mixing ( $R_M = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). The DCS ratio becomes  $0.0090^{+0.0120}_{-0.0109} \pm 0.0044$  when mixing is allowed.
- 102 This LI 05A result assumes no mixing or  $CP$  violation. If mixing but no  $CP$  violation is allowed,  $R_D = (2.87 \pm 0.37) \times 10^{-3}$ .
- 103 This LINK 01 result assumes no mixing or  $CP$  violation; see Fig. 4 of the paper for the DCS value as a function of the (unknown) mixing parameters  $x'$  and  $y'$ . See also the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings for results on  $x'$  and  $y'$  from FOCUS and other experiments.
- 104 BARATE 98W gets  $0.0177^{+0.0060}_{-0.0056} \pm 0.0031$  for the DCS ratio when mixing is allowed, assuming no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings).
- 105 CINABRO 94, AMMAR 91, and ALBRECHT 87K cannot distinguish between doubly Cabibbo-suppressed decay and  $D^0$ - $\bar{D}^0$  mixing.
- 106 ANJOS 88C allows mixing but assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.049.

### $\Gamma(K^+\pi^-(\text{via } \bar{D}^0))/\Gamma(K^-\pi^+)$

### $\Gamma_{188}/\Gamma_{29}$

This is  $R_M$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings. The experiments here (1) use the charge of the pion in  $D^*(2010)^\pm \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^\pm$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born; and (2) use the decay-time distribution to disentangle doubly Cabibbo-suppressed decay and mixing. For the limits on  $|m_1 - m_2|$  and  $(\Gamma_1 - \Gamma_2)/\Gamma$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<0.00040	95	107	ZHANG	06	BELL $e^+e^-$
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>					
<0.00046	95	108	LI	05A	BELL See ZHANG 06
<0.0063	95	109	LINK	05H	FOCS $\gamma$ nucleus
<0.0013	95	110	AUBERT	03Z	BABR $e^+e^-$ , 10.6 GeV
<0.00041	95	111	GODANG	00	CLE2 $e^+e^-$
<0.0092	95	112	BARATE	98W	ALEP $e^+e^-$ at $Z^0$
<0.005	90	1 ± 4	ANJOS	88C	E691 Photoproduction

- 107 This ZHANG 06 result allows  $CP$  violation, but the result does not change if  $CP$  violation is not allowed.  
 108 This LI 05A result allows  $CP$  violation. The limit becomes  $< 0.00042$  (95% CL) if  $CP$  violation is not allowed.  
 109 LINK 05H obtains the same result whether or not  $CP$  violation is allowed.  
 110 This AUBERT 03Z result allows  $CP$  violation and assumes that the strong phase between  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^+ \pi^-$  is small, and limits only  $D^0 \rightarrow \bar{D}^0$  transitions via off-shell intermediate states. The limit on transitions via on-shell intermediate states is 0.0016.  
 111 This GODANG 00 result allows  $CP$  violation and assumes that the strong phase between  $D^0 \rightarrow K^+ \pi^-$  and  $\bar{D}^0 \rightarrow K^+ \pi^-$  is small, and limits only  $D^0 \rightarrow \bar{D}^0$  transitions via off-shell intermediate states. The limit on transitions via on-shell intermediate states is 0.0017.  
 112 This BARATE 98W result assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.036 (95%CL).  
 113 This ANJOS 88C result assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.019.

 $\Gamma(K_S^0 \pi^+ \pi^- \text{ (in } D^0 \rightarrow \bar{D}^0\text{)}) / \Gamma(K_S^0 \pi^+ \pi^-)$  $\Gamma_{189}/\Gamma_{31}$ 

This is  $R_M$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings. The experiments here (1) use the charge of the pion in  $D^*(2010)^{\pm} \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^{\pm}$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born; and (2) use the decay-time distribution to disentangle doubly Cabibbo-suppressed decay and mixing. For the limits on  $|m_1 - m_2|$  and  $(\Gamma_1 - \Gamma_2)/\Gamma$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.0063</b>	95	114 ASNER	05 CLEO	$e^+ e^- \approx 10 \text{ GeV}$

- 114 This ASNER 05 limit allows  $CP$  violation. If  $CP$  violation is not allowed, the limit is 0.0042 at 95% CL.

 $\Gamma(K^+ \pi^- \pi^0) / \Gamma(K^- \pi^+ \pi^0)$  $\Gamma_{191}/\Gamma_{43}$ 

The experiments here use the charge of the pion in  $D^*(2010)^{\pm} \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^{\pm}$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born. The  $D^0 \rightarrow K^+ \pi^- \pi^0$  decay can occur directly by doubly Cabibbo-suppressed (DCS) decay, or indirectly by  $D^0 \rightarrow \bar{D}^0$  mixing followed by  $\bar{D}^0 \rightarrow K^+ \pi^- \pi^0$  decay.

VALUE (units $10^{-3}$ )	EVTS	DOCUMENT ID	TECN	COMMENT
<b>2.34<sup>+0.20</sup><sub>-0.17</sub> OUR AVERAGE</b>				
$2.29 \pm 0.15^{+0.13}_{-0.09}$	$1978 \pm 104$	TIAN	05 BELL	$e^+ e^- \approx \gamma(4S)$
$4.3^{+1.1}_{-1.0} \pm 0.7$	38	BRANDENB...	01 CLE2	$e^+ e^- \approx \gamma(4S)$

 $\Gamma(K^+ \pi^- \pi^+ \pi^-) / \Gamma(K^- \pi^+ \pi^+ \pi^-)$  $\Gamma_{192}/\Gamma_{55}$ 

The experiments here use the charge of the pion in  $D^*(2010)^{\pm} \rightarrow (D^0 \text{ or } \bar{D}^0) \pi^{\pm}$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born. The  $D^0 \rightarrow K^+ \pi^- \pi^+ \pi^-$  decay can occur directly by doubly Cabibbo-suppressed (DCS) decay, or indirectly by  $D^0 \rightarrow \bar{D}^0$  mixing followed by  $\bar{D}^0 \rightarrow K^+ \pi^- \pi^+ \pi^-$  decay. Some of the experiments can use the decay-time information to disentangle the two mechanisms. Here, we list the

experimental branching ratio, which if there is no mixing is the DCS ratio; in the next data block we give the limits on the mixing ratio.

Some early limits have been omitted from this Listing; see our 1998 edition (EPJ **C3** 1).

<u>VALUE</u> (units $10^{-3}$ )	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>3.23<math>^{+0.25}_{-0.22}</math> OUR AVERAGE</b>					
$3.20 \pm 0.18^{+0.18}_{-0.13}$		$1721 \pm 75$	115 TIAN	05 BELL	$e^+ e^- \approx \gamma(4S)$
$4.4^{+1.3}_{-1.2} \pm 0.6$		54	115 DYTMAN	01 CLE2	$e^+ e^- \approx \gamma(4S)$
$2.5^{+3.6}_{-3.4} \pm 0.3$			116 AITALA	98 E791	$\pi^-$ nucl., 500 GeV
• • • We do not use the following data for averages, fits, limits, etc. • • •					
<18	90		115 AMMAR	91 CLEO	$e^+ e^- \approx 10.5$ GeV
<18	90	$5 \pm 12$	117 ANJOS	88C E691	Photoproduction

115 AMMAR 91 cannot and DYTMAN 01 and TIAN 05 do not distinguish between doubly Cabibbo-suppressed decay and  $D^0$ - $\bar{D}^0$  mixing.

116 This AITALA 98 result assumes no  $D^0$ - $\bar{D}^0$  mixing ( $R_M$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing”). It becomes  $-0.0020^{+0.0117}_{-0.0106} \pm 0.0035$  when mixing is allowed and decay-time information is used to distinguish doubly Cabibbo-suppressed decays from mixing.

117 ANJOS 88C uses decay-time information to distinguish doubly Cabibbo-suppressed (DCS) decays from  $D^0$ - $\bar{D}^0$  mixing. However, the result assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.033.

### $\Gamma(K^+\pi^-\pi^+\pi^- \text{ (via } \bar{D}^0\text{)})/\Gamma(K^-\pi^+\pi^+\pi^-)$ $\Gamma_{193}/\Gamma_{55}$

This is a  $D^0$ - $\bar{D}^0$  mixing limit. The experiments here (1) use the charge of the pion in  $D^*(2010)^\pm \rightarrow (D^0 \text{ or } \bar{D}^0)\pi^\pm$  decay to tell whether a  $D^0$  or a  $\bar{D}^0$  was born; and (2) use the decay-time distribution to disentangle doubly Cabibbo-suppressed decay and mixing. For the limits on  $|m_{D_1^0} - m_{D_2^0}|$  and  $(\Gamma_{D_1^0} - \Gamma_{D_2^0})/\Gamma_{D^0}$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
<b>&lt;0.005</b>	90	$0 \pm 4$	118 ANJOS	88C E691	Photoproduction

118 ANJOS 88C uses decay-time information to distinguish doubly Cabibbo-suppressed (DCS) decays from  $D^0$ - $\bar{D}^0$  mixing. However, the result assumes no interference between the DCS and mixing amplitudes ( $y' = 0$  in the note on “ $D^0$ - $\bar{D}^0$  Mixing” near the start of the  $D^0$  Listings). When interference is allowed, the limit degrades to 0.007.

### $\Gamma(K^+\pi^- \text{ or } K^+\pi^-\pi^+\pi^- \text{ (via } \bar{D}^0\text{)})/\Gamma(K^-\pi^+ \text{ or } K^-\pi^+\pi^+\pi^-)$ $\Gamma_{194}/\Gamma_0$

This is a  $D^0$ - $\bar{D}^0$  mixing limit. For the limits on  $|m_{D_1^0} - m_{D_2^0}|$  and  $(\Gamma_{D_1^0} - \Gamma_{D_2^0})/\Gamma_{D^0}$  that come from the best mixing limit, see near the beginning of these  $D^0$  Listings.

<u>VALUE</u>	<u>CL%</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
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• • • We do not use the following data for averages, fits, limits, etc. • • •

<0.0085	90	119 AITALA	98 E791	$\pi^-$ nucleus, 500 GeV
<0.0037	90	120 ANJOS	88C E691	Photoproduction

- 119 AITALA 98 uses decay-time information to distinguish doubly Cabibbo-suppressed decays from  $D^0-\bar{D}^0$  mixing. The fit allows interference between the two amplitudes, and also allows  $CP$  violation in this term. The central value obtained is  $0.0039^{+0.0036}_{-0.0032} \pm 0.0016$ . When interference is disallowed, the result becomes  $0.0021 \pm 0.0009 \pm 0.0002$ .
- 120 This combines results of ANJOS 88C on  $K^+\pi^-$  and  $K^+\pi^-\pi^+\pi^-$  (via  $\bar{D}^0$ ) reported in the data block above (see footnotes there). It assumes no interference.

 **$\Gamma(\mu^- \text{anything (via } \bar{D}^0\text{)})/\Gamma(\mu^+ \text{anything})$**  **$\Gamma_{195}/\Gamma_6$** This is a  $D^0-\bar{D}^0$  mixing limit. See the somewhat better limits above.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.0056</b>	90	LOUIS	86	SPEC $\pi^-$ W 225 GeV
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>				
<0.012	90	BENVENUTI	85	CNTR $\mu$ C, 200 GeV
<0.044	90	BODEK	82	SPEC $\pi^-$ , $p$ Fe $\rightarrow D^0$

**Rare or forbidden modes** **$\Gamma(\gamma\gamma)/\Gamma(\pi^0\pi^0)$**  **$\Gamma_{196}/\Gamma_{131}$**  $D^0 \rightarrow \gamma\gamma$  is a flavor-changing neutral-current decay, forbidden in the Standard Model at the tree level.

VALUE	CL%	DOCUMENT ID	TECN	COMMENT
<b>&lt;0.033</b>	90	COAN	03	CLE2 $e^+e^- \approx \gamma(4S)$

 **$\Gamma(e^+e^-)/\Gamma_{\text{total}}$**  **$\Gamma_{197}/\Gamma$** A test for the  $\Delta C = 1$  weak neutral current. Allowed by first-order weak interaction combined with electromagnetic interaction.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>&lt;1.2 × 10<sup>-6</sup></b>	90	3	AUBERT,B	04Y BABR	$e^+e^- \approx \gamma(4S)$
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>					
<8.19 × 10 <sup>-6</sup>	90		PRIPSTEIN	00 E789	$p$ nucleus, 800 GeV
<6.2 × 10 <sup>-6</sup>	90		AITALA	99G E791	$\pi^- N$ 500 GeV
<1.3 × 10 <sup>-5</sup>	90	0	FREYBERGER	96 CLE2	$e^+e^- \approx \gamma(4S)$
<1.3 × 10 <sup>-4</sup>	90		ADLER	88 MRK3	$e^+e^-$ 3.77 GeV
<1.7 × 10 <sup>-4</sup>	90	7	ALBRECHT	88G ARG	$e^+e^-$ 10 GeV
<2.2 × 10 <sup>-4</sup>	90	8	HAAS	88 CLEO	$e^+e^-$ 10 GeV

 **$\Gamma(\mu^+\mu^-)/\Gamma_{\text{total}}$**  **$\Gamma_{198}/\Gamma$** A test for the  $\Delta C = 1$  weak neutral current. Allowed by first-order weak interaction combined with electromagnetic interaction.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
<b>&lt;1.3 × 10<sup>-6</sup></b>	90	1	AUBERT,B	04Y BABR	$e^+e^- \approx \gamma(4S)$
<b>• • •</b> We do not use the following data for averages, fits, limits, etc. <b>• • •</b>					
<2.0 × 10 <sup>-6</sup>	90		ABT	04 HERB	$pA$ , 920 GeV
<2.5 × 10 <sup>-6</sup>	90		ACOSTA	03F CDF	$p\bar{p}$ , $\sqrt{s} = 1.96$ TeV
<1.56 × 10 <sup>-5</sup>	90		PRIPSTEIN	00 E789	$p$ nucleus, 800 GeV
<5.2 × 10 <sup>-6</sup>	90		AITALA	99G E791	$\pi^- N$ 500 GeV
<4.1 × 10 <sup>-6</sup>	90		ADAMOVICH	97 BEAT	$\pi^-$ Cu, W 350 GeV
<4.2 × 10 <sup>-6</sup>	90		ALEXOPOU...	96 E771	$p$ Si, 800 GeV
<3.4 × 10 <sup>-5</sup>	90	1	FREYBERGER	96 CLE2	$e^+e^- \approx \gamma(4S)$

$<7.6 \times 10^{-6}$	90	0	ADAMOVICH	95	BEAT	See ADAMOVICH 97
$<4.4 \times 10^{-5}$	90	0	KODAMA	95	E653	$\pi^-$ emulsion 600 GeV
$<3.1 \times 10^{-5}$	90	121	MISHRA	94	E789	$-4.1 \pm 4.8$ events
$<7.0 \times 10^{-5}$	90	3	ALBRECHT	88G	ARG	$e^+ e^-$ 10 GeV
$<1.1 \times 10^{-5}$	90		LOUIS	86	SPEC	$\pi^-$ W 225 GeV
$<3.4 \times 10^{-4}$	90		AUBERT	85	EMC	Deep inelast. $\mu^- N$

121 Here MISHRA 94 uses "the statistical approach advocated by the PDG." For an alternate approach, giving a limit of  $9 \times 10^{-6}$  at 90% confidence level, see the paper.

### $\Gamma(\pi^0 e^+ e^-)/\Gamma_{\text{total}}$

### $\Gamma_{199}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<4.5 \times 10^{-5}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

### $\Gamma(\pi^0 \mu^+ \mu^-)/\Gamma_{\text{total}}$

### $\Gamma_{200}/\Gamma$

A test for the  $\Delta C=1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.8 \times 10^{-4}$	90	2	KODAMA	95	E653 $\pi^-$ emulsion 600 GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$					
$<5.4 \times 10^{-4}$	90	3	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

### $\Gamma(\eta e^+ e^-)/\Gamma_{\text{total}}$

### $\Gamma_{201}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.1 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

### $\Gamma(\eta \mu^+ \mu^-)/\Gamma_{\text{total}}$

### $\Gamma_{202}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<5.3 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

### $\Gamma(\pi^+ \pi^- e^+ e^-)/\Gamma_{\text{total}}$

### $\Gamma_{203}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<3.73 \times 10^{-4}$	90	9	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

### $\Gamma(\rho^0 e^+ e^-)/\Gamma_{\text{total}}$

### $\Gamma_{204}/\Gamma$

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.0 \times 10^{-4}$	90	2	122 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\bullet \bullet \bullet$  We do not use the following data for averages, fits, limits, etc.  $\bullet \bullet \bullet$

$<1.24 \times 10^{-4}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV
$<4.5 \times 10^{-4}$	90	2	HAAS	88 CLEO	$e^+ e^-$ 10 GeV

122 This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 1.8 \times 10^{-4}$  using a photon pole amplitude model.

$\Gamma(\pi^+\pi^-\mu^+\mu^-)/\Gamma_{\text{total}}$  $\Gamma_{205}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<3.0 \times 10^{-5}$	90	2	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(\rho^0\mu^+\mu^-)/\Gamma_{\text{total}}$  $\Gamma_{206}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<2.2 \times 10^{-5}$	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$					
$<4.9 \times 10^{-4}$	90	1	123 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$
$<2.3 \times 10^{-4}$	90	0	KODAMA	95 E653	$\pi^-$ emulsion 600 GeV
$<8.1 \times 10^{-4}$	90	5	HAAS	88 CLEO	$e^+ e^-$ 10 GeV

123 This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 4.5 \times 10^{-4}$  using a photon pole amplitude model.

 $\Gamma(\omega e^+ e^-)/\Gamma_{\text{total}}$  $\Gamma_{207}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.8 \times 10^{-4}$	90	1	124 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

124 This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 2.7 \times 10^{-4}$  using a photon pole amplitude model.

 $\Gamma(\omega\mu^+\mu^-)/\Gamma_{\text{total}}$  $\Gamma_{208}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<8.3 \times 10^{-4}$	90	0	125 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

125 This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 6.5 \times 10^{-4}$  using a photon pole amplitude model.

 $\Gamma(K^-K^+e^+e^-)/\Gamma_{\text{total}}$  $\Gamma_{209}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<3.15 \times 10^{-4}$	90	9	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(\phi e^+ e^-)/\Gamma_{\text{total}}$  $\Gamma_{210}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<5.2 \times 10^{-5}$	90	2	126 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\bullet \bullet \bullet$  We do not use the following data for averages, fits, limits, etc.  $\bullet \bullet \bullet$

$<5.9 \times 10^{-5}$  90 0 AITALA 01C E791  $\pi^-$  nucleus, 500 GeV

126 This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 7.6 \times 10^{-5}$  using a photon pole amplitude model.

$\Gamma(K^- K^+ \mu^+ \mu^-)/\Gamma_{\text{total}}$  $\Gamma_{211}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<3.3 \times 10^{-5}$	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(\phi \mu^+ \mu^-)/\Gamma_{\text{total}}$  $\Gamma_{212}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<3.1 \times 10^{-5}$	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<4.1 \times 10^{-4}$	90	0	127 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$
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127 This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 2.4 \times 10^{-4}$  using a photon pole amplitude model.

 $\Gamma(\bar{K}^0 e^+ e^-)/\Gamma_{\text{total}}$  $\Gamma_{213}/\Gamma$ 

Not a useful test for  $\Delta C = 1$  weak neutral current because both quarks must change flavor.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.1 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<1.7 \times 10^{-3}$	90		ADLER	89C MRK3	$e^+ e^-$ 3.77 GeV
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 $\Gamma(\bar{K}^0 \mu^+ \mu^-)/\Gamma_{\text{total}}$  $\Gamma_{214}/\Gamma$ 

Not a useful test for  $\Delta C = 1$  weak neutral current because both quarks must change flavor.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<2.6 \times 10^{-4}$	90	2	KODAMA	95 E653	$\pi^-$ emulsion 600 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<6.7 \times 10^{-4}$	90	1	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$
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 $\Gamma(K^- \pi^+ e^+ e^-)/\Gamma_{\text{total}}$  $\Gamma_{215}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<3.85 \times 10^{-4}$	90	6	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(\bar{K}^*(892)^0 e^+ e^-)/\Gamma_{\text{total}}$  $\Gamma_{216}/\Gamma$ 

Not a useful test for  $\Delta C = 1$  weak neutral current because both quarks must change flavor.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<4.7 \times 10^{-5}$	90	2	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<1.4 \times 10^{-4}$	90	1	128 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$
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128 This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 2.0 \times 10^{-4}$  using a photon pole amplitude model.

$\Gamma(K^-\pi^+\mu^+\mu^-)/\Gamma_{\text{total}}$  $\Gamma_{217}/\Gamma$ 

A test for the  $\Delta C = 1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<3.59 \times 10^{-4}$	90	12	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(\bar{K}^*(892)^0\mu^+\mu^-)/\Gamma_{\text{total}}$  $\Gamma_{218}/\Gamma$ 

Not a useful test for  $\Delta C = 1$  weak neutral current because both quarks must change flavor.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<2.4 \times 10^{-5}$	90	3	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$<1.18 \times 10^{-3}$  90 1 129 FREYBERGER 96 CLE2  $e^+ e^- \approx \gamma(4S)$

129 This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 1.0 \times 10^{-3}$  using a photon pole amplitude model.

 $\Gamma(\pi^+\pi^-\pi^0\mu^+\mu^-)/\Gamma_{\text{total}}$  $\Gamma_{219}/\Gamma$ 

A test for the  $\Delta C=1$  weak neutral current. Allowed by higher-order electroweak interactions.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<8.1 \times 10^{-4}$	90	1	KODAMA	95 E653	$\pi^-$ emulsion 600 GeV

 $\Gamma(\mu^\pm e^\mp)/\Gamma_{\text{total}}$  $\Gamma_{220}/\Gamma$ 

A test of lepton family number conservation.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$< 8.1 \times 10^{-7}$	90	0	AUBERT,B	04Y BABR	$e^+ e^- \approx \gamma(4S)$

• • • We do not use the following data for averages, fits, limits, etc. • • •

$< 1.72 \times 10^{-5}$	90		PRIPSTEIN	00 E789	$p$ nucleus, 800 GeV
$< 8.1 \times 10^{-6}$	90		AITALA	99G E791	$\pi^- N$ 500 GeV
$< 1.9 \times 10^{-5}$	90	2	130 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$
$< 1.0 \times 10^{-4}$	90	4	ALBRECHT	88G ARG	$e^+ e^-$ 10 GeV
$< 2.7 \times 10^{-4}$	90	9	HAAS	88 CLEO	$e^+ e^-$ 10 GeV
$< 1.2 \times 10^{-4}$	90		BECKER	87C MRK3	$e^+ e^-$ 3.77 GeV
$< 9 \times 10^{-4}$	90		PALKA	87 SILI	200 GeV $\pi p$
$< 21 \times 10^{-4}$	90	0	131 RILES	87 MRK2	$e^+ e^-$ 29 GeV

130 This is the corrected result given in the erratum to FREYBERGER 96.

131 RILES 87 assumes  $B(D \rightarrow K\pi) = 3.0\%$  and has production model dependency.

 $\Gamma(\pi^0 e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{221}/\Gamma$ 

A test of lepton family number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<8.6 \times 10^{-5}$	90	2	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

 $\Gamma(\eta e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{222}/\Gamma$ 

A test of lepton family number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<1.0 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\Gamma(\pi^+ \pi^- e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{223}/\Gamma$ 

A test of lepton family-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.5 \times 10^{-5}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(\rho^0 e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{224}/\Gamma$ 

A test of lepton family number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<4.9 \times 10^{-5}$	90	0	132 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$
$\bullet \bullet \bullet$ We do not use the following data for averages, fits, limits, etc. $\bullet \bullet \bullet$					
$<6.6 \times 10^{-5}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

<sup>132</sup>This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 5.0 \times 10^{-5}$  using a photon pole amplitude model.

 $\Gamma(\omega e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{225}/\Gamma$ 

A test of lepton family number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.2 \times 10^{-4}$	90	0	133 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

<sup>133</sup>This FREYBERGER 96 limit is obtained using a phase-space model. The same limit is obtained using a photon pole amplitude model.

 $\Gamma(K^- K^+ e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{226}/\Gamma$ 

A test of lepton family-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.8 \times 10^{-4}$	90	5	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(\phi e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{227}/\Gamma$ 

A test of lepton family number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<3.4 \times 10^{-5}$	90	0	134 FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

$\bullet \bullet \bullet$  We do not use the following data for averages, fits, limits, etc.  $\bullet \bullet \bullet$

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<4.7 \times 10^{-5}$	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

<sup>134</sup>This FREYBERGER 96 limit is obtained using a phase-space model. The limit changes to  $< 3.3 \times 10^{-5}$  using a photon pole amplitude model.

 $\Gamma(\bar{K}^0 e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{228}/\Gamma$ 

A test of lepton family number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.0 \times 10^{-4}$	90	0	FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

 $\Gamma(K^- \pi^+ e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{229}/\Gamma$ 

A test of lepton family-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<5.53 \times 10^{-4}$	90	15	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(K^*(892)^0 e^\pm \mu^\mp)/\Gamma_{\text{total}}$  $\Gamma_{230}/\Gamma$ 

A test of lepton family number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<8.3 \times 10^{-5}$	90	9	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV
$<1.0 \times 10^{-4}$	90	0	<sup>135</sup> FREYBERGER 96	CLE2	$e^+ e^- \approx \gamma(4S)$

<sup>135</sup>This FREYBERGER 96 limit is obtained using a phase-space model. The same limit is obtained using a photon pole amplitude model.

 $\Gamma(\pi^- \pi^- e^+ e^+ + \text{c.c.})/\Gamma_{\text{total}}$  $\Gamma_{231}/\Gamma$ 

A test of lepton-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.12 \times 10^{-4}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(\pi^- \pi^- \mu^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$  $\Gamma_{232}/\Gamma$ 

A test of lepton-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<2.9 \times 10^{-5}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(K^- \pi^- e^+ e^+ + \text{c.c.})/\Gamma_{\text{total}}$  $\Gamma_{233}/\Gamma$ 

A test of lepton-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<2.06 \times 10^{-4}$	90	2	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(K^- \pi^- \mu^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$  $\Gamma_{234}/\Gamma$ 

A test of lepton-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<3.9 \times 10^{-4}$	90	14	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(K^- K^- e^+ e^+ + \text{c.c.})/\Gamma_{\text{total}}$  $\Gamma_{235}/\Gamma$ 

A test of lepton-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<1.52 \times 10^{-4}$	90	2	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(K^- K^- \mu^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$  $\Gamma_{236}/\Gamma$ 

A test of lepton-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<9.4 \times 10^{-5}$	90	1	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(\pi^- \pi^- \mu^+ \mu^+ + \text{c.c.})/\Gamma_{\text{total}}$  $\Gamma_{237}/\Gamma$ 

A test of lepton-number conservation. The value is for the sum of the two charge states.

<u>VALUE</u>	<u>CL%</u>	<u>EVTS</u>	<u>DOCUMENT ID</u>	<u>TECN</u>	<u>COMMENT</u>
$<7.9 \times 10^{-5}$	90	4	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

$\Gamma(K^-\pi^-e^+\mu^++\text{c.c.})/\Gamma_{\text{total}}$  $\Gamma_{238}/\Gamma$ 

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<2.18 \times 10^{-4}$	90	7	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 $\Gamma(K^+K^-e^+\mu^++\text{c.c.})/\Gamma_{\text{total}}$  $\Gamma_{239}/\Gamma$ 

A test of lepton-number conservation. The value is for the sum of the two charge states.

VALUE	CL%	EVTS	DOCUMENT ID	TECN	COMMENT
$<5.7 \times 10^{-5}$	90	0	AITALA	01C E791	$\pi^-$ nucleus, 500 GeV

 **$D^0$  CP-VIOLATING DECAY-RATE ASYMMETRIES** $A_{CP}(K^+K^-)$  in  $D^0, \bar{D}^0 \rightarrow K^+K^-$ 

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.014±0.010 OUR AVERAGE</b>				
+0.020±0.012±0.006	136	ACOSTA	05C CDF	$p\bar{p}, \sqrt{s}=1.96$ TeV
0.000±0.022±0.008	3023	136 CSORNA	02 CLE2	$e^+e^- \approx \Upsilon(4S)$
-0.001±0.022±0.015	3330	136 LINK	00B FOCS	
-0.010±0.049±0.012	609	136 AITALA	98C E791	$-0.093 < A_{CP} < +0.073$ (90% CL)
+0.080±0.061		BARTELTT	95 CLE2	$-0.022 < A_{CP} < +0.18$ (90% CL)
+0.024±0.084		136 FRABETTI	94I E687	$-0.11 < A_{CP} < +0.16$ (90% CL)

136 FRABETTI 94I, AITALA 98C, LINK 00B, CSORNA 02, and ACOSTA 05C measure  $N(D^0 \rightarrow K^+K^-)/N(D^0 \rightarrow K^-\pi^+)$ , the ratio of numbers of events observed, and similarly for the  $\bar{D}^0$ .

 $A_{CP}(K_S^0K_S^0)$  in  $D^0, \bar{D}^0 \rightarrow K_S^0K_S^0$ 

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>-0.23±0.19</b>	65	BONVICINI	01 CLE2	$e^+e^- \approx 10.6$ GeV

 $A_{CP}(\pi^+\pi^-)$  in  $D^0, \bar{D}^0 \rightarrow \pi^+\pi^-$ 

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.013±0.012 OUR AVERAGE</b>				
+0.010±0.013±0.006	137	ACOSTA	05C CDF	$p\bar{p}, \sqrt{s}=1.96$ TeV
0.019±0.032±0.008	1136	137 CSORNA	02 CLE2	$e^+e^- \approx \Upsilon(4S)$
+0.048±0.039±0.025	1177	137 LINK	00B FOCS	
-0.049±0.078±0.030	343	137 AITALA	98C E791	$-0.186 < A_{CP} < +0.088$ (90% CL)

<sup>137</sup>AITALA 98C, LINK 00B, CSORNA 02, and ACOSTA 05C measure  $N(D^0 \rightarrow \pi^+ \pi^-)/N(D^0 \rightarrow K^- \pi^+)$ , the ratio of numbers of events observed, and similarly for the  $\bar{D}^0$ .

### $A_{CP}(\pi^0 \pi^0)$ in $D^0, \bar{D}^0 \rightarrow \pi^0 \pi^0$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*: D^{*+} \rightarrow D^0 \pi^+$  and  $D^{*-} \rightarrow \bar{D}^0 \pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>+0.001±0.048</b>	810	BONVICINI	01	CLE2 $e^+ e^- \approx 10.6$ GeV

### $A_{CP}(\pi^+ \pi^- \pi^0)$ in $D^0, \bar{D}^0 \rightarrow \pi^+ \pi^- \pi^0$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*: D^{*+} \rightarrow D^0 \pi^+$  and  $D^{*-} \rightarrow \bar{D}^0 \pi^-$ .

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.01<sup>+0.09</sup><sub>-0.07</sub>±0.05</b>	CRONIN-HEN..05	CLEO	$e^+ e^- \approx 10$ GeV

### $A_{CP}(K_S^0 \phi)$ in $D^0, \bar{D}^0 \rightarrow K_S^0 \phi$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*: D^{*+} \rightarrow D^0 \pi^+$  and  $D^{*-} \rightarrow \bar{D}^0 \pi^-$ .

VALUE	DOCUMENT ID	TECN	COMMENT
<b>-0.028±0.094</b>	BARTELTT	95	CLE2 $-0.182 < A_{CP} < +0.126$ (90%CL)

### $A_{CP}(K_S^0 \pi^0)$ in $D^0, \bar{D}^0 \rightarrow K_S^0 \pi^0$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*: D^{*+} \rightarrow D^0 \pi^+$  and  $D^{*-} \rightarrow \bar{D}^0 \pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>+0.001±0.013</b>	9099	BONVICINI	01	CLE2 $e^+ e^- \approx 10.6$ GeV

• • • We do not use the following data for averages, fits, limits, etc. • • •

$-0.018 \pm 0.030$  BARTELTT 95 CLE2 See BONVICINI 01

### $A_{CP}(K^\pm \pi^\mp)$ in $D^0 \rightarrow K^+ \pi^-, \bar{D}^0 \rightarrow K^- \pi^+$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*: D^{*+} \rightarrow D^0 \pi^+$  and  $D^{*-} \rightarrow \bar{D}^0 \pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.05 ± 0.04 OUR AVERAGE</b>				
$+0.023 \pm 0.047$	$4024 \pm 88$	138 ZHANG	06 BELL	$e^+ e^-$
$+0.18 \pm 0.14 \pm 0.04$		139 LINK	05H FOCS	$\gamma$ nucleus
$+0.095 \pm 0.061 \pm 0.083$		140 AUBERT	03Z BABR	$e^+ e^-$ , 10.6 GeV
$+0.02 \pm 0.19 \pm 0.01$	45	141 GODANG	00 CLE2	$-0.43 < A_{CP} < +0.34$ (95%CL)

• • • We do not use the following data for averages, fits, limits, etc. • • •

$-0.080 \pm 0.077$  845  $\pm 40$  142 LI 05A BELL See ZHANG 06

138 This ZHANG 06 result allows mixing.

139 This LINK 05H result assumes no mixing. If mixing is allowed, it becomes  $0.13^{+0.33}_{-0.25} \pm 0.10$ .

140 This AUBERT 03Z limit assumes no mixing. If mixing is allowed, the 95% confidence-level interval is  $(-2.8 < A_D < 4.9) \times 10^{-3}$ .

141 This GODANG 00 result assumes no  $D^0$ - $\bar{D}^0$  mixing; it becomes  $-0.01^{+0.16}_{-0.17} \pm 0.01$  when mixing is allowed.

142 This LI 05A result allows mixing.

### $A_{CP}(K^\mp\pi^\pm\pi^0)$ in $D^0 \rightarrow K^-\pi^+\pi^0, \bar{D}^0 \rightarrow K^+\pi^-\pi^0$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE		DOCUMENT ID	TECN	COMMENT
<b><math>-0.031 \pm 0.086</math></b>	143	KOPP	01	CLE2 $e^+e^- \approx 10.6$ GeV

143 KOPP 01 fits separately the  $D^0$  and  $\bar{D}^0$  Dalitz plots and then calculates the integrated difference of normalized densities divided by the integrated sum.

### $A_{CP}(K^\pm\pi^\mp\pi^0)$ in $D^0 \rightarrow K^+\pi^-\pi^0, \bar{D}^0 \rightarrow K^-\pi^+\pi^0$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.00 <math>\pm 0.05</math> OUR AVERAGE</b>				
$-0.006 \pm 0.053$	$1978 \pm 104$	TIAN	05	BELL $e^+e^- \approx \gamma(4S)$

### $A_{CP}(K_S^0\pi^+\pi^-)$ in $D^0, \bar{D}^0 \rightarrow K_S^0\pi^+\pi^-$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>-0.009 \pm 0.021^{+0.016}_{-0.057}</math></b>	4854	144	ASNER	04A CLEO $e^+e^- \approx 10$ GeV

144 This is the overall result of ASNER 04A;  $CP$ -violating limits are also given for each of the 10 resonant submodes found in an amplitude analysis of the  $D^0$  and  $\bar{D}^0 \rightarrow K_S^0\pi^+\pi^-$  Dalitz plots. These limits range from  $< 3.5 \times 10^{-4}$  to  $28.4 \times 10^{-4}$  at 95% CL.

### $A_{CP}(K^\pm\pi^\mp\pi^+\pi^-)$ in $D^0 \rightarrow K^+\pi^-\pi^+\pi^-, \bar{D}^0 \rightarrow K^-\pi^+\pi^+\pi^-$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>-0.018 \pm 0.044</math></b>	$1721 \pm 75$	TIAN	05	BELL $e^+e^- \approx \gamma(4S)$

### $A_{CP}(K^+K^-\pi^+\pi^-)$ in $D^0, \bar{D}^0 \rightarrow K^+K^-\pi^+\pi^-$

This is the difference between  $D^0$  and  $\bar{D}^0$  partial widths for these modes divided by the sum of the widths. The  $D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow \bar{D}^0\pi^-$ .

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>-0.082 \pm 0.056 \pm 0.047</math></b>	$828 \pm 46$	LINK	05E FOCS	$\gamma A, \bar{E}_\gamma \approx 180$ GeV

## **$D^0$ - $\bar{D}^0$ T-VIOLATING DECAY-RATE ASYMMETRIES**

$D^0$  and  $\bar{D}^0$  are distinguished by the charge of the parent  $D^*$ :  $D^{*+} \rightarrow D^0\pi^+$  and  $D^{*-} \rightarrow D^0\pi^-$ .

### **$A_{T\text{viol}}(K^+K^-\pi^+\pi^-)$ in $D^0, \bar{D}^0 \rightarrow K^+K^-\pi^+\pi^-$**

$C_T \equiv \vec{p}_{K^+} \cdot (\vec{p}_{\pi^+} \times \vec{p}_{\pi^-})$  is a  $T$ -odd correlation of the  $K^+$ ,  $\pi^+$ , and  $\pi^-$  momenta for the  $D^0$ .  $\bar{C}_T \equiv \vec{p}_{K^-} \cdot (\vec{p}_{\pi^-} \times \vec{p}_{\pi^+})$  is the corresponding quantity for the  $\bar{D}^0$ .  $A_T \equiv [\Gamma(C_T > 0) - \Gamma(C_T < 0)] / [\Gamma(C_T > 0) + \Gamma(C_T < 0)]$  would, in the absence of strong phases, test for  $T$  violation in  $D^0$  decays (the  $\Gamma$ 's are partial widths). With  $\bar{A}_T \equiv [\Gamma(-\bar{C}_T > 0) - \Gamma(-\bar{C}_T < 0)] / [\Gamma(-\bar{C}_T > 0) + \Gamma(-\bar{C}_T < 0)]$ , the asymmetry  $A_{T\text{viol}} \equiv \frac{1}{2}(A_T - \bar{A}_T)$  tests for  $T$  violation even with nonzero strong phases.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>+0.010±0.057±0.037</b>	828 ± 46	LINK	05E FOCS	$\gamma$ A, $\bar{E}_\gamma \approx 180$ GeV

## **$D^0$ CPT-VIOLATING DECAY-RATE ASYMMETRIES**

### **$A_{CPT}(K^\mp\pi^\pm)$ in $D^0 \rightarrow K^-\pi^+, \bar{D}^0 \rightarrow K^+\pi^-$**

$A_{CPT}(t)$  is defined in terms of the time-dependent decay probabilities  $P(D^0 \rightarrow K^-\pi^+)$  and  $\bar{P}(\bar{D}^0 \rightarrow K^+\pi^-)$  by  $A_{CPT}(t) = (\bar{P} - P)/(\bar{P} + P)$ . For small mixing parameters  $x \equiv \Delta m/\Gamma$  and  $y \equiv \Delta\Gamma/2\Gamma$  (as is the case), and times  $t$ ,  $A_{CPT}(t)$  reduces to  $[y \operatorname{Re} \xi - x \operatorname{Im} \xi] \Gamma t$ , where  $\xi$  is the  $CPT$ -violating parameter.

The following is actually  $y \operatorname{Re} \xi - x \operatorname{Im} \xi$ .

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.0083±0.0065±0.0041</b>	LINK	03B FOCS	$\gamma$ nucleus, $\bar{E}_\gamma \approx 180$ GeV

## **$D^0 \rightarrow K^*(892)^-\ell^+\nu_\ell$ FORM FACTORS**

### **$r_V \equiv V(0)/A_1(0)$ in $D^0 \rightarrow K^*(892)^-\ell^+\nu_\ell$**

VALUE	DOCUMENT ID	TECN	COMMENT
<b>1.71±0.68±0.34</b>	LINK	05B FOCS	$K^*(892)^-\mu^+\nu_\mu$

### **$r_2 \equiv A_2(0)/A_1(0)$ in $D^0 \rightarrow K^*(892)^-\ell^+\nu_\ell$**

VALUE	DOCUMENT ID	TECN	COMMENT
<b>0.91±0.37±0.10</b>	LINK	05B FOCS	$K^*(892)^-\mu^+\nu_\mu$

## **$D^0$ REFERENCES**

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ASNER	04A	PR D70 091101R	D.M. Asner <i>et al.</i>	(CLEO Collab.)
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AITALA	98	PR D57 13	E.M. Aitala <i>et al.</i>	(FNAL E791 Collab.)
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Also		PRL 77 2147 (erratum)	A. Freyberger <i>et al.</i>	(CLEO Collab.)
KUBOTA	96B	PR D54 2994	Y. Kubota <i>et al.</i>	(CLEO Collab.)

ADAMOVICH	95	PL B353 563	M.I. Adamovich <i>et al.</i>	(CERN BEATRICE Collab.)
BARTEL	95	PR D52 4860	J.E. Bartelt <i>et al.</i>	(CLEO Collab.)
BUTLER	95	PR D52 2656	F. Butler <i>et al.</i>	(CLEO Collab.)
FRAEBETTI	95C	PL B354 486	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRAEBETTI	95G	PL B364 127	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
KODAMA	95	PL B345 85	K. Kodama <i>et al.</i>	(FNAL E653 Collab.)
ALBRECHT	94	PL B324 249	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALBRECHT	94F	PL B340 125	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALBRECHT	94I	ZPHY C64 375	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
CINABRO	94	PRL 72 1406	D. Cinabro <i>et al.</i>	(CLEO Collab.)
FRAEBETTI	94C	PL B321 295	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRAEBETTI	94D	PL B323 459	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRAEBETTI	94G	PL B331 217	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRAEBETTI	94I	PR D50 R2953	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRAEBETTI	94J	PL B340 254	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
KODAMA	94	PL B336 605	K. Kodama <i>et al.</i>	(FNAL E653 Collab.)
MISHRA	94	PR D50 R9	C.S. Mishra <i>et al.</i>	(FNAL E789 Collab.)
AKERIB	93	PRL 71 3070	D.S. Akerib <i>et al.</i>	(CLEO Collab.)
ALBRECHT	93D	PL B308 435	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ANJOS	93	PR D48 56	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
BEAN	93C	PL B317 647	A. Bean <i>et al.</i>	(CLEO Collab.)
FRAEBETTI	93I	PL B315 203	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
KODAMA	93B	PL B313 260	K. Kodama <i>et al.</i>	(FNAL E653 Collab.)
PROCARIO	93B	PR D48 4007	M. Procario <i>et al.</i>	(CLEO Collab.)
SELEN	93	PRL 71 1973	M.A. Selen <i>et al.</i>	(CLEO Collab.)
ADAMOVICH	92	PL B280 163	M.I. Adamovich <i>et al.</i>	(CERN WA82 Collab.)
ALBRECHT	92P	ZPHY C56 7	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ANJOS	92B	PR D46 R1	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
ANJOS	92C	PR D46 1941	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
BARLAG	92C	ZPHY C55 383	S. Barlag <i>et al.</i>	(ACCMOR Collab.)
Also		ZPHY C48 29	S. Barlag <i>et al.</i>	(ACCMOR Collab.)
COFFMAN	92B	PR D45 2196	D.M. Coffman <i>et al.</i>	(Mark III Collab.)
Also		PRL 64 2615	J. Adler <i>et al.</i>	(Mark III Collab.)
FRAEBETTI	92	PL B281 167	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
FRAEBETTI	92B	PL B286 195	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
ALVAREZ	91B	ZPHY C50 11	M.P. Alvarez <i>et al.</i>	(CERN NA14/2 Collab.)
AMMAR	91	PR D44 3383	R. Ammar <i>et al.</i>	(CLEO Collab.)
ANJOS	91	PR D43 R635	J.C. Anjos <i>et al.</i>	(FNAL-TPS Collab.)
ANJOS	91D	PR D44 R3371	J.C. Anjos <i>et al.</i>	(FNAL-TPS Collab.)
BAI	91	PRL 66 1011	Z. Bai <i>et al.</i>	(Mark III Collab.)
COFFMAN	91	PL B263 135	D.M. Coffman <i>et al.</i>	(Mark III Collab.)
CRAWFORD	91B	PR D44 3394	G. Crawford <i>et al.</i>	(CLEO Collab.)
DECAMP	91J	PL B266 218	D. Decamp <i>et al.</i>	(ALEPH Collab.)
FRAEBETTI	91	PL B263 584	P.L. Frabetti <i>et al.</i>	(FNAL E687 Collab.)
KINOSHITA	91	PR D43 2836	K. Kinoshita <i>et al.</i>	(CLEO Collab.)
KODAMA	91	PRL 66 1819	K. Kodama <i>et al.</i>	(FNAL E653 Collab.)
ALBRECHT	90C	ZPHY C46 9	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALEXANDER	90	PRL 65 1184	J. Alexander <i>et al.</i>	(CLEO Collab.)
ALEXANDER	90B	PRL 65 1531	J. Alexander <i>et al.</i>	(CLEO Collab.)
ALVAREZ	90	ZPHY C47 539	M.P. Alvarez <i>et al.</i>	(CERN NA14/2 Collab.)
ANJOS	90D	PR D42 2414	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
BARLAG	90C	ZPHY C46 563	S. Barlag <i>et al.</i>	(ACCMOR Collab.)
ADLER	89	PRL 62 1821	J. Adler <i>et al.</i>	(Mark III Collab.)
ADLER	89C	PR D40 906	J. Adler <i>et al.</i>	(Mark III Collab.)
ALBRECHT	89D	ZPHY C43 181	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ANJOS	89F	PRL 62 1587	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
ABACHI	88	PL B205 411	S. Abachi <i>et al.</i>	(HRS Collab.)
ADLER	88	PR D37 2023	J. Adler <i>et al.</i>	(Mark III Collab.)
ADLER	88C	PRL 60 89	J. Adler <i>et al.</i>	(Mark III Collab.)
ALBRECHT	88G	PL B209 380	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALBRECHT	88I	PL B210 267	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ANJOS	88C	PRL 60 1239	J.C. Anjos <i>et al.</i>	(FNAL E691 Collab.)
BORTOLETTO	88	PR D37 1719	D. Bortoletto <i>et al.</i>	(CLEO Collab.)
Also		PR D39 1471 (erratum)	D. Bortoletto <i>et al.</i>	(CLEO Collab.)
HAAS	88	PRL 60 1614	P. Haas <i>et al.</i>	(CLEO Collab.)
RAAB	88	PR D37 2391	J.R. Raab <i>et al.</i>	(FNAL E691 Collab.)
ADAMOVICH	87	EPL 4 887	M.I. Adamovich <i>et al.</i>	(Photon Emulsion Collab.)
ADLER	87	PL B196 107	J. Adler <i>et al.</i>	(Mark III Collab.)
AGUILAR-...	87E	ZPHY C36 551	M. Aguilar-Benitez <i>et al.</i>	(LEBC-EHS Collab.)
Also		ZPHY C40 321	M. Aguilar-Benitez <i>et al.</i>	(LEBC-EHS Collab.)

AGUILAR-	87F	ZPHY C36 559	M. Aguilar-Benitez <i>et al.</i>	(LEBC-EHS Collab.)
Also		ZPHY C38 520 (erratum)	M. Aguilar-Benitez <i>et al.</i>	(LEBC-EHS Collab.)
ALBRECHT	87K	PL B199 447	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
BARLAG	87B	ZPHY C37 17	S. Barlag <i>et al.</i>	(ACCMOR Collab.)
BECKER	87C	PL B193 147	J.J. Becker <i>et al.</i>	(Mark III Collab.)
Also		PL B198 590 (erratum)	J.J. Becker <i>et al.</i>	(Mark III Collab.)
PALKA	87	PL B189 238	H. Palka <i>et al.</i>	(ACCMOR Collab.)
RILES	87	PR D35 2914	K. Riles <i>et al.</i>	(Mark II Collab.)
BAILEY	86	ZPHY C30 51	R. Bailey <i>et al.</i>	(ACCMOR Collab.)
BEBEK	86	PRL 56 1893	C. Bebek <i>et al.</i>	(CLEO Collab.)
LOUIS	86	PRL 56 1027	W.C. Louis <i>et al.</i>	(PRIN, CHIC, ISU)
ALBRECHT	85B	PL 158B 525	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
ALBRECHT	85F	PL 150B 235	H. Albrecht <i>et al.</i>	(ARGUS Collab.)
AUBERT	85	PL 155B 461	J.J. Aubert <i>et al.</i>	(EMC Collab.)
BALTRUSAIT...	85B	PRL 54 1976	R.M. Baltrusaitis <i>et al.</i>	(Mark III Collab.)
BALTRUSAIT...	85E	PRL 55 150	R.M. Baltrusaitis <i>et al.</i>	(Mark III Collab.)
BENVENUTI	85	PL 158B 531	A.C. Benvenuti <i>et al.</i>	(BCDMS Collab.)
ADAMOVICH	84B	PL 140B 123	M.I. Adamovich <i>et al.</i>	(CERN WA58 Collab.)
DERRICK	84	PR 53 1971	M. Derrick <i>et al.</i>	(HRS Collab.)
SUMMERS	84	PRL 52 410	D.J. Summers <i>et al.</i>	(UCSB, CARL, COLO+)
BAILEY	83B	PL 132B 237	R. Bailey <i>et al.</i>	(ACCMOR Collab.)
BODEK	82	PL 113B 82	A. Bodek <i>et al.</i>	(ROCH, CIT, CHIC, FNAL+)
FIORINO	81	LNC 30 166	A. Fiorino <i>et al.</i>	
SCHINDLER	81	PR D24 78	R.H. Schindler <i>et al.</i>	(Mark II Collab.)
TRILLING	81	PRPL 75 57	G.H. Trilling	(LBL, UCB) J
ASTON	80E	PL 94B 113	D. Aston <i>et al.</i>	(BONN, CERN, EPOL, GLAS+)
AVERY	80	PRL 44 1309	P. Avery <i>et al.</i>	(ILL, FNAL, COLU)
ZHOLENTZ	80	PL 96B 214	A.A. Zholents <i>et al.</i>	(NOVO)
Also		SJNP 34 814	A.A. Zholents <i>et al.</i>	(NOVO)
		Translated from YAF 34	1471.	
ABRAMS	79D	PRL 43 481	G.S. Abrams <i>et al.</i>	(Mark II Collab.)
ATIYA	79	PRL 43 414	M.S. Atiya <i>et al.</i>	(COLU, ILL, FNAL)
BALTAY	78C	PRL 41 73	C. Baltay <i>et al.</i>	(COLU, BNL)
VUILLEMEN	78	PRL 41 1149	V. Vuillemin <i>et al.</i>	(Mark I Collab.)
GOLDHABER	77	PL 69B 503	G. Goldhaber <i>et al.</i>	(Mark I Collab.)
PERUZZI	77	PRL 39 1301	I. Peruzzi <i>et al.</i>	(Mark I Collab.)
PICCOLO	77	PL 70B 260	M. Piccolo <i>et al.</i>	(Mark I Collab.)
GOLDHABER	76	PRL 37 255	G. Goldhaber <i>et al.</i>	(Mark I Collab.)

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